



THE UNIVERSITY
of EDINBURGH

Connecting scales in the SMEFT at the LHC and future colliders

DAMTP seminar

Based on arXiv: 2502.20453, 2504.05974

In collaboration with Luca Mantani, Juan
Rojo, Alejo Rossia, Eleni Vryonidou

Jaco ter Hoeve

02/05/25

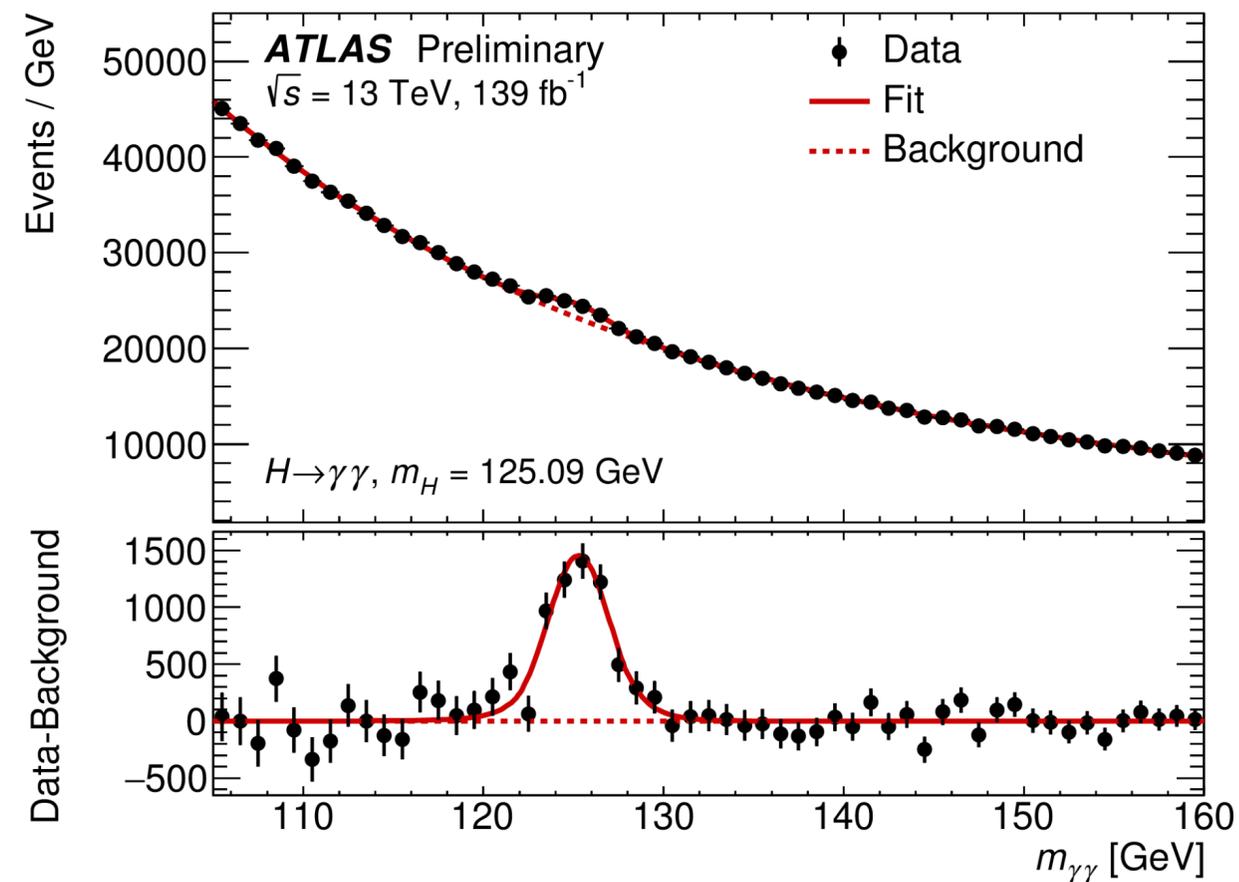


This talk

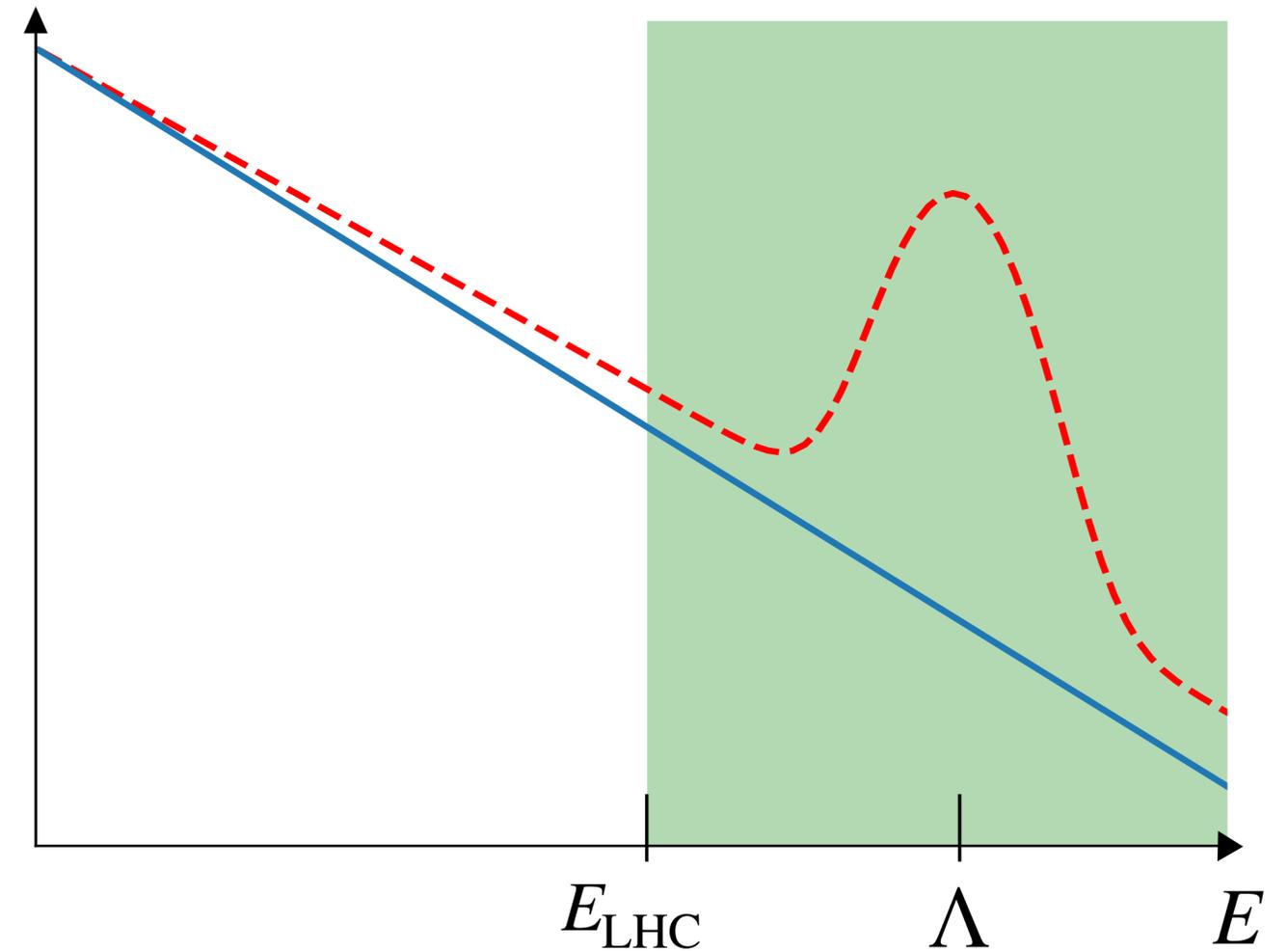
- A brief introduction to the SMEFT
- Renormalisation Group Evolution in global SMEFT fits
- The SMEFT at future colliders
- The Higgs self-coupling at FCC-ee
- Summary

How to look for New Physics

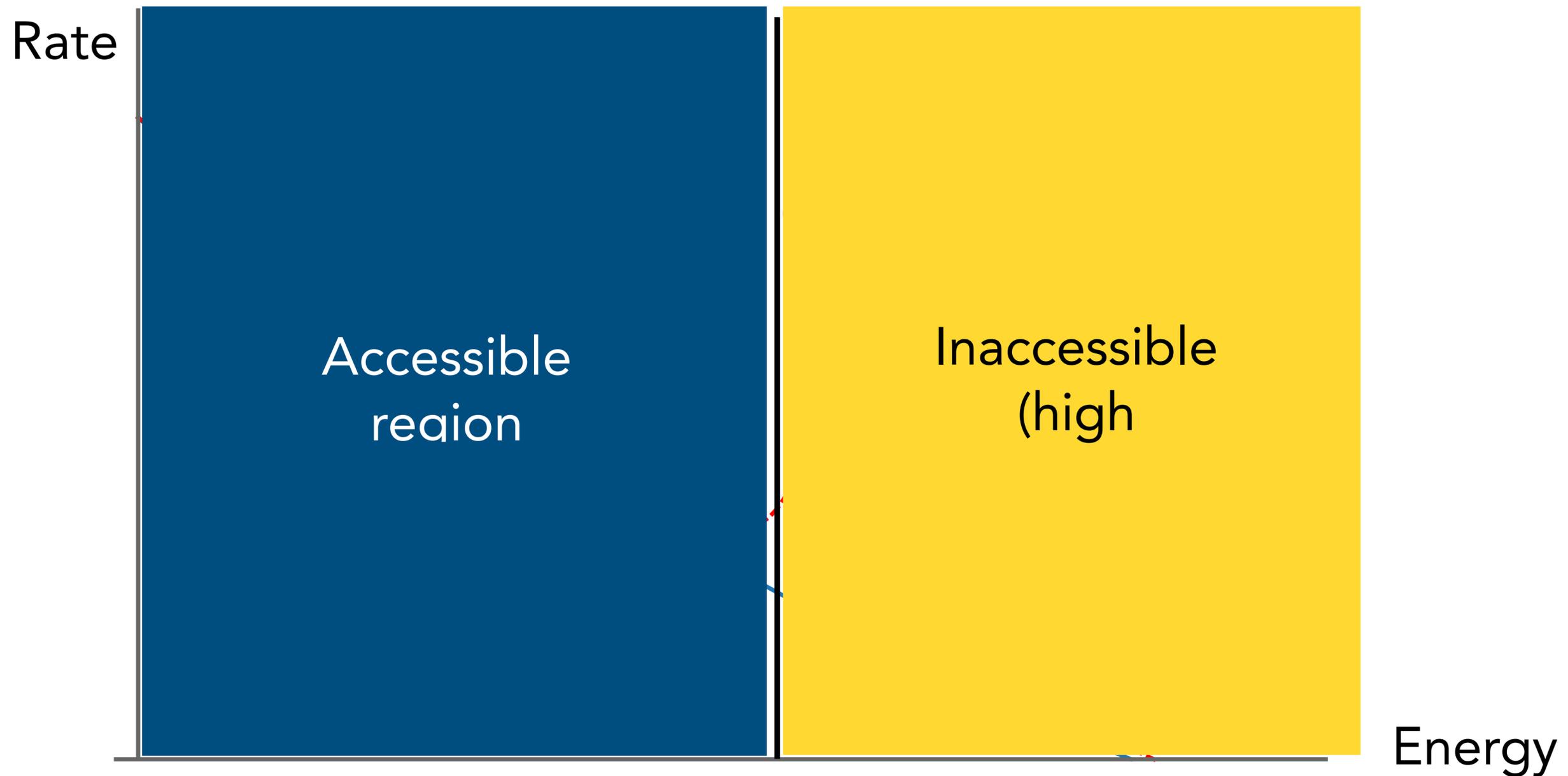
1. Directly: bump hunting



2. Indirectly: tiny deviations in tails



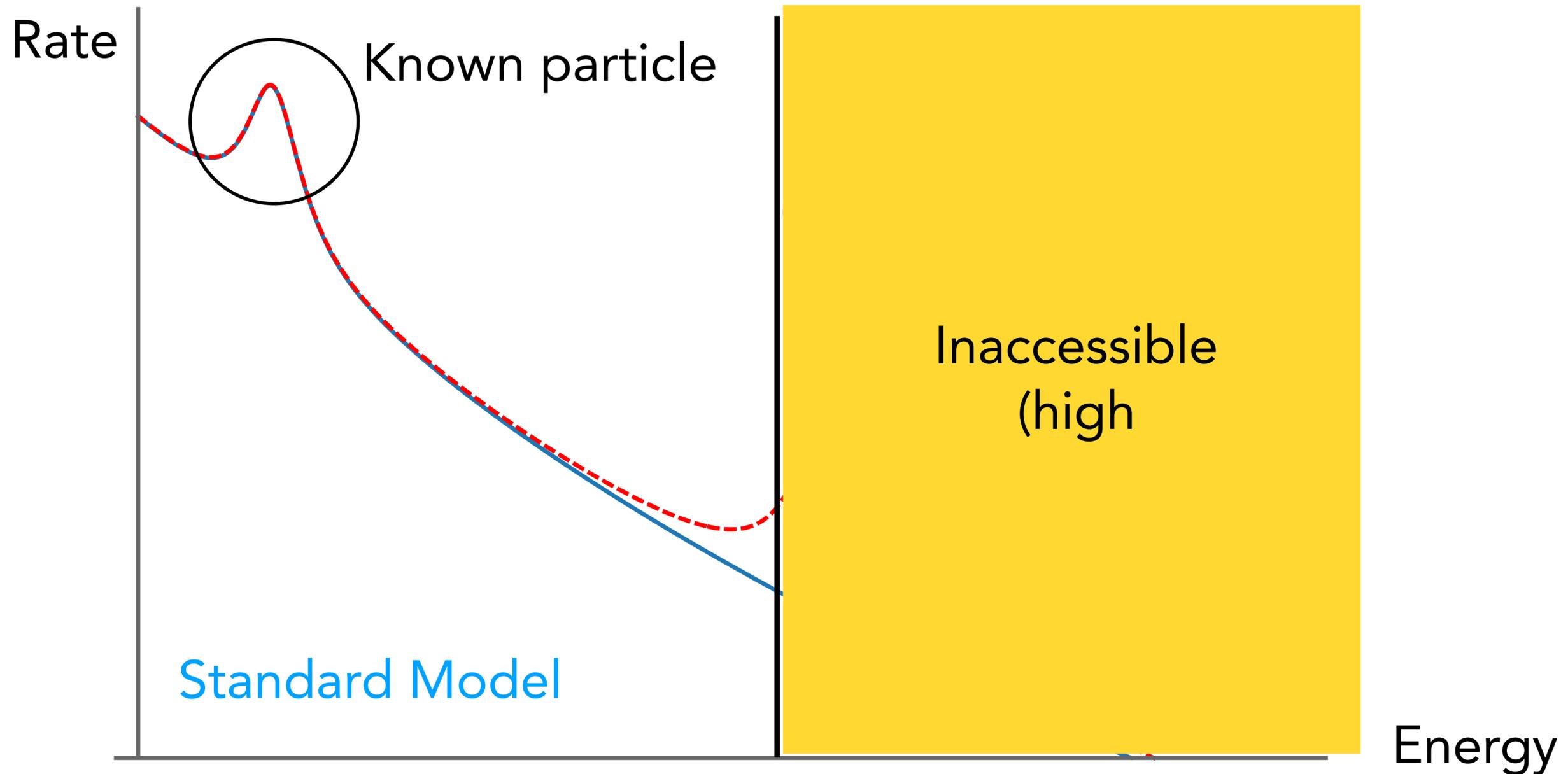
The Standard Model Effective Field Theory



*Adapted from
E. Vryonidou*

**The SMEFT: the way to probe new physics beyond
the direct collider energy reach**

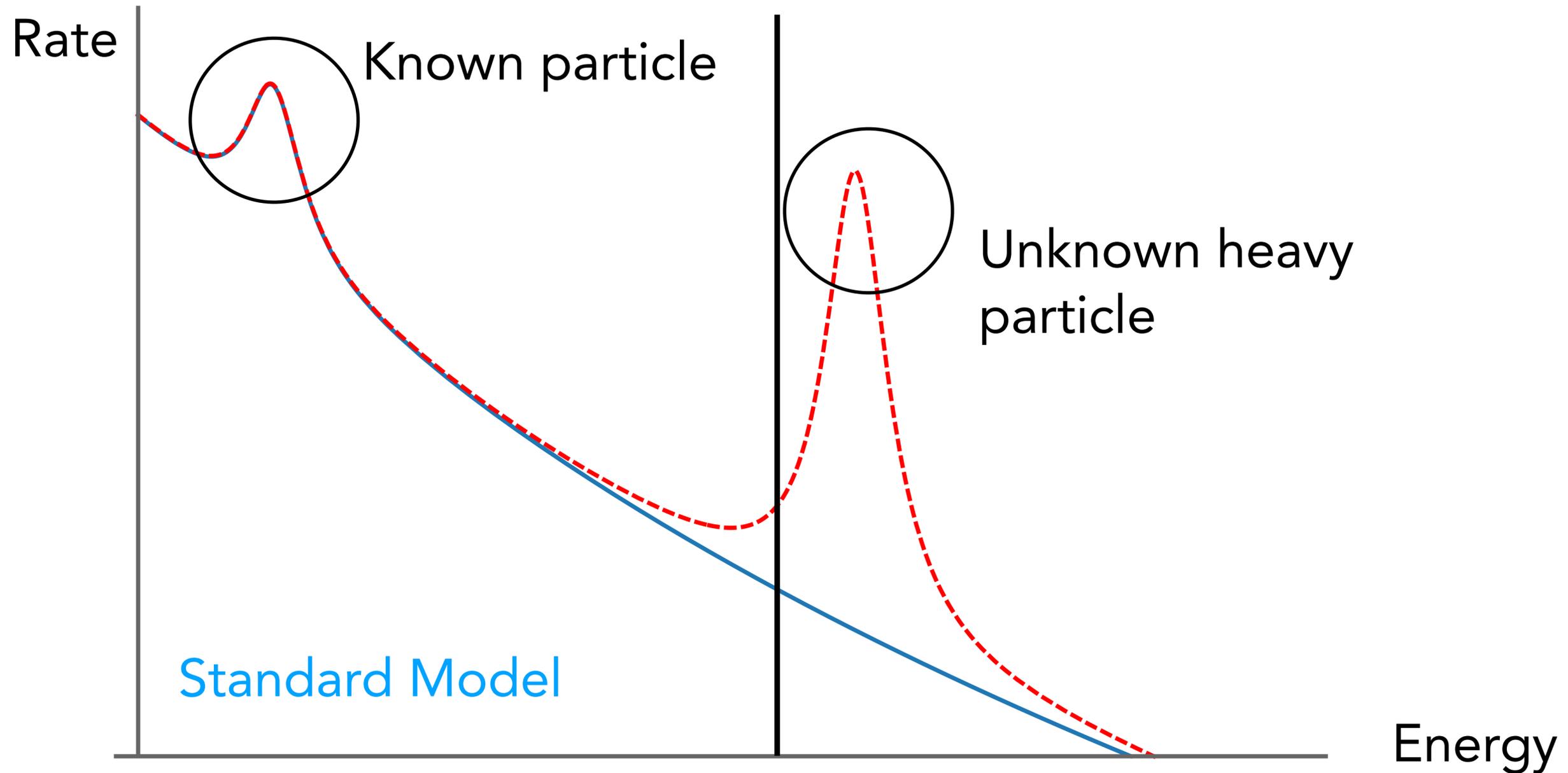
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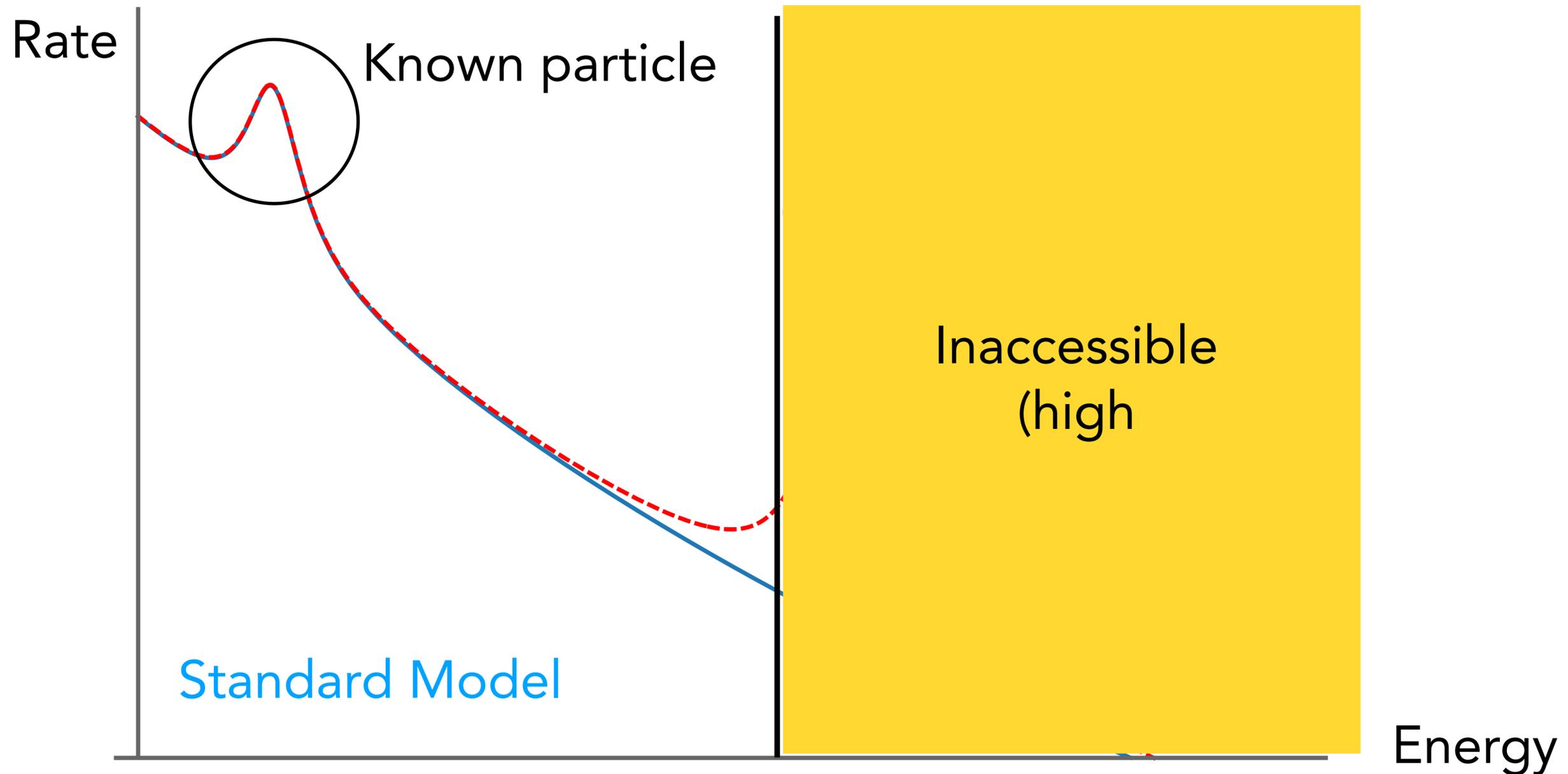
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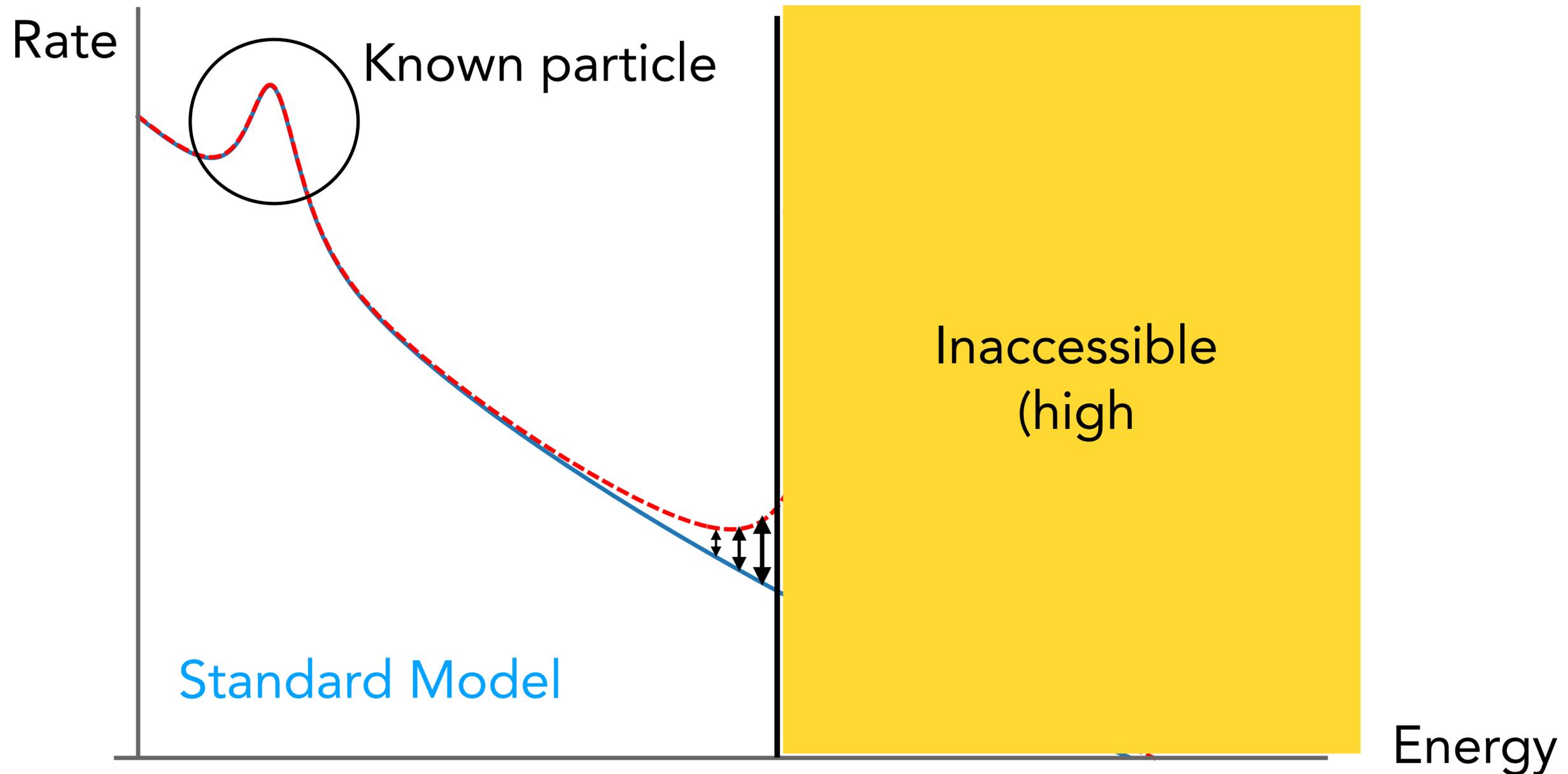
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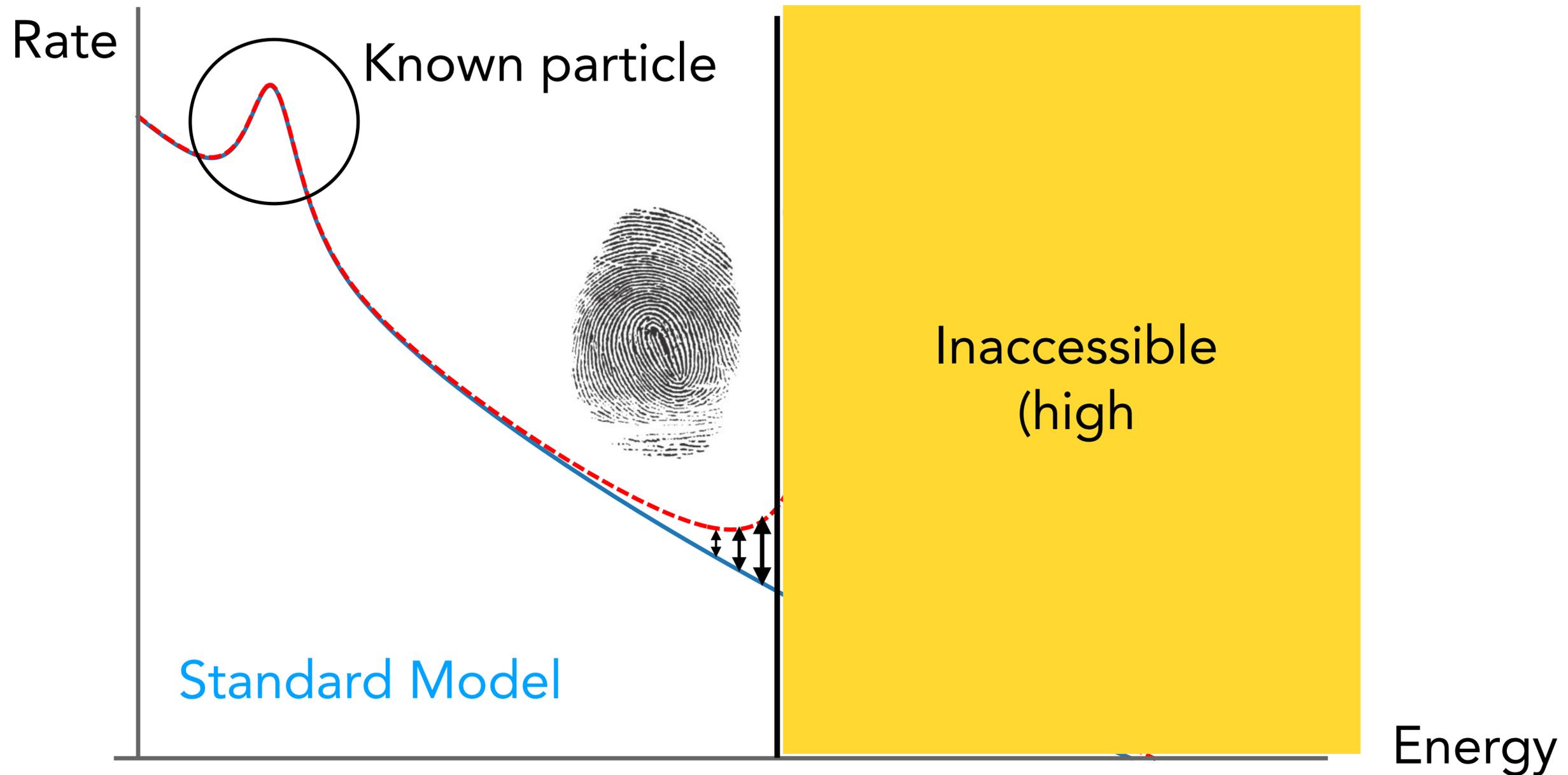
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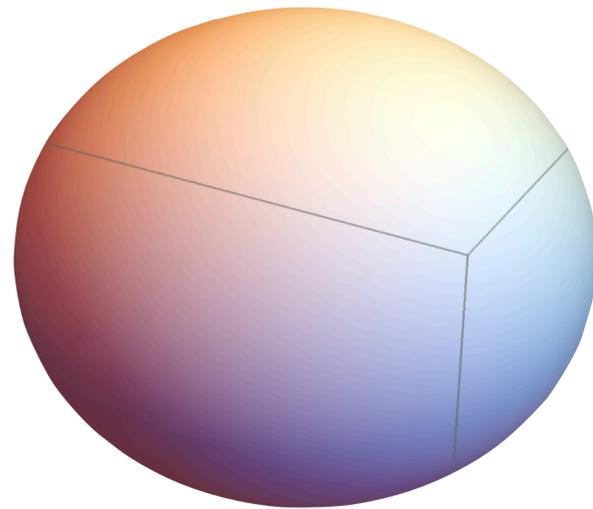


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**The SMEFT: the way to probe new physics beyond
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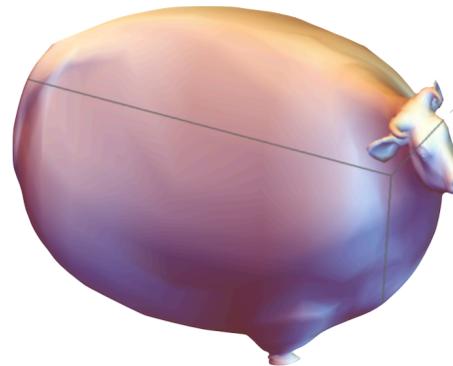
The SMEFT philosophy

The SM



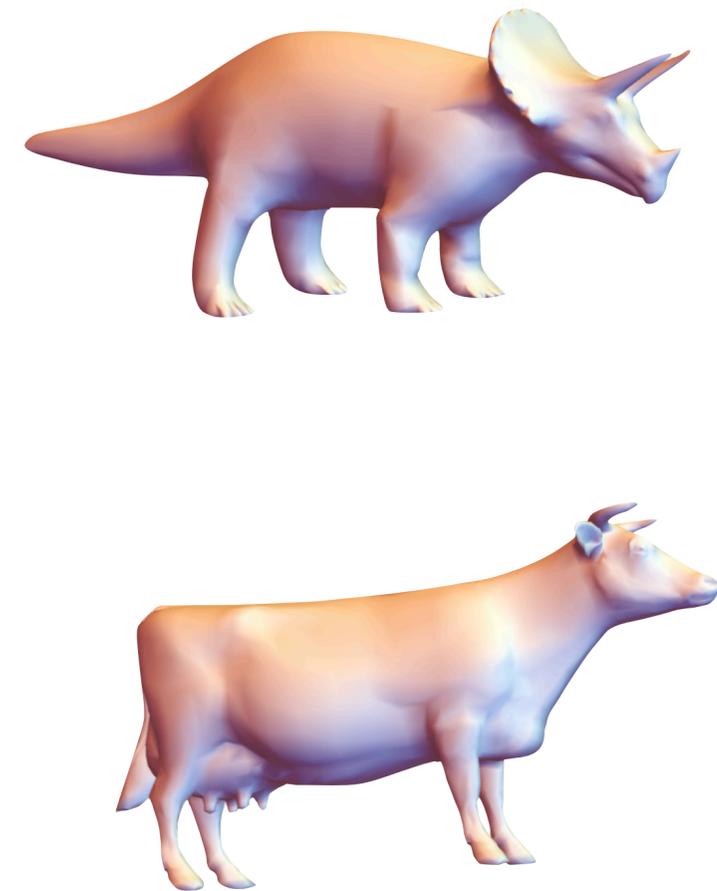
+

$c \cdot$



\approx

New Physics



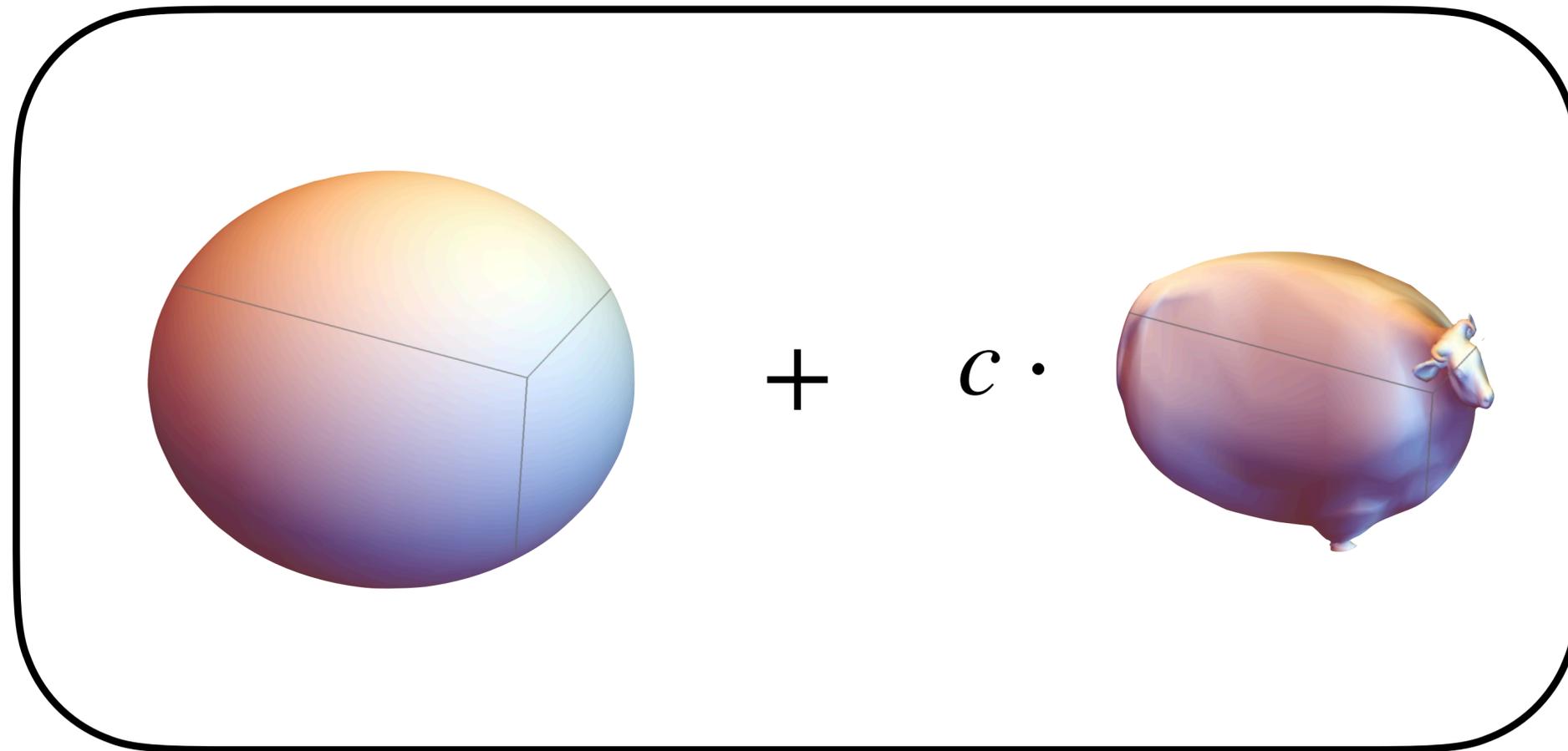
Goal: find the value of c , and precisely!

The SMEFT philosophy

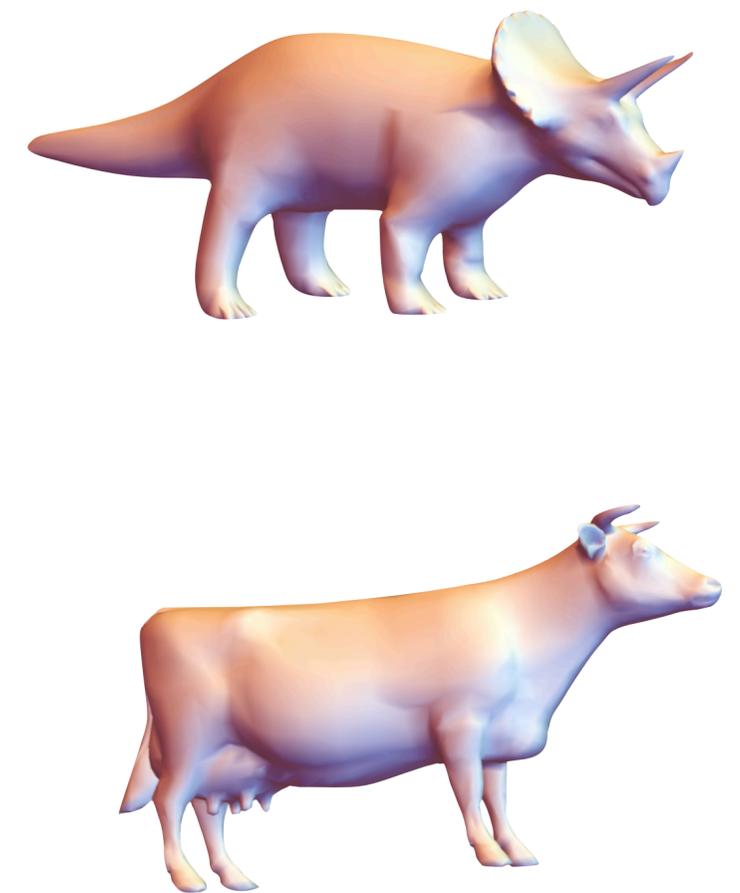
The SM

EFT

New Physics



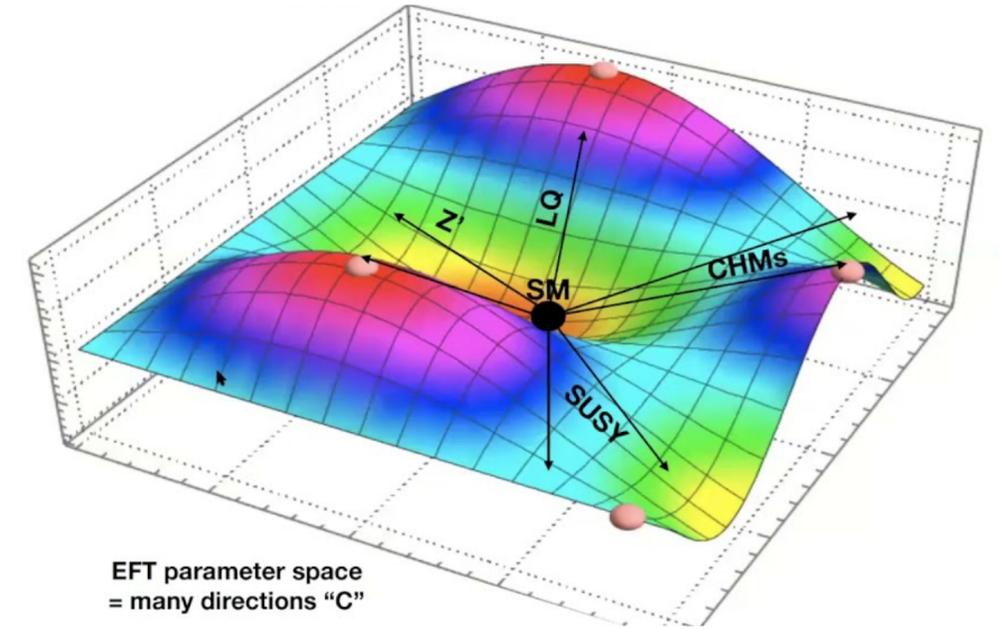
\approx



Goal: find the value of c , and precisely!

SMEFT requirements

$$\mathcal{L}_{\text{SMEFT}} = \mathcal{L}_{\text{SM}} + \sum_i^{N_{d5}} \frac{c_i}{\Lambda} \mathcal{O}_i^{(5)} + \sum_i^{N_{d6}} \frac{c_i}{\Lambda^2} \mathcal{O}_i^{(6)} + \dots$$



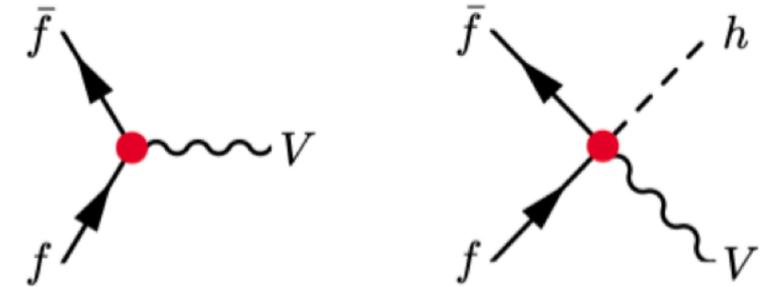
- The SMEFT parameterises the **theory space around the SM**
- Defines a perfectly valid QFT that is **renormalisable** order by order in Λ
- Operators are constructed out of **SM fields only**
- Respects the SM **symmetries**
- Forms a **complete basis** at any given mass dimension

The Warsaw basis

arXiv: 1008.4884

X^3		φ^6 and $\varphi^4 D^2$		$\psi^2 \varphi^3$	
Q_G	$f^{ABC} G_\mu^{A\nu} G_\nu^{B\rho} G_\rho^{C\mu}$	Q_φ	$(\varphi^\dagger \varphi)^3$	$Q_{e\varphi}$	$(\varphi^\dagger \varphi)(\bar{l}_p e_r \varphi)$
$Q_{\tilde{G}}$	$f^{ABC} \tilde{G}_\mu^{A\nu} G_\nu^{B\rho} G_\rho^{C\mu}$	$Q_{\varphi\Box}$	$(\varphi^\dagger \varphi)\Box(\varphi^\dagger \varphi)$	$Q_{u\varphi}$	$(\varphi^\dagger \varphi)(\bar{q}_p u_r \tilde{\varphi})$
Q_W	$\varepsilon^{IJK} W_\mu^{I\nu} W_\nu^{J\rho} W_\rho^{K\mu}$	$Q_{\varphi D}$	$(\varphi^\dagger D^\mu \varphi)^* (\varphi^\dagger D_\mu \varphi)$	$Q_{d\varphi}$	$(\varphi^\dagger \varphi)(\bar{q}_p d_r \varphi)$
$Q_{\tilde{W}}$	$\varepsilon^{IJK} \tilde{W}_\mu^{I\nu} W_\nu^{J\rho} W_\rho^{K\mu}$				
$X^2 \varphi^2$		$\psi^2 X \varphi$		$\psi^2 \varphi^2 D$	
$Q_{\varphi G}$	$\varphi^\dagger \varphi G_{\mu\nu}^A G^{A\mu\nu}$	Q_{eW}	$(\bar{l}_p \sigma^{\mu\nu} e_r) \tau^I \varphi W_{\mu\nu}^I$	$Q_{\varphi l}^{(1)}$	$(\varphi^\dagger i \overleftrightarrow{D}_\mu \varphi)(\bar{l}_p \gamma^\mu l_r)$
$Q_{\varphi \tilde{G}}$	$\varphi^\dagger \varphi \tilde{G}_{\mu\nu}^A G^{A\mu\nu}$	Q_{eB}	$(\bar{l}_p \sigma^{\mu\nu} e_r) \varphi B_{\mu\nu}$	$Q_{\varphi l}^{(3)}$	$(\varphi^\dagger i \overleftrightarrow{D}_\mu^I \varphi)(\bar{l}_p \tau^I \gamma^\mu l_r)$
$Q_{\varphi W}$	$\varphi^\dagger \varphi W_{\mu\nu}^I W^{I\mu\nu}$	Q_{uG}	$(\bar{q}_p \sigma^{\mu\nu} T^A u_r) \tilde{\varphi} G_{\mu\nu}^A$	$Q_{\varphi e}$	$(\varphi^\dagger i \overleftrightarrow{D}_\mu \varphi)(\bar{e}_p \gamma^\mu e_r)$
$Q_{\varphi \tilde{W}}$	$\varphi^\dagger \varphi \tilde{W}_{\mu\nu}^I W^{I\mu\nu}$	Q_{uW}	$(\bar{q}_p \sigma^{\mu\nu} u_r) \tau^I \tilde{\varphi} W_{\mu\nu}^I$	$Q_{\varphi q}^{(1)}$	$(\varphi^\dagger i \overleftrightarrow{D}_\mu \varphi)(\bar{q}_p \gamma^\mu q_r)$
$Q_{\varphi B}$	$\varphi^\dagger \varphi B_{\mu\nu} B^{\mu\nu}$	Q_{uB}	$(\bar{q}_p \sigma^{\mu\nu} u_r) \tilde{\varphi} B_{\mu\nu}$	$Q_{\varphi q}^{(3)}$	$(\varphi^\dagger i \overleftrightarrow{D}_\mu^I \varphi)(\bar{q}_p \tau^I \gamma^\mu q_r)$
$Q_{\varphi \tilde{B}}$	$\varphi^\dagger \varphi \tilde{B}_{\mu\nu} B^{\mu\nu}$	Q_{dG}	$(\bar{q}_p \sigma^{\mu\nu} T^A d_r) \varphi G_{\mu\nu}^A$	$Q_{\varphi u}$	$(\varphi^\dagger i \overleftrightarrow{D}_\mu \varphi)(\bar{u}_p \gamma^\mu u_r)$
$Q_{\varphi WB}$	$\varphi^\dagger \tau^I \varphi W_{\mu\nu}^I B^{\mu\nu}$	Q_{dW}	$(\bar{q}_p \sigma^{\mu\nu} d_r) \tau^I \varphi W_{\mu\nu}^I$	$Q_{\varphi d}$	$(\varphi^\dagger i \overleftrightarrow{D}_\mu \varphi)(\bar{d}_p \gamma^\mu d_r)$
$Q_{\varphi \tilde{W}B}$	$\varphi^\dagger \tau^I \varphi \tilde{W}_{\mu\nu}^I B^{\mu\nu}$	Q_{dB}	$(\bar{q}_p \sigma^{\mu\nu} d_r) \varphi B_{\mu\nu}$	$Q_{\varphi ud}$	$i(\tilde{\varphi}^\dagger D_\mu \varphi)(\bar{u}_p \gamma^\mu d_r)$

Current operators



Adapted from
K. Mimasu

- Shift SM ffV couplings
- $ffVh$ contact interactions

Contact interactions grow with energy and thus provide a sensitive probe in the TeV region

The Warsaw basis

arXiv: 1008.4884

Adapted from
K. Mimasu

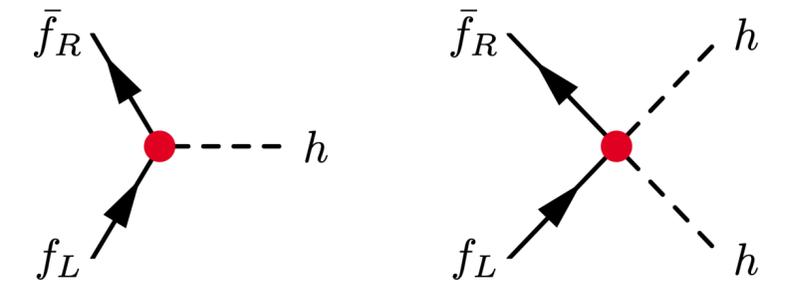
X^3		φ^6 and $\varphi^4 D^2$		$\psi^2 \varphi^3$	
Q_G	$f^{ABC} G_\mu^{A\nu} G_\nu^{B\rho} G_\rho^{C\mu}$	Q_φ	$(\varphi^\dagger \varphi)^3$	$Q_{e\varphi}$	$(\varphi^\dagger \varphi)(\bar{l}_p e_r \varphi)$
$Q_{\tilde{G}}$	$f^{ABC} \tilde{G}_\mu^{A\nu} G_\nu^{B\rho} G_\rho^{C\mu}$	$Q_{\varphi\Box}$	$(\varphi^\dagger \varphi)\Box(\varphi^\dagger \varphi)$	$Q_{u\varphi}$	$(\varphi^\dagger \varphi)(\bar{q}_p u_r \tilde{\varphi})$
Q_W	$\varepsilon^{IJK} W_\mu^{I\nu} W_\nu^{J\rho} W_\rho^{K\mu}$	$Q_{\varphi D}$	$(\varphi^\dagger D^\mu \varphi)^* (\varphi^\dagger D_\mu \varphi)$	$Q_{d\varphi}$	$(\varphi^\dagger \varphi)(\bar{q}_p d_r \varphi)$
$Q_{\tilde{W}}$	$\varepsilon^{IJK} \tilde{W}_\mu^{I\nu} W_\nu^{J\rho} W_\rho^{K\mu}$				
$X^2 \varphi^2$		$\psi^2 X \varphi$		$\psi^2 \varphi^2 D$	
$Q_{\varphi G}$	$\varphi^\dagger \varphi G_{\mu\nu}^A G^{A\mu\nu}$	Q_{eW}	$(\bar{l}_p \sigma^{\mu\nu} e_r) \tau^I \varphi W_{\mu\nu}^I$	$Q_{\varphi l}^{(1)}$	$(\varphi^\dagger i \overleftrightarrow{D}_\mu \varphi)(\bar{l}_p \gamma^\mu l_r)$
$Q_{\varphi \tilde{G}}$	$\varphi^\dagger \varphi \tilde{G}_{\mu\nu}^A G^{A\mu\nu}$	Q_{eB}	$(\bar{l}_p \sigma^{\mu\nu} e_r) \varphi B_{\mu\nu}$	$Q_{\varphi l}^{(3)}$	$(\varphi^\dagger i \overleftrightarrow{D}_\mu^I \varphi)(\bar{l}_p \tau^I \gamma^\mu l_r)$
$Q_{\varphi W}$	$\varphi^\dagger \varphi W_{\mu\nu}^I W^{I\mu\nu}$	Q_{uG}	$(\bar{q}_p \sigma^{\mu\nu} T^A u_r) \tilde{\varphi} G_{\mu\nu}^A$	$Q_{\varphi e}$	$(\varphi^\dagger i \overleftrightarrow{D}_\mu \varphi)(\bar{e}_p \gamma^\mu e_r)$
$Q_{\varphi \tilde{W}}$	$\varphi^\dagger \varphi \tilde{W}_{\mu\nu}^I W^{I\mu\nu}$	Q_{uW}	$(\bar{q}_p \sigma^{\mu\nu} u_r) \tau^I \tilde{\varphi} W_{\mu\nu}^I$	$Q_{\varphi q}^{(1)}$	$(\varphi^\dagger i \overleftrightarrow{D}_\mu \varphi)(\bar{q}_p \gamma^\mu q_r)$
$Q_{\varphi B}$	$\varphi^\dagger \varphi B_{\mu\nu} B^{\mu\nu}$	Q_{uB}	$(\bar{q}_p \sigma^{\mu\nu} u_r) \tilde{\varphi} B_{\mu\nu}$	$Q_{\varphi q}^{(3)}$	$(\varphi^\dagger i \overleftrightarrow{D}_\mu^I \varphi)(\bar{q}_p \tau^I \gamma^\mu q_r)$
$Q_{\varphi \tilde{B}}$	$\varphi^\dagger \varphi \tilde{B}_{\mu\nu} B^{\mu\nu}$	Q_{dG}	$(\bar{q}_p \sigma^{\mu\nu} T^A d_r) \varphi G_{\mu\nu}^A$	$Q_{\varphi u}$	$(\varphi^\dagger i \overleftrightarrow{D}_\mu \varphi)(\bar{u}_p \gamma^\mu u_r)$
$Q_{\varphi WB}$	$\varphi^\dagger \tau^I \varphi W_{\mu\nu}^I B^{\mu\nu}$	Q_{dW}	$(\bar{q}_p \sigma^{\mu\nu} d_r) \tau^I \varphi W_{\mu\nu}^I$	$Q_{\varphi d}$	$(\varphi^\dagger i \overleftrightarrow{D}_\mu \varphi)(\bar{d}_p \gamma^\mu d_r)$
$Q_{\varphi \tilde{W}B}$	$\varphi^\dagger \tau^I \varphi \tilde{W}_{\mu\nu}^I B^{\mu\nu}$	Q_{dB}	$(\bar{q}_p \sigma^{\mu\nu} d_r) \varphi B_{\mu\nu}$	$Q_{\varphi ud}$	$i(\tilde{\varphi}^\dagger D_\mu \varphi)(\bar{u}_p \gamma^\mu d_r)$

Contact interactions grow with energy and thus provide a sensitive probe in the TeV region

The Warsaw basis

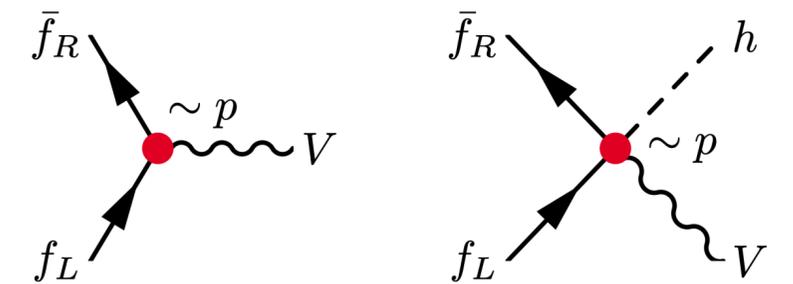
X^3		φ^6 and $\varphi^4 D^2$		$\psi^2 \varphi^3$	
Q_G	$f^{ABC} G_\mu^{A\nu} G_\nu^{B\rho} G_\rho^{C\mu}$	Q_φ	$(\varphi^\dagger \varphi)^3$	$Q_{e\varphi}$	$(\varphi^\dagger \varphi)(\bar{l}_p e_r \varphi)$
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$Q_{\varphi \tilde{G}}$	$\varphi^\dagger \varphi \tilde{G}_{\mu\nu}^A G^{A\mu\nu}$	Q_{eB}	$(\bar{l}_p \sigma^{\mu\nu} e_r) \varphi B_{\mu\nu}$	$Q_{\varphi l}^{(3)}$	$(\varphi^\dagger i \overleftrightarrow{D}_\mu^I \varphi)(\bar{l}_p \tau^I \gamma^\mu l_r)$
$Q_{\varphi W}$	$\varphi^\dagger \varphi W_{\mu\nu}^I W^{I\mu\nu}$	Q_{uG}	$(\bar{q}_p \sigma^{\mu\nu} T^A u_r) \tilde{\varphi} G_{\mu\nu}^A$	$Q_{\varphi e}$	$(\varphi^\dagger i \overleftrightarrow{D}_\mu \varphi)(\bar{e}_p \gamma^\mu e_r)$
$Q_{\varphi \tilde{W}}$	$\varphi^\dagger \varphi \tilde{W}_{\mu\nu}^I W^{I\mu\nu}$	Q_{uW}	$(\bar{q}_p \sigma^{\mu\nu} u_r) \tau^I \tilde{\varphi} W_{\mu\nu}^I$	$Q_{\varphi q}^{(1)}$	$(\varphi^\dagger i \overleftrightarrow{D}_\mu \varphi)(\bar{q}_p \gamma^\mu q_r)$
$Q_{\varphi B}$	$\varphi^\dagger \varphi B_{\mu\nu} B^{\mu\nu}$	Q_{uB}	$(\bar{q}_p \sigma^{\mu\nu} u_r) \tilde{\varphi} B_{\mu\nu}$	$Q_{\varphi q}^{(3)}$	$(\varphi^\dagger i \overleftrightarrow{D}_\mu^I \varphi)(\bar{q}_p \tau^I \gamma^\mu q_r)$
$Q_{\varphi \tilde{B}}$	$\varphi^\dagger \varphi \tilde{B}_{\mu\nu} B^{\mu\nu}$	Q_{dG}	$(\bar{q}_p \sigma^{\mu\nu} T^A d_r) \varphi G_{\mu\nu}^A$	$Q_{\varphi u}$	$(\varphi^\dagger i \overleftrightarrow{D}_\mu \varphi)(\bar{u}_p \gamma^\mu u_r)$
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$Q_{\varphi \tilde{W}B}$	$\varphi^\dagger \tau^I \varphi \tilde{W}_{\mu\nu}^I B^{\mu\nu}$	Q_{dB}	$(\bar{q}_p \sigma^{\mu\nu} d_r) \varphi B_{\mu\nu}$	$Q_{\varphi ud}$	$i(\tilde{\varphi}^\dagger D_\mu \varphi)(\bar{u}_p \gamma^\mu d_r)$

Yukawa operators



- Decouple m_f and y_f
- $ff\bar{h}h(h)$ contact interactions

Dipole operators

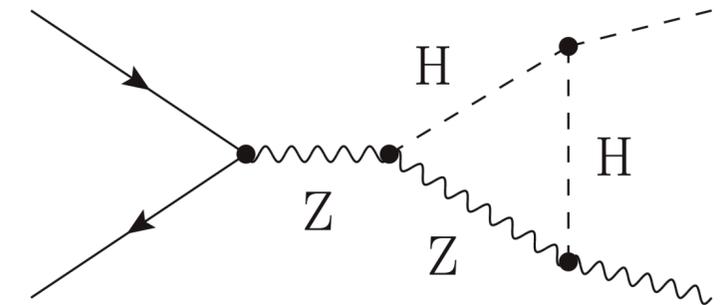


- Chirality flipping $ff\bar{V}$ couplings
- $ff\bar{V}h(h)$ contact interactions
- Momentum dependent vertex

The Warsaw basis

X^3		φ^6 and $\varphi^4 D^2$		$\psi^2 \varphi^3$	
Q_G	$f^{ABC} G_\mu^{A\nu} G_\nu^{B\rho} G_\rho^{C\mu}$	Q_φ	$(\varphi^\dagger \varphi)^3$	$Q_{e\varphi}$	$(\varphi^\dagger \varphi)(\bar{l}_p e_r \varphi)$
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$Q_{\varphi \tilde{G}}$	$\varphi^\dagger \varphi \tilde{G}_{\mu\nu}^A G^{A\mu\nu}$	Q_{eB}	$(\bar{l}_p \sigma^{\mu\nu} e_r) \varphi B_{\mu\nu}$	$Q_{\varphi l}^{(3)}$	$(\varphi^\dagger i \overleftrightarrow{D}_\mu^I \varphi) (\bar{l}_p \tau^I \gamma^\mu l_r)$
$Q_{\varphi W}$	$\varphi^\dagger \varphi W_{\mu\nu}^I W^{I\mu\nu}$	Q_{uG}	$(\bar{q}_p \sigma^{\mu\nu} T^A u_r) \tilde{\varphi} G_{\mu\nu}^A$	$Q_{\varphi e}$	$(\varphi^\dagger i \overleftrightarrow{D}_\mu \varphi) (\bar{e}_p \gamma^\mu e_r)$
$Q_{\varphi \tilde{W}}$	$\varphi^\dagger \varphi \tilde{W}_{\mu\nu}^I W^{I\mu\nu}$	Q_{uW}	$(\bar{q}_p \sigma^{\mu\nu} u_r) \tau^I \tilde{\varphi} W_{\mu\nu}^I$	$Q_{\varphi q}^{(1)}$	$(\varphi^\dagger i \overleftrightarrow{D}_\mu \varphi) (\bar{q}_p \gamma^\mu q_r)$
$Q_{\varphi B}$	$\varphi^\dagger \varphi B_{\mu\nu} B^{\mu\nu}$	Q_{uB}	$(\bar{q}_p \sigma^{\mu\nu} u_r) \tilde{\varphi} B_{\mu\nu}$	$Q_{\varphi q}^{(3)}$	$(\varphi^\dagger i \overleftrightarrow{D}_\mu^I \varphi) (\bar{q}_p \tau^I \gamma^\mu q_r)$
$Q_{\varphi \tilde{B}}$	$\varphi^\dagger \varphi \tilde{B}_{\mu\nu} B^{\mu\nu}$	Q_{dG}	$(\bar{q}_p \sigma^{\mu\nu} T^A d_r) \varphi G_{\mu\nu}^A$	$Q_{\varphi u}$	$(\varphi^\dagger i \overleftrightarrow{D}_\mu \varphi) (\bar{u}_p \gamma^\mu u_r)$
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Bosonic operators



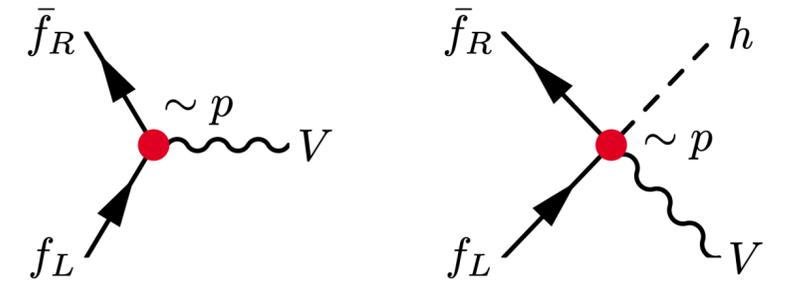
- H^6 modifies the Higgs self-coupling
- Enters at NLO EW in $e^+e^- \rightarrow ZH$

arXiv: 2406.03557

The Warsaw basis

	X^3	φ^6 and $\varphi^4 D^2$	$\psi^2 \varphi^3$
Q_G	$(\bar{L}L)(\bar{L}L)$	$(\bar{R}R)(\bar{R}R)$	$(\bar{L}L)(\bar{R}R)$
$Q_{\tilde{G}}$	$Q_{ll} \quad (\bar{l}_p \gamma_\mu l_r)(\bar{l}_s \gamma^\mu l_t)$	$Q_{ee} \quad (\bar{e}_p \gamma_\mu e_r)(\bar{e}_s \gamma^\mu e_t)$	$Q_{le} \quad (\bar{l}_p \gamma_\mu l_r)(\bar{e}_s \gamma^\mu e_t)$
Q_W	$Q_{qq}^{(1)} \quad (\bar{q}_p \gamma_\mu q_r)(\bar{q}_s \gamma^\mu q_t)$	$Q_{uu} \quad (\bar{u}_p \gamma_\mu u_r)(\bar{u}_s \gamma^\mu u_t)$	$Q_{lu} \quad (\bar{l}_p \gamma_\mu l_r)(\bar{u}_s \gamma^\mu u_t)$
$Q_{\tilde{W}}$	$Q_{qq}^{(3)} \quad (\bar{q}_p \gamma_\mu \tau^I q_r)(\bar{q}_s \gamma^\mu \tau^I q_t)$	$Q_{dd} \quad (\bar{d}_p \gamma_\mu d_r)(\bar{d}_s \gamma^\mu d_t)$	$Q_{ld} \quad (\bar{l}_p \gamma_\mu l_r)(\bar{d}_s \gamma^\mu d_t)$
	$Q_{lq}^{(1)} \quad (\bar{l}_p \gamma_\mu l_r)(\bar{q}_s \gamma^\mu q_t)$	$Q_{eu} \quad (\bar{e}_p \gamma_\mu e_r)(\bar{u}_s \gamma^\mu u_t)$	$Q_{qe} \quad (\bar{q}_p \gamma_\mu q_r)(\bar{e}_s \gamma^\mu e_t)$
	$Q_{lq}^{(3)} \quad (\bar{l}_p \gamma_\mu \tau^I l_r)(\bar{q}_s \gamma^\mu \tau^I q_t)$	$Q_{ed} \quad (\bar{e}_p \gamma_\mu e_r)(\bar{d}_s \gamma^\mu d_t)$	$Q_{qu}^{(1)} \quad (\bar{q}_p \gamma_\mu q_r)(\bar{u}_s \gamma^\mu u_t)$
$Q_{\varphi G}$		$Q_{ud}^{(1)} \quad (\bar{u}_p \gamma_\mu u_r)(\bar{d}_s \gamma^\mu d_t)$	$Q_{qu}^{(8)} \quad (\bar{q}_p \gamma_\mu T^A q_r)(\bar{u}_s \gamma^\mu T^A u_t)$
$Q_{\varphi \tilde{G}}$		$Q_{ud}^{(8)} \quad (\bar{u}_p \gamma_\mu T^A u_r)(\bar{d}_s \gamma^\mu T^A d_t)$	$Q_{qd}^{(1)} \quad (\bar{q}_p \gamma_\mu q_r)(\bar{d}_s \gamma^\mu d_t)$
$Q_{\varphi W}$			$Q_{qd}^{(8)} \quad (\bar{q}_p \gamma_\mu T^A q_r)(\bar{d}_s \gamma^\mu T^A d_t)$
$Q_{\varphi \tilde{W}}$	$(\bar{L}R)(\bar{R}L)$ and $(\bar{L}R)(\bar{L}R)$	B -violating	
$Q_{\varphi B}$	$Q_{ledq} \quad (\bar{l}_p^j e_r)(\bar{d}_s q_t^j)$	$Q_{duq} \quad \varepsilon^{\alpha\beta\gamma} \varepsilon_{jk} [(d_p^\alpha)^T C u_r^\beta] [(q_s^j)^T C l_t^k]$	
$Q_{\varphi \tilde{B}}$	$Q_{quqd}^{(1)} \quad (\bar{q}_p^j u_r) \varepsilon_{jk} (\bar{q}_s^k d_t)$	$Q_{qqu} \quad \varepsilon^{\alpha\beta\gamma} \varepsilon_{jk} [(q_p^{\alpha j})^T C q_r^{\beta k}] [(u_s^\gamma)^T C e_t]$	
$Q_{\varphi WB}$	$Q_{quqd}^{(8)} \quad (\bar{q}_p^j T^A u_r) \varepsilon_{jk} (\bar{q}_s^k T^A d_t)$	$Q_{qqq} \quad \varepsilon^{\alpha\beta\gamma} \varepsilon_{jn} \varepsilon_{km} [(q_p^{\alpha j})^T C q_r^{\beta k}] [(q_s^{\gamma m})^T C l_t^n]$	
$Q_{\varphi \tilde{W}B}$	$Q_{lequ}^{(1)} \quad (\bar{l}_p^j e_r) \varepsilon_{jk} (\bar{q}_s^k u_t)$	$Q_{duu} \quad \varepsilon^{\alpha\beta\gamma} [(d_p^\alpha)^T C u_r^\beta] [(u_s^\gamma)^T C e_t]$	
	$Q_{lequ}^{(3)} \quad (\bar{l}_p^j \sigma_{\mu\nu} e_r) \varepsilon_{jk} (\bar{q}_s^k \sigma^{\mu\nu} u_t)$		

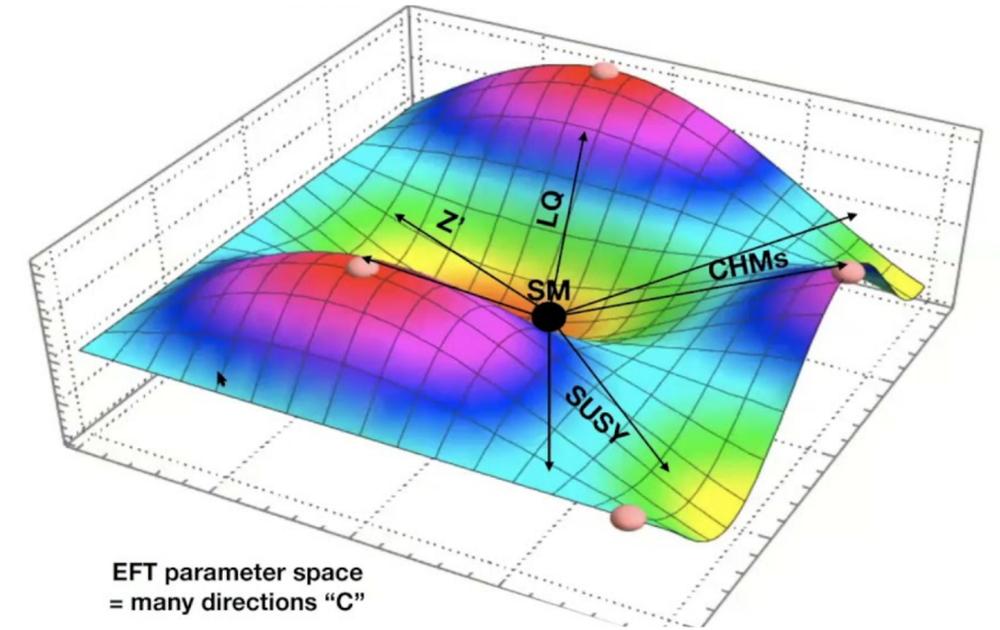
Four-fermion operators



- ▶ Contact interactions
- ▶ 2205 B/L preserving

SMEFT predictions

$$\mathcal{L}_{\text{SMEFT}} = \mathcal{L}_{\text{SM}} + \sum_i^{N_{d5}} \frac{c_i}{\Lambda} \mathcal{O}_i^{(5)} + \sum_i^{N_{d6}} \frac{c_i}{\Lambda^2} \mathcal{O}_i^{(6)} + \dots$$



$$\sigma_{\text{SMEFT}}(c, \Lambda) = \sigma_{\text{SM}} \times \left(1 + \sum_i^{N_{d6}} \kappa_i \frac{c_i}{\Lambda^2} + \sum_{i < j}^{N_{d6}} \tilde{\kappa}_{ij} \frac{c_i \cdot c_j}{\Lambda^4} + \mathcal{O}(\Lambda^{-6}) \right)$$

Linear EFT corrections:
interference with SM

Quadratic EFT corrections

Building the likelihood

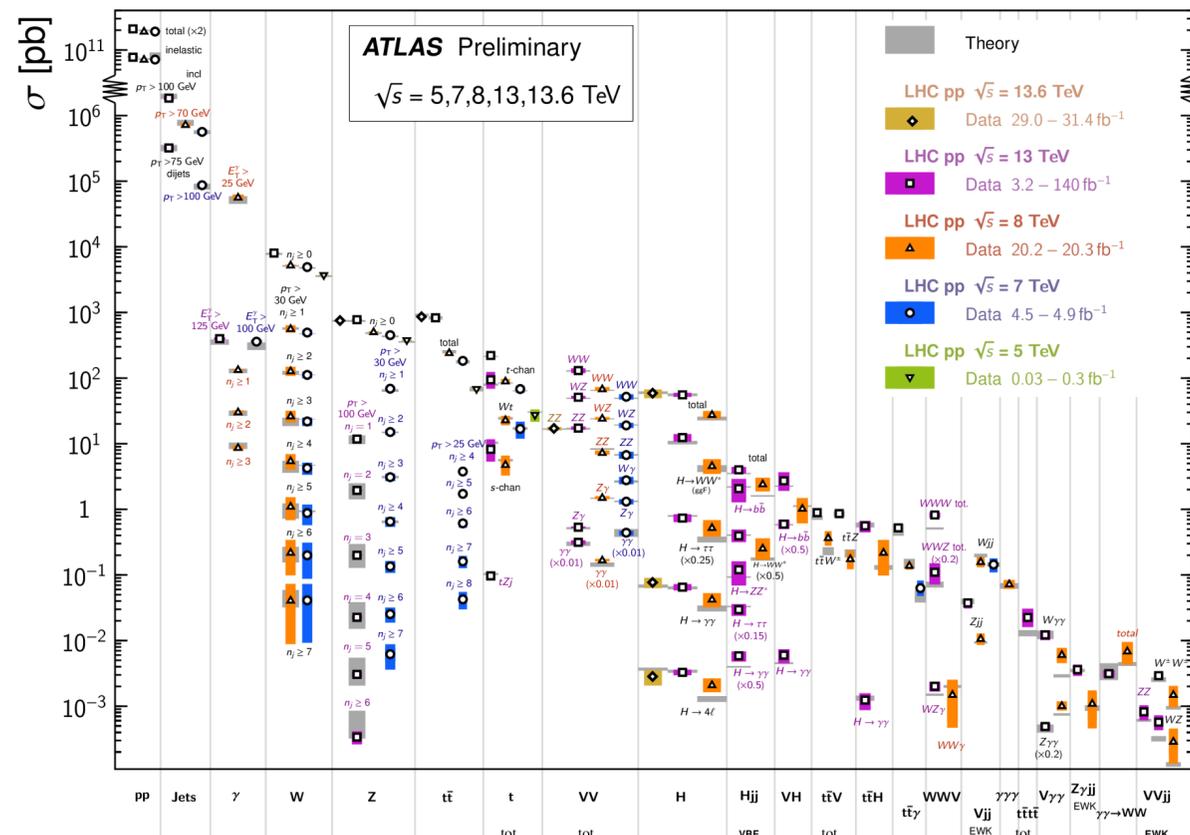
$$\chi^2 = \frac{1}{n_{\text{dat}}} \sum_{i,j=1}^{n_{\text{dat}}} \left(\sigma_{i,\text{SMEFT}}(c) - \sigma_{i,\text{exp}} \right) (\text{cov}^{-1})_{ij} \left(\sigma_{j,\text{SMEFT}}(c) - \sigma_{j,\text{exp}} \right)$$

[ATL-PHYS-PUB-2023-039]

$$\text{cov}^{(\text{tot})}_{ij} = \text{cov}^{(\text{th})}_{ij} + \text{cov}^{(\text{exp})}_{ij}$$

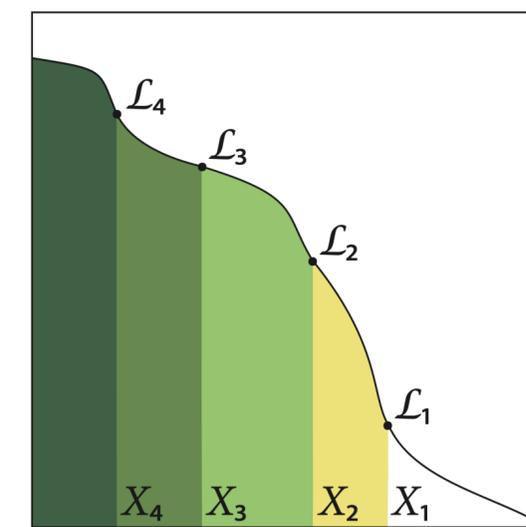
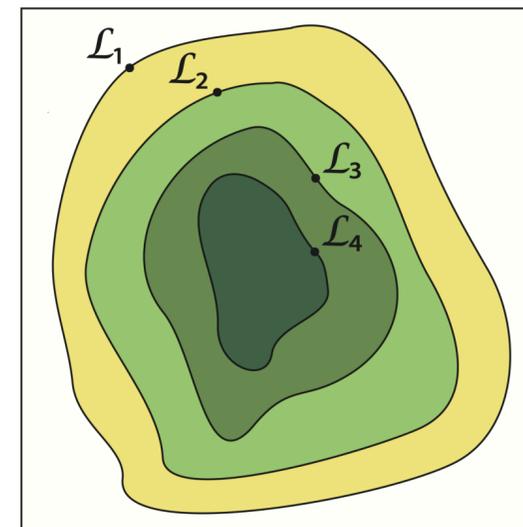
Standard Model Production Cross Section Measurements

Status: October 2023



Analytic if $\mathcal{O}(\Lambda^{-2})$, fast!

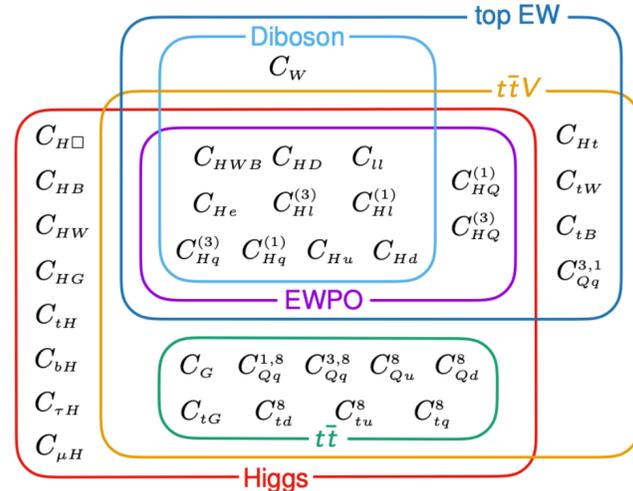
Nested Sampling $\mathcal{O}(\Lambda^{-2})$ or $\mathcal{O}(\Lambda^{-4})$



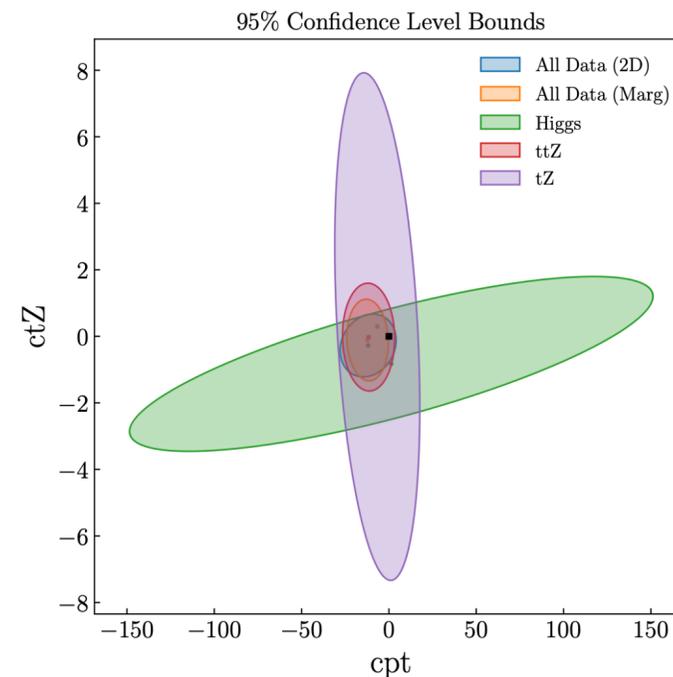
Feroz et al [1306.2144]

Why global SMEFT fits?

- ▶ **Cross-talk** between Higgs, top, diboson and EWPO (and flavour and low energy observables) requires a global analysis
- ▶ **Challenge:** a large number of operators, with many datasets needed to break degeneracies

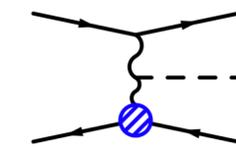
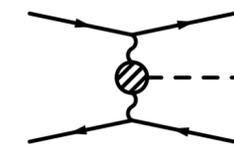


[2012.02779] Fitmaker collaboration

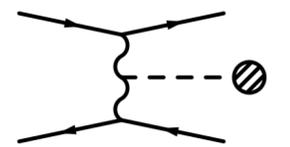


[2105.00006] SMEFiT collaboration

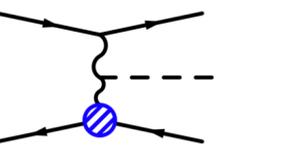
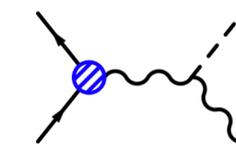
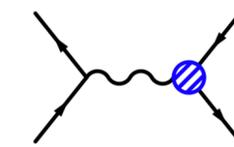
One observable can be influenced by many operators



Higgs decay



One operator can contribute to many different observables

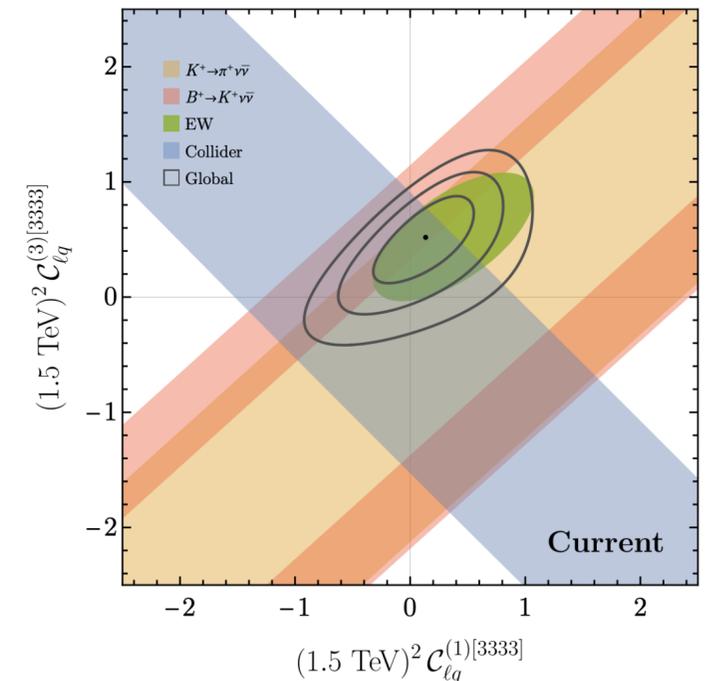


Anke Biekötter - HET seminar Brookhaven

Previously on global fits...



- **SMEFiT**: EW + Higgs + diboson + top + projections, NLO, quadratic - [2105.00006, 2309.04523, 2404.12809]
- **ATLAS**: EW + Higgs, LO, quadratic - [ATL-PHYS-PUB-2022-037]
- **simuNET**: simultaneous EFT + PDF fit in EW + Higgs + diboson + top, NLO, linear - [2402.03308]
- **Fitmaker**: EW + Higgs + top + diboson, linear - [2012.02779, 2204.05260]
- **SFitter**: EW + Higgs, top, NLO, quadratic - [1812.07587, 1910.03606]
- **HEPfit**: EW, flavour, projections, LO, linear - [1910.14012]
- **TopFitter**: top, linear, LO - [1901.03164]
- **EFTfitter**: top + DY + flavour, LO, quadratic, RG effects - [1605.05585, 2304.12837]
- **Mainz group**: EW + Higgs + top + flavour + dijet + PV + lepton scattering, NLO, linear - [2311.04963]
- **Zurich group**: EW + flavour + (DY, LEP II, Jet observables), individual, RG effects - [2311.00020]
-

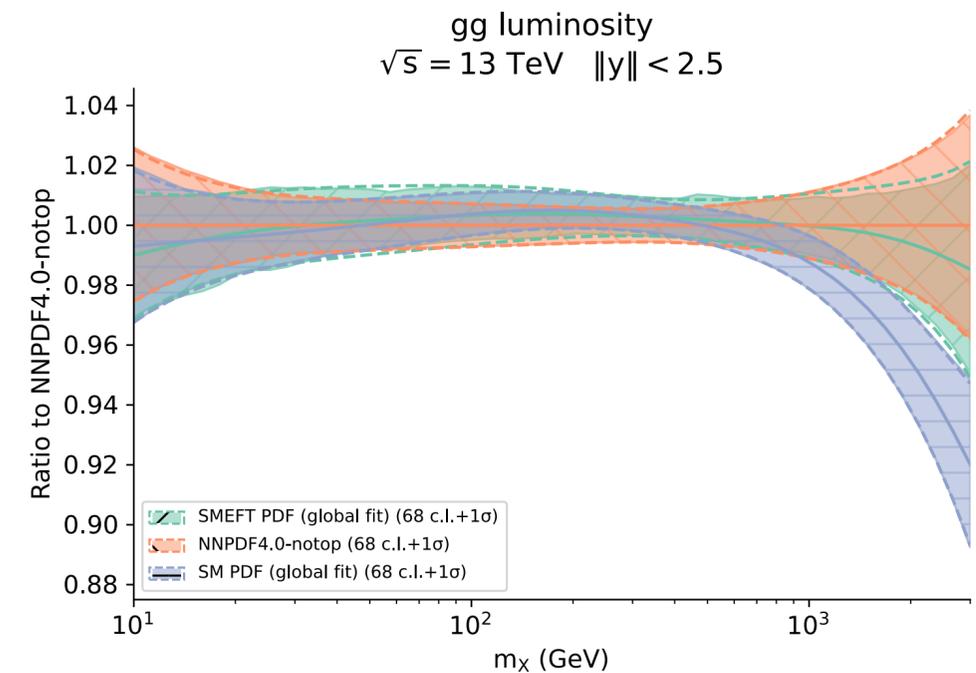
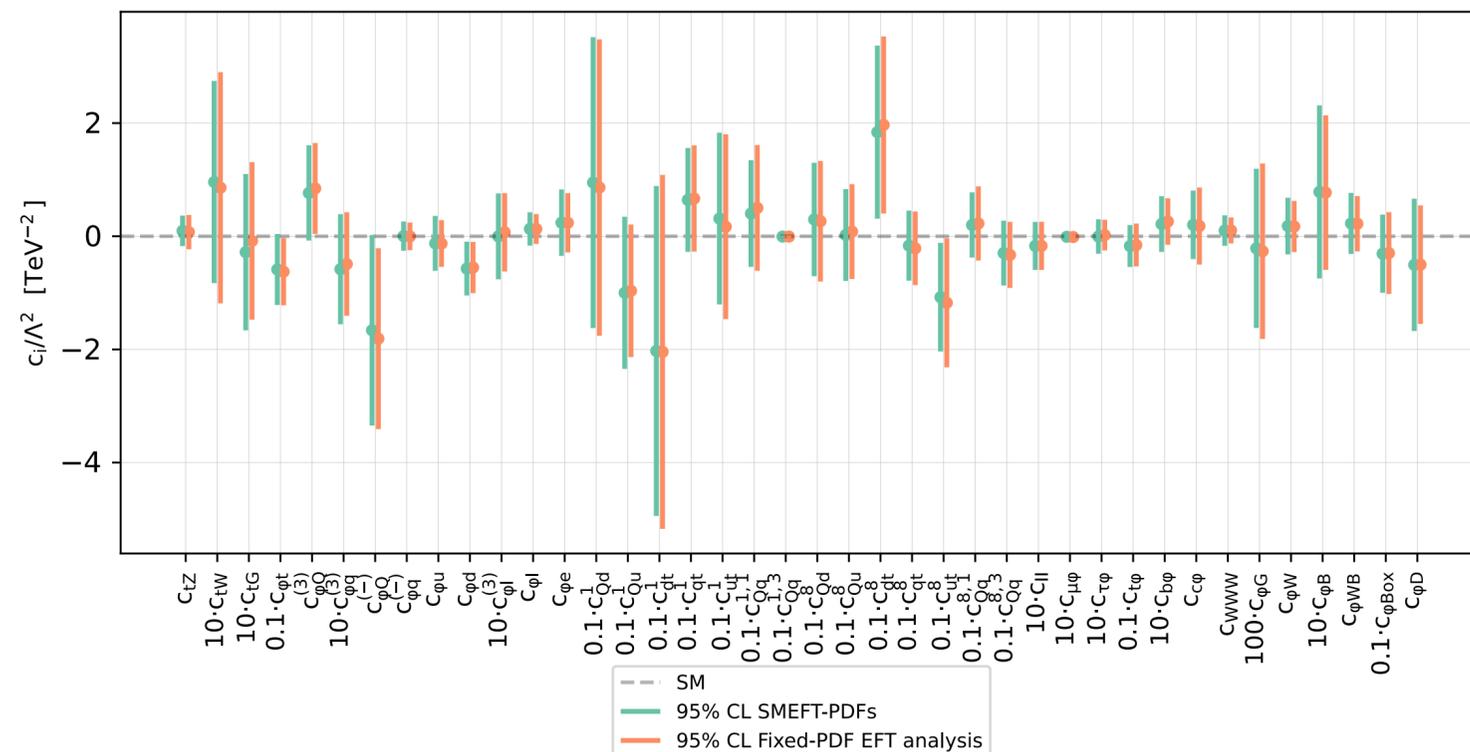
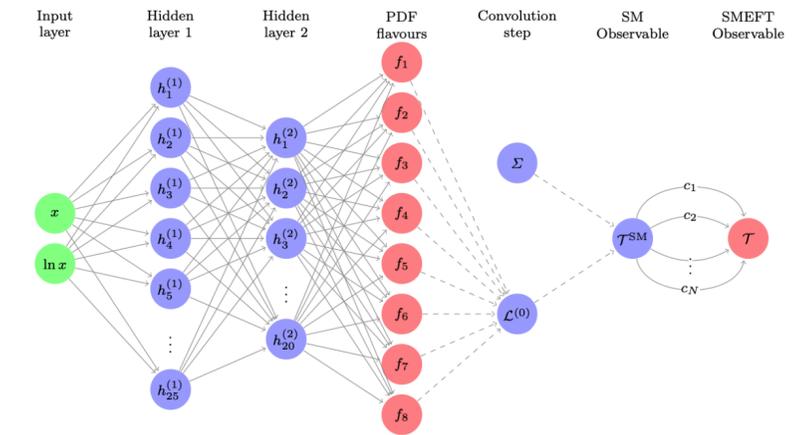


arXiv:2311.00020

simuNET: a simultaneous PDF + EFT fit

Costantini, Hammou, Madigan,
Mantani, Moore, Morales, Ubiali
[2402.03308]

- Most EFT global fits assume a fixed PDF set. Ideally, a full treatment fits the EFT and PDF parameters **simultaneously**, as done by simuNET
- EFT parameters are stable, while the PDF fits undergo **shifts** at high invariant mass in e.g. the gluon-gluon luminosity



RGE in the SMEFT



Connecting scales

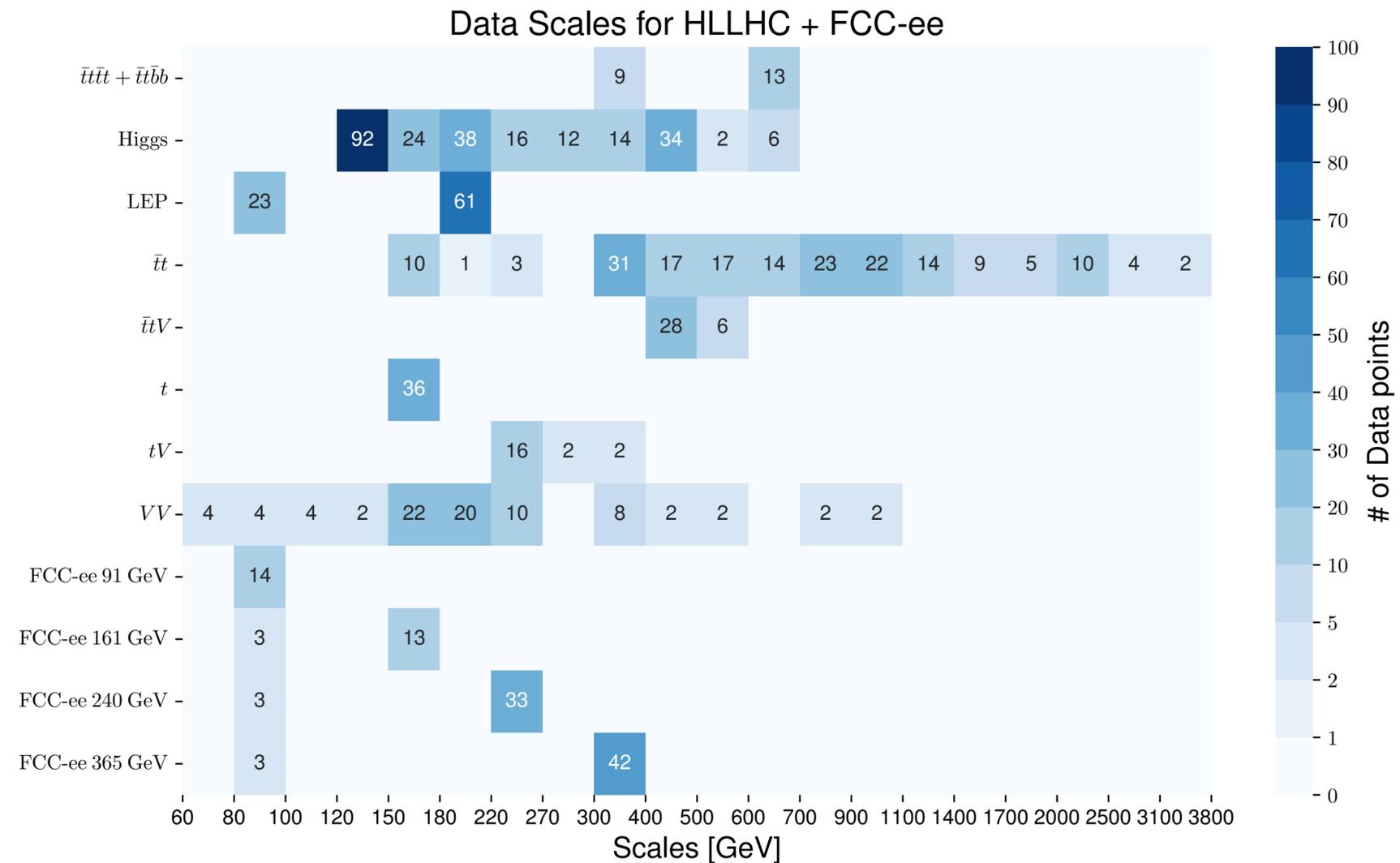
- Experimental input to global fits spans a **wide range of different energy scales**, from m_Z at LEP to $m_{t\bar{t}} \sim 3 \text{ TeV}$ in tails at LHC

- Wilson coefficients run with energy

$$\frac{dc_i(\mu)}{d \ln \mu} = \sum_{j=1}^{n_{\text{op}}} \gamma_{ij}^{(6)}(\bar{g}) c_j(\mu)$$

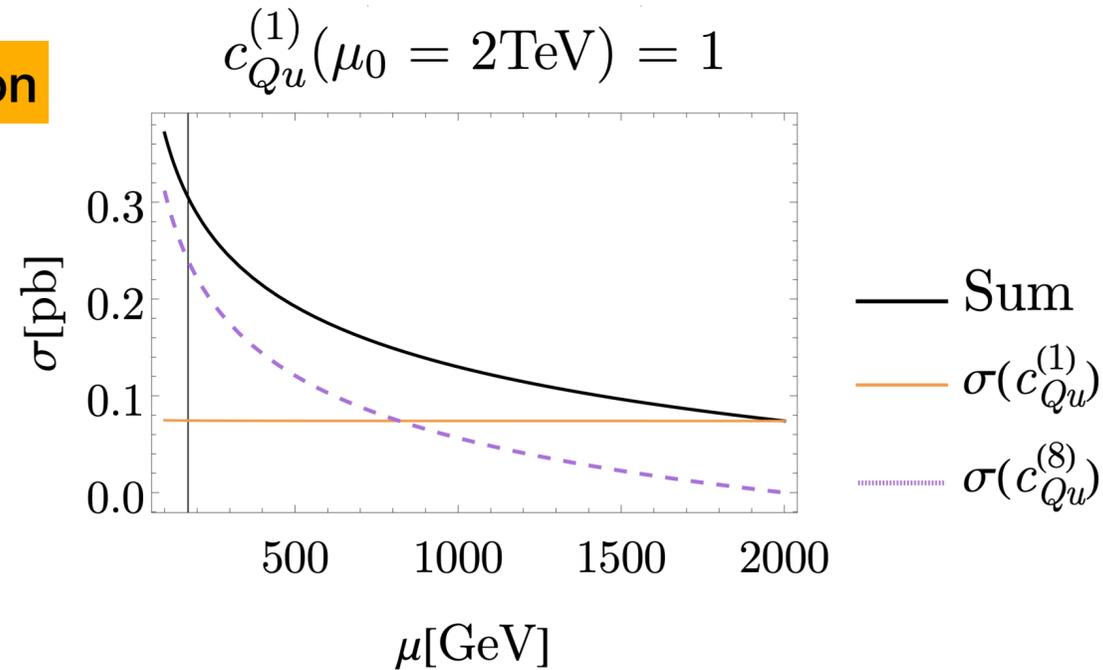
Jenkins, Manohar, Trott, Alonso
arXiv:1308.2627, 1310.4838, 1312.2014

- Operator mixing through the anomalous dimension

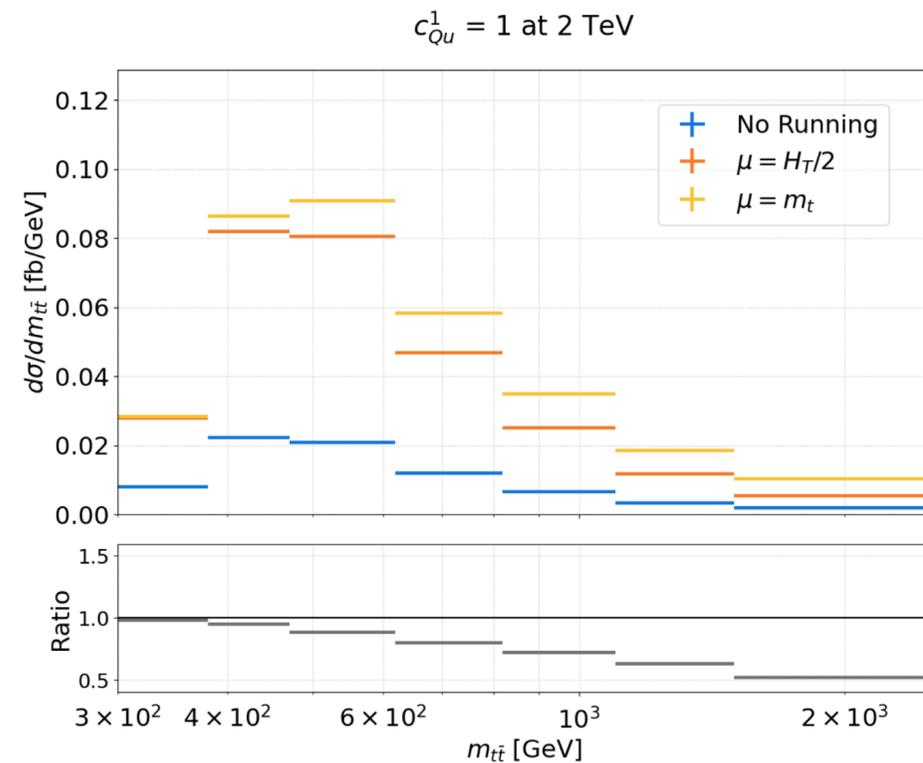


RGE in SMEFT phenomenology

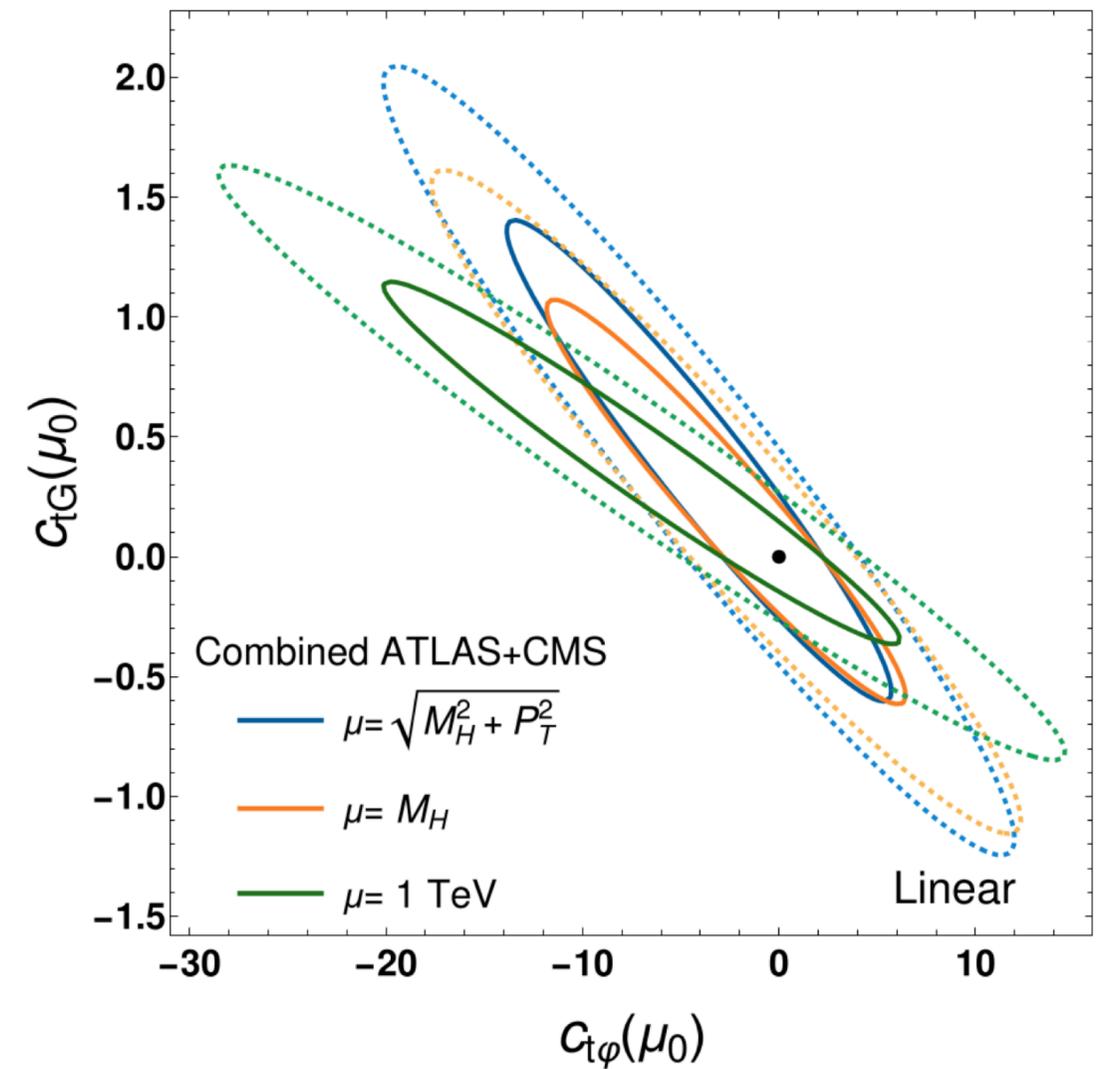
Top pair production



Aoude et al.
arXiv: 2212.05067



Higgs production



Maltoni, Ventura, Vryonidou
arXiv: 2406.06670

- ▶ We include RG effects in the **Matrix evolution approximation** interfaced to Wilson

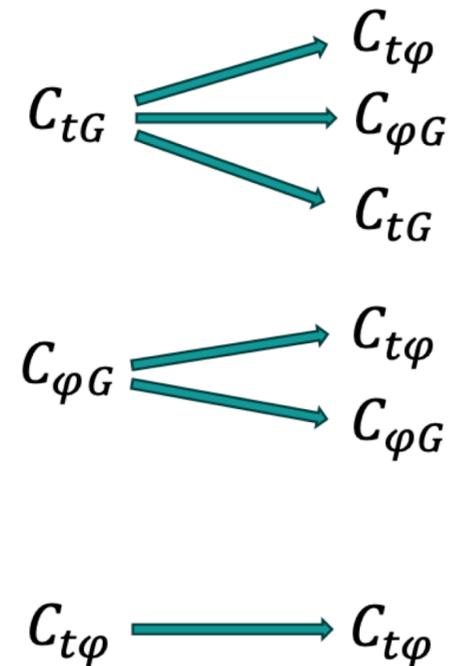
$$c_i(\mu) = \sum_{j=1}^{n_{\text{op}}} \Gamma_{ij}(\mu, \mu_0; \alpha_s, \alpha) c_j(\mu_0)$$



[1804.05033] Aebischer,
Kumar, Straub

- ▶ Express theory predictions at a common scale μ_0

$$\begin{aligned} T_{\text{EFT}}(\mathbf{c}(\mu)/\Lambda^2) &= T_{\text{SM}} + \sum_{i=1}^{n_{\text{op}}} \kappa_i \frac{c_i(\mu)}{\Lambda^2} + \sum_{i,j=1}^{n_{\text{op}}} \tilde{\kappa}_{ij} \frac{c_i(\mu)c_j(\mu)}{\Lambda^4} \\ &= T_{\text{SM}} + \sum_{i,j=1}^{n_{\text{op}}} \kappa_i \Gamma_{ij} \frac{c_j(\mu_0)}{\Lambda^2} + \sum_{i,j,k,l=1}^{n_{\text{op}}} \tilde{\kappa}_{ij} \Gamma_{ik} \Gamma_{jl} \frac{c_k(\mu_0)c_l(\mu_0)}{\Lambda^4} \\ &= T_{\text{SM}} + \sum_{j=1}^{n_{\text{op}}} \kappa'_j \frac{c_j(\mu_0)}{\Lambda^2} + \sum_{k,l=1}^{n_{\text{op}}} \tilde{\kappa}'_{kl} \frac{c_k(\mu_0)c_l(\mu_0)}{\Lambda^4} \end{aligned}$$



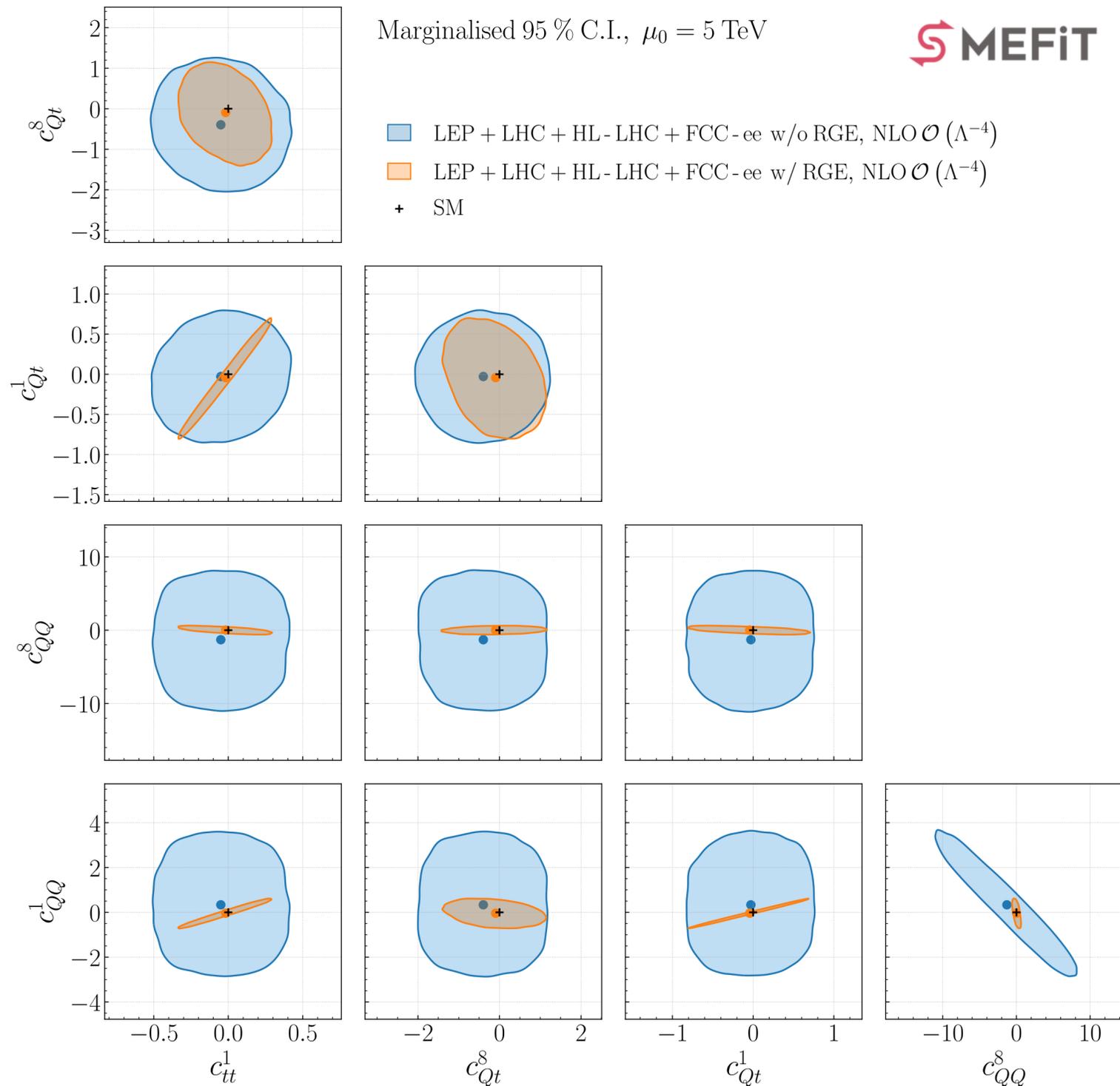
RGE in SMEFiT

- ▶ We associate a fixed characteristic scale choice μ to each bin
- ▶ "Event by event" running was performed in by Aoude et al in arXiv: 2212.05067
- ▶ We have studied the impact of our scale choice by varying $\mu \rightarrow \tilde{\mu} = \kappa\mu$

Process	Scale Choice μ	Process	Scale Choice μ
Higgs (ggF)	$\sqrt{m_H^2 + (p_T^H)^2}$	$t\bar{t}b\bar{b}$	$2m_t$
Higgs (VBF)	$\sqrt{m_H^2 + (p_T^H)^2}$	$t\bar{t}V$	$\sqrt{(2m_t + m_V)^2 + (p_T^V)^2}$
VH	$\sqrt{(m_V + m_H)^2 + (p_T^V)^2}$	tV	$m_t + m_V$ or $\sqrt{(m_t + m_V)^2 + (p_T^t)^2}$
$t\bar{t}H$	$\sqrt{(2m_t + m_H)^2 + (p_T^H)^2}$	W-helicities	m_t
tH	$m_t + m_H$	WZ	m_T^{WZ} or $\sqrt{(m_Z + m_W)^2 + (p_T^Z)^2}$
$t\bar{t}$	$m_{t\bar{t}}$	WW	$m_{e\mu}$
Single- t	m_t	V pole (incl. EWPOs)	m_V
$t\bar{t}\gamma$	$2m_t$	Bhabha scattering	\sqrt{s}
$t\bar{t}t\bar{t}$	$4m_t$	$e^+e^- \rightarrow WW / t\bar{t} / f\bar{f}$	\sqrt{s}
HH	$2m_H$	$e^+e^- \rightarrow ZH$	\sqrt{s}

RG in restricted operator basis

Mantani, Rojo, Rossia, Vryonidou, **JtH**
arXiv: 2502.20453

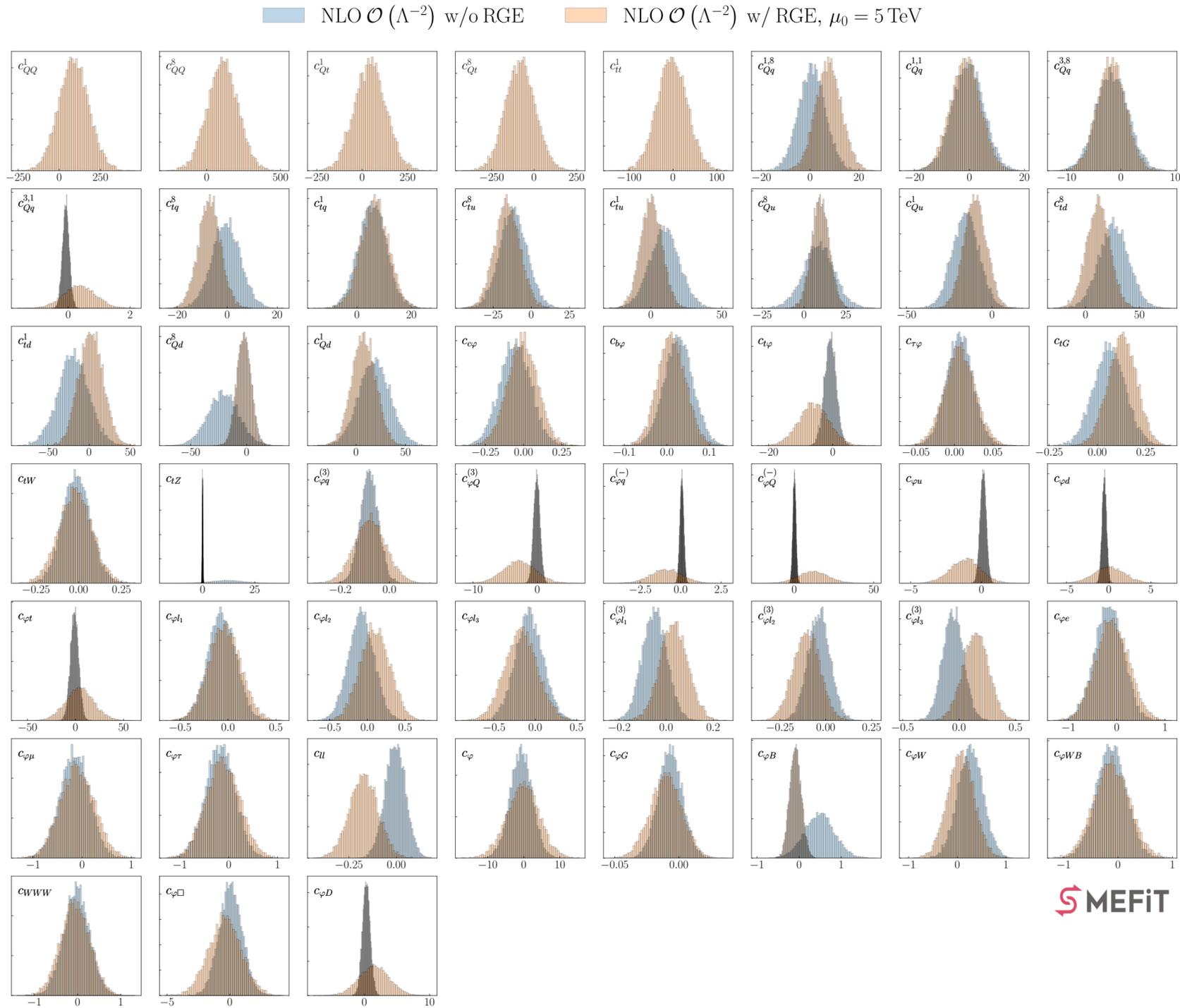


Example: 4 heavy operators flow into operators sensitive to the EWPOs at low energy

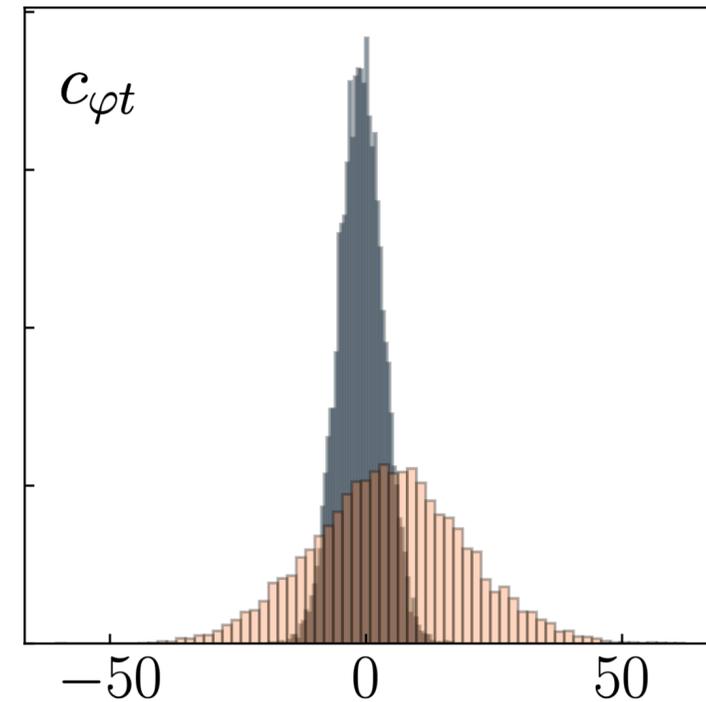
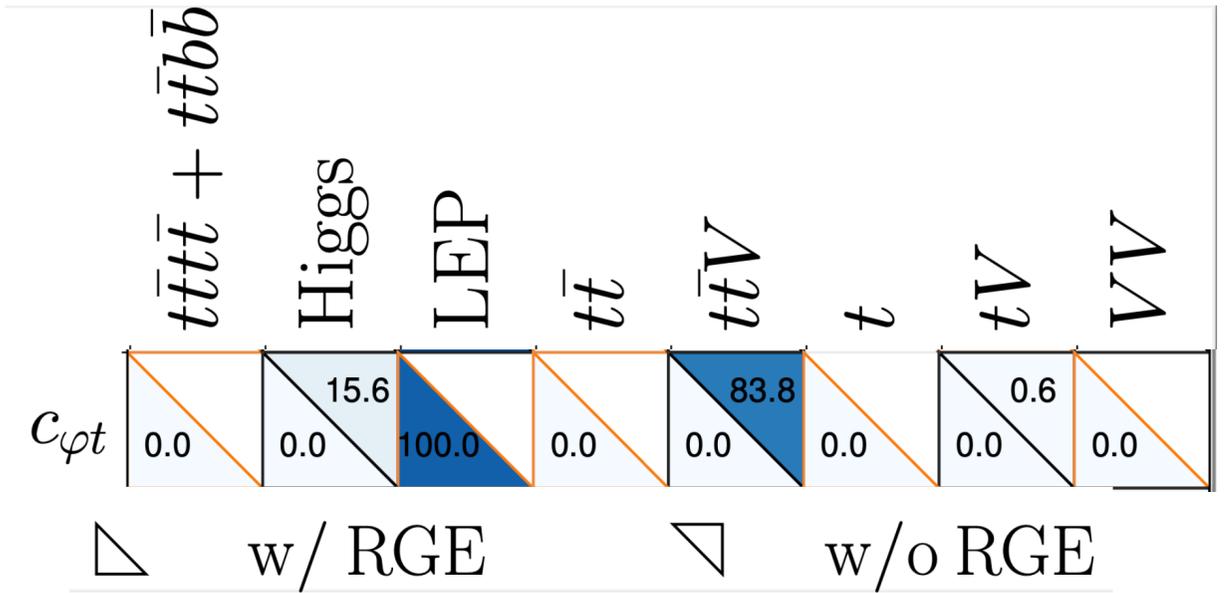
In general one may see two competing effects:

- ▶ Ill constrained operators may flow into a precisely determined observable
- ▶ More operators enter the same observable, making bounds weaker

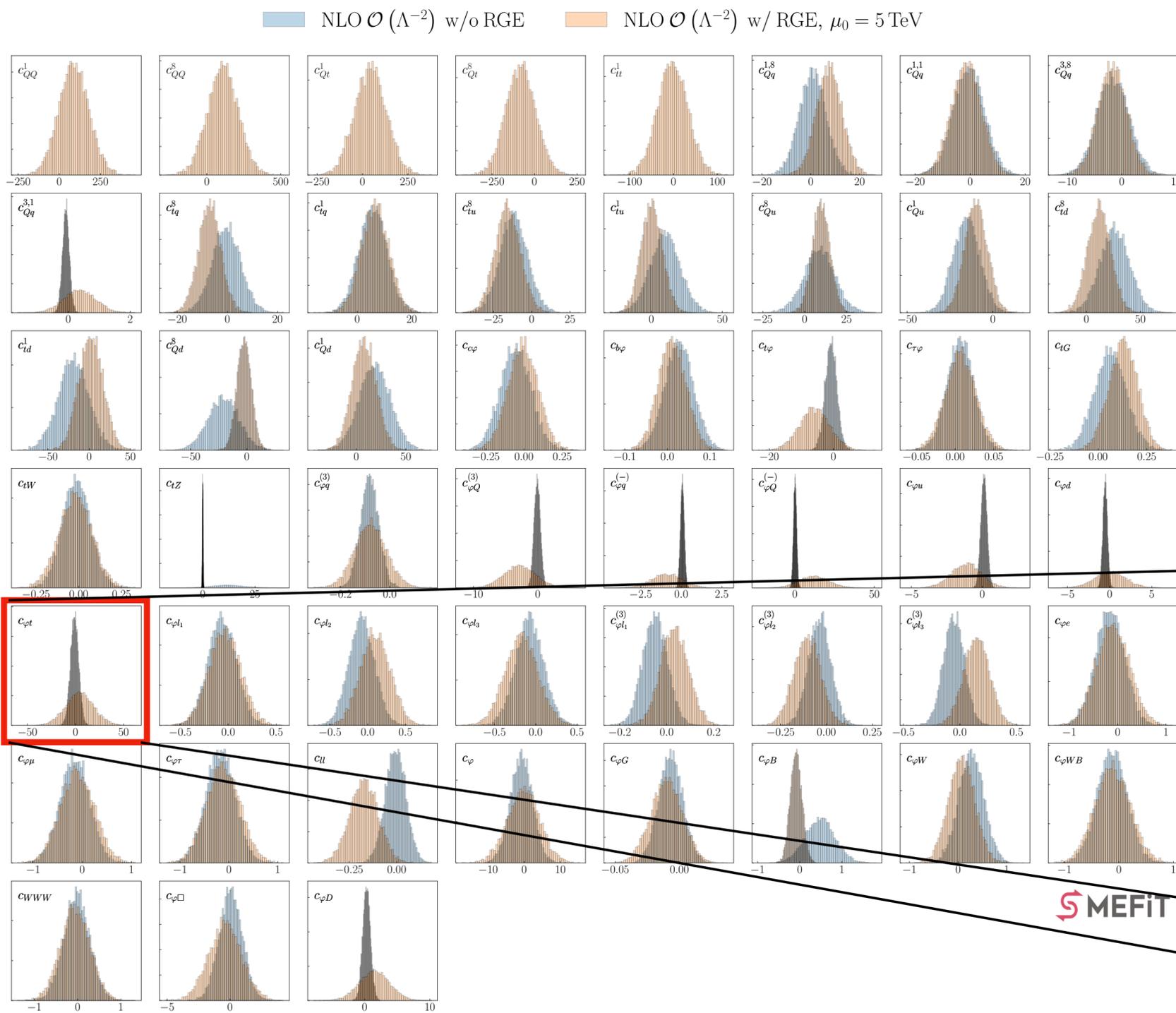
RG effects in the global fit



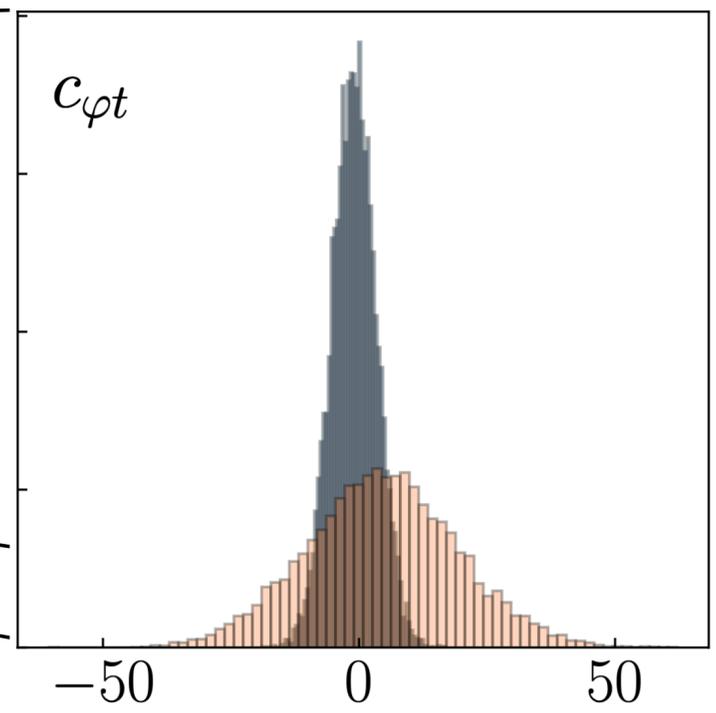
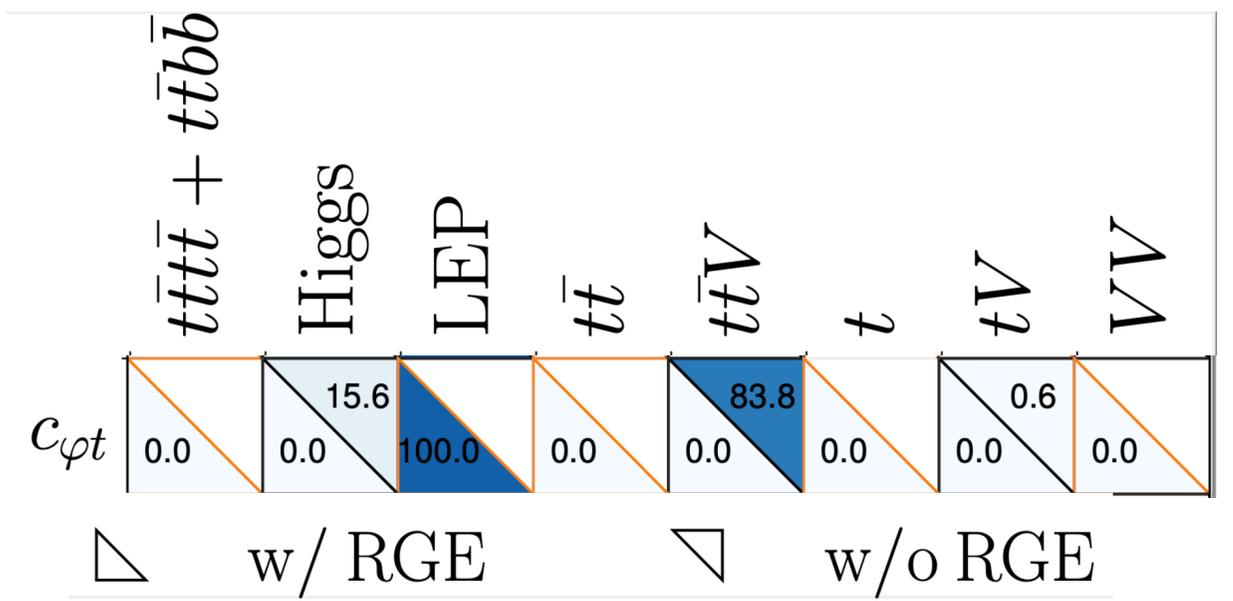
Fisher information moves from $t\bar{t}V$ to LEP



RG effects in the global fit



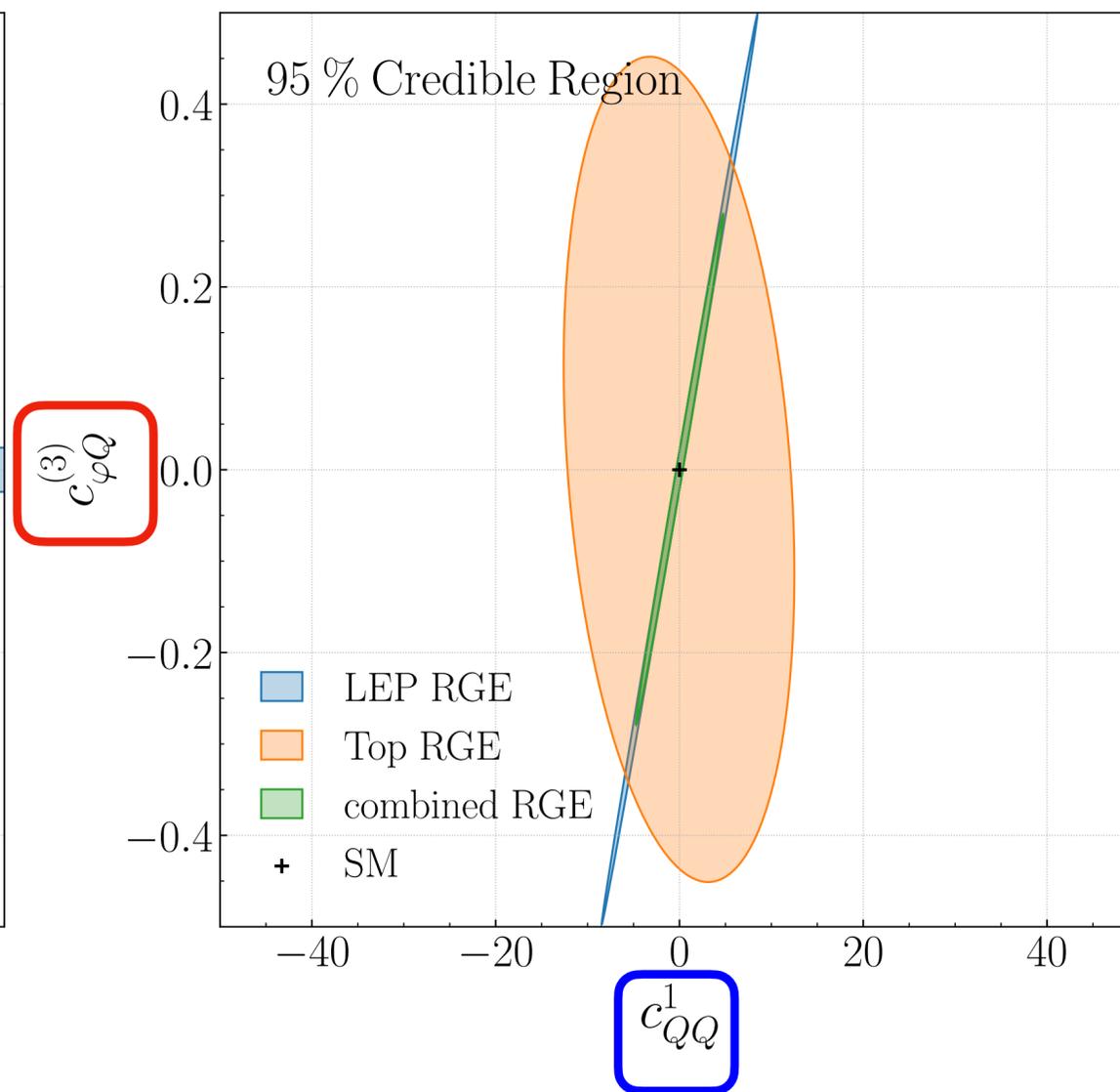
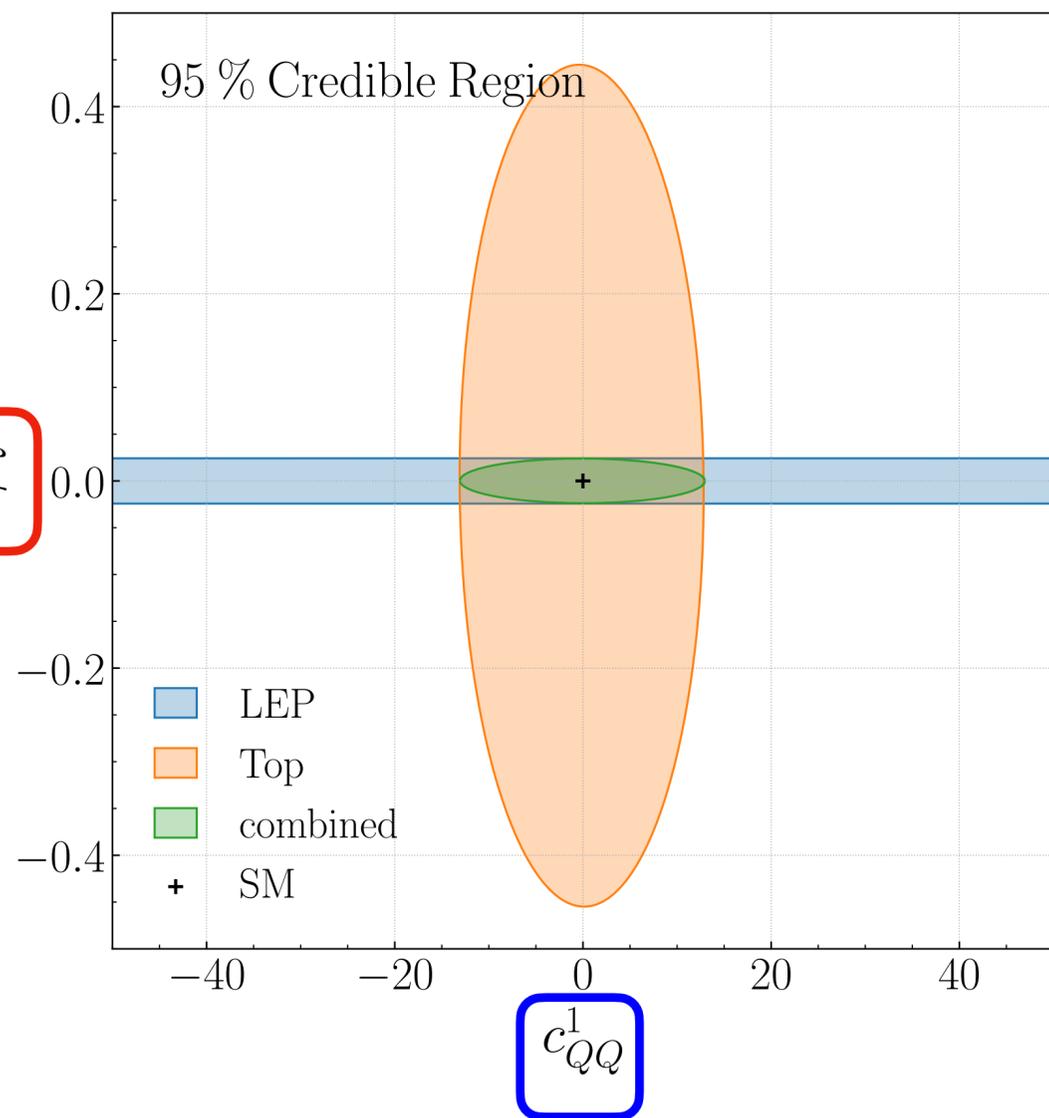
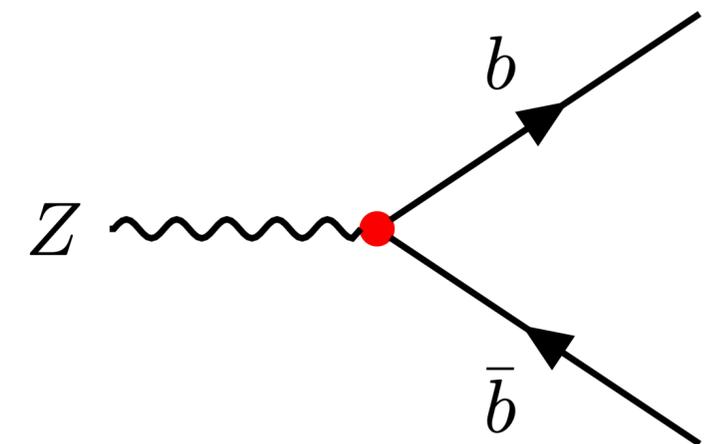
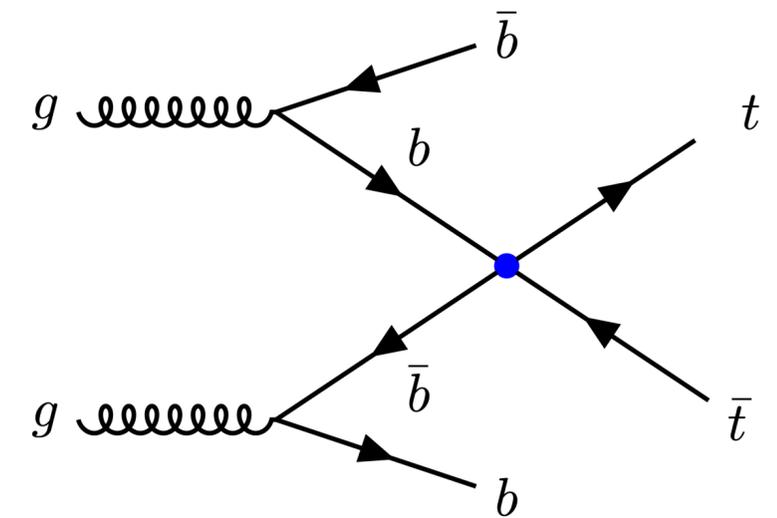
Fisher information moves from $t\bar{t}V$ to LEP



MEFIT

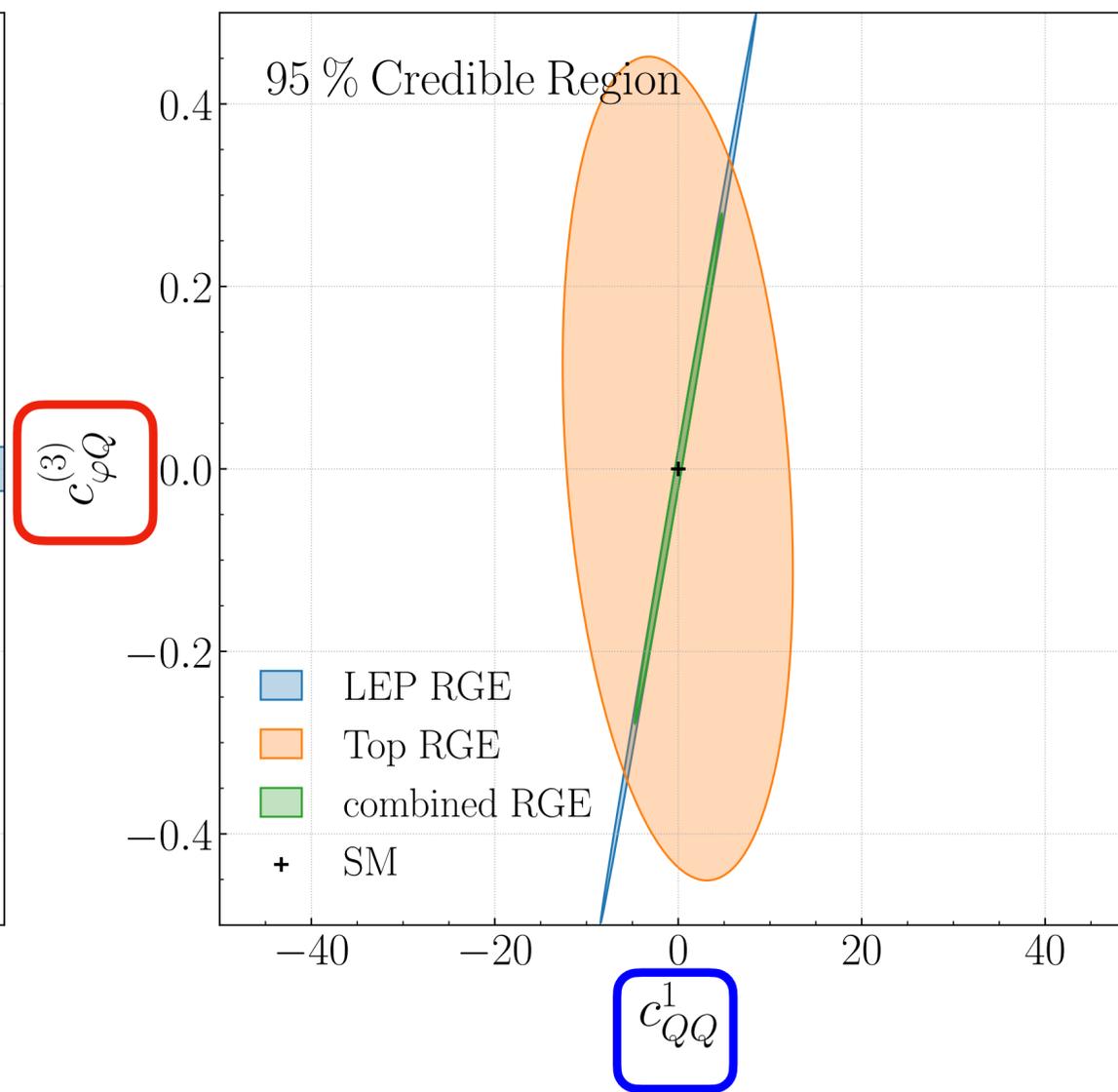
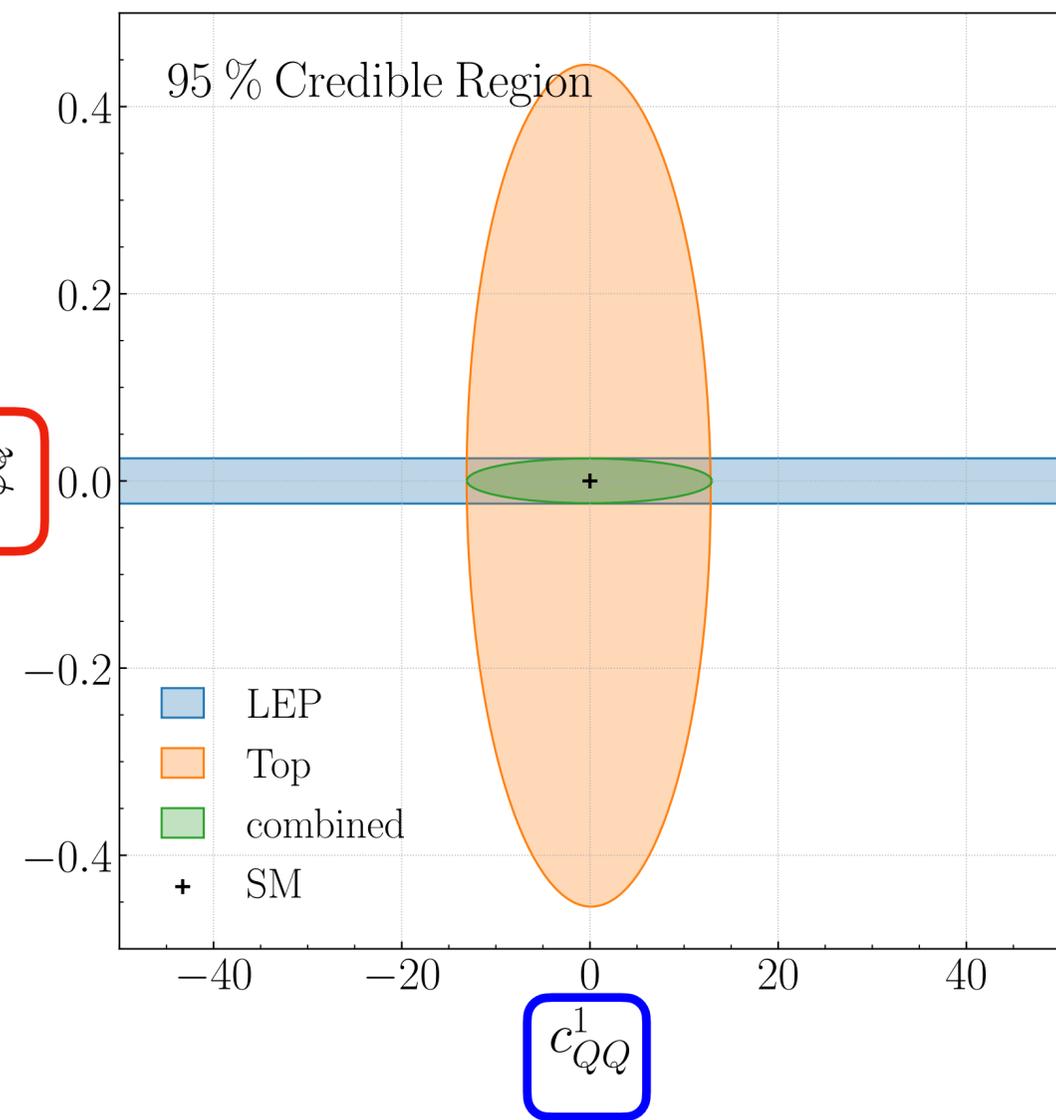
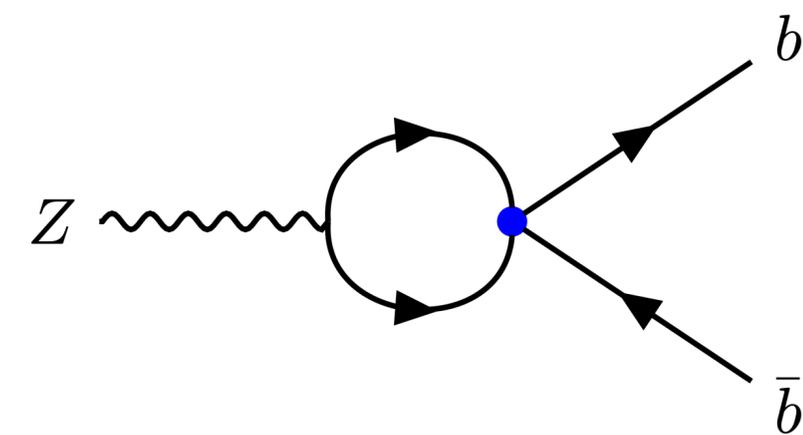
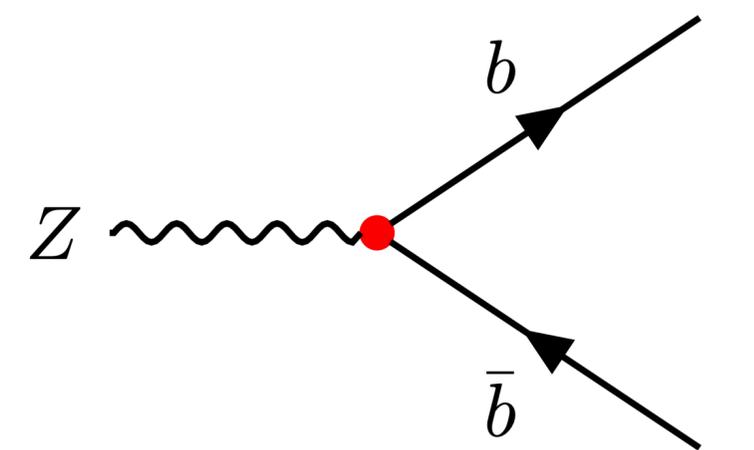
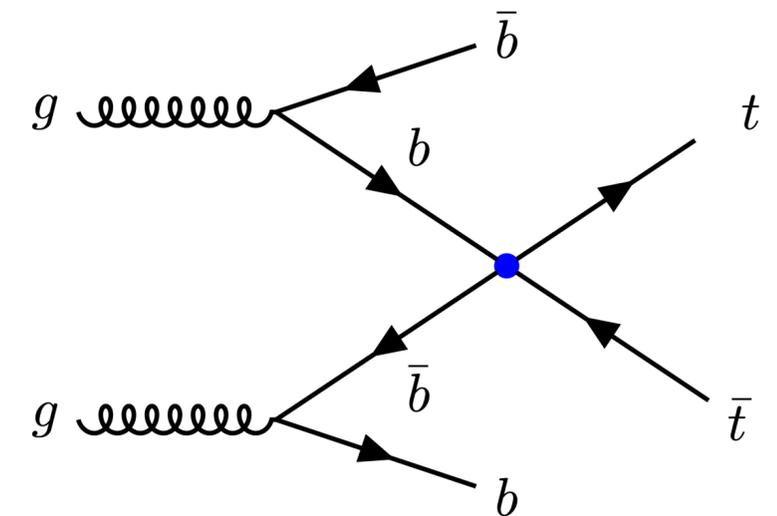
RG effects in the global fit

Global fits neglecting RGE effects can severely overestimate the bounds!



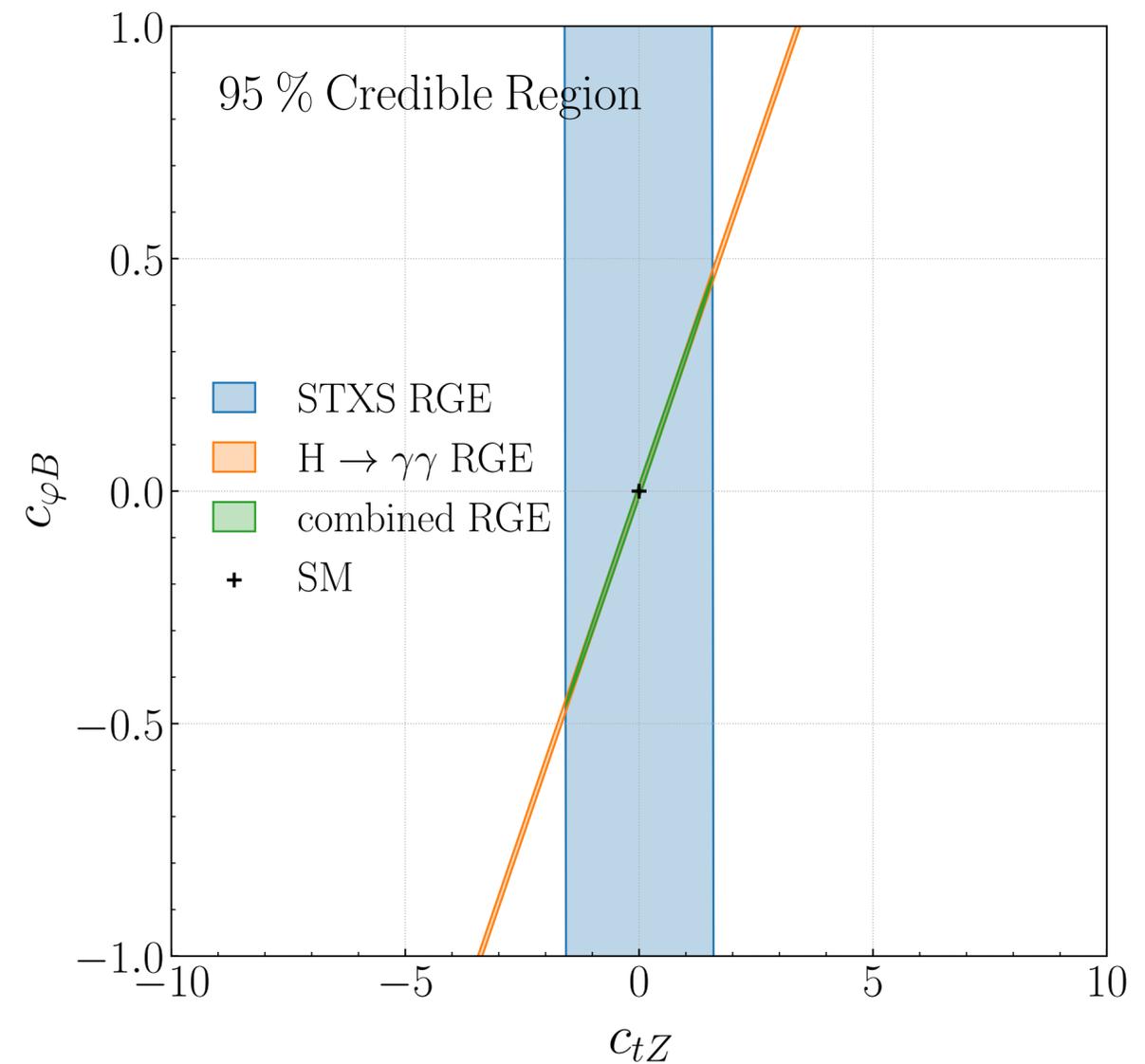
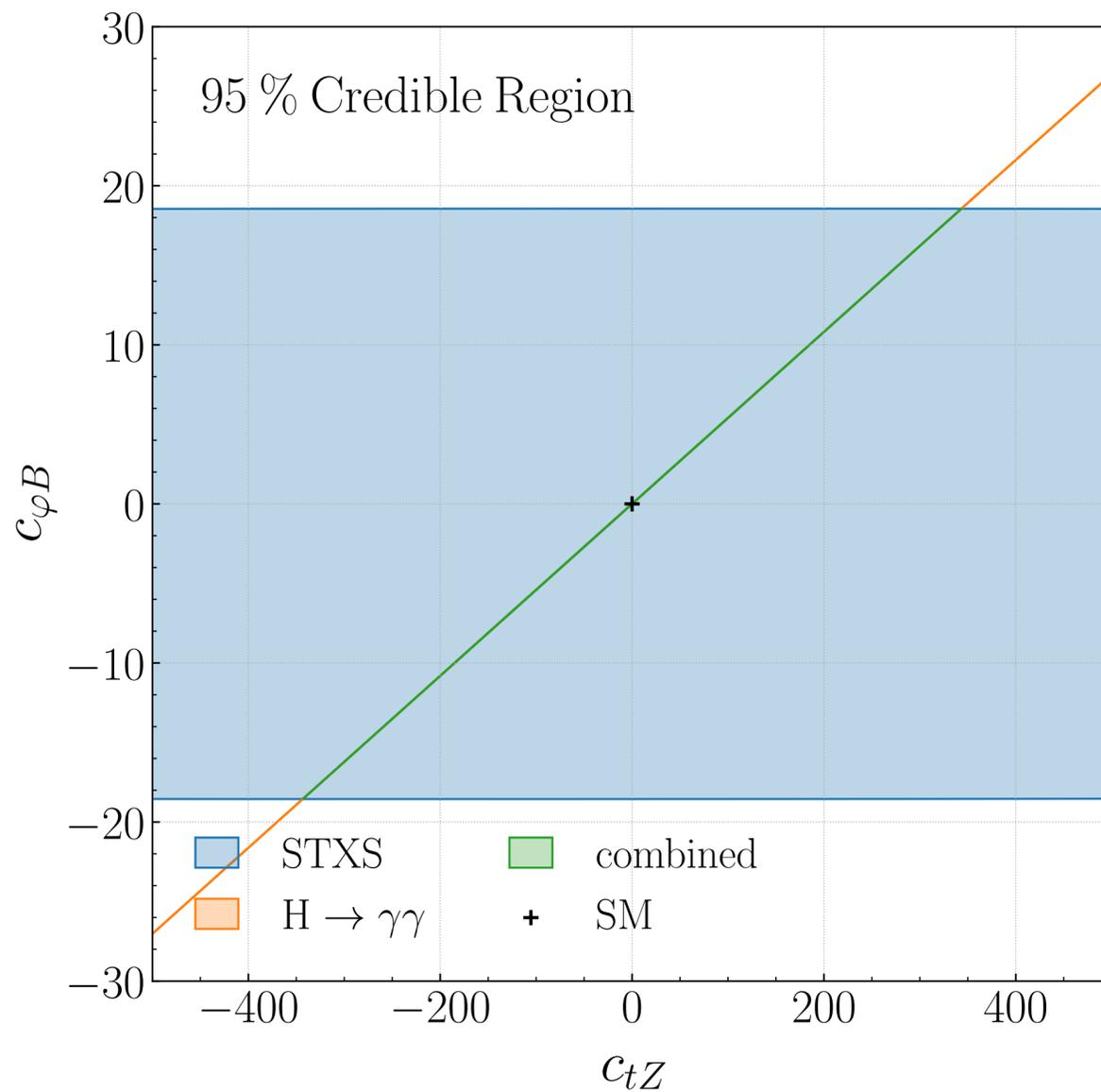
RG effects in the global fit

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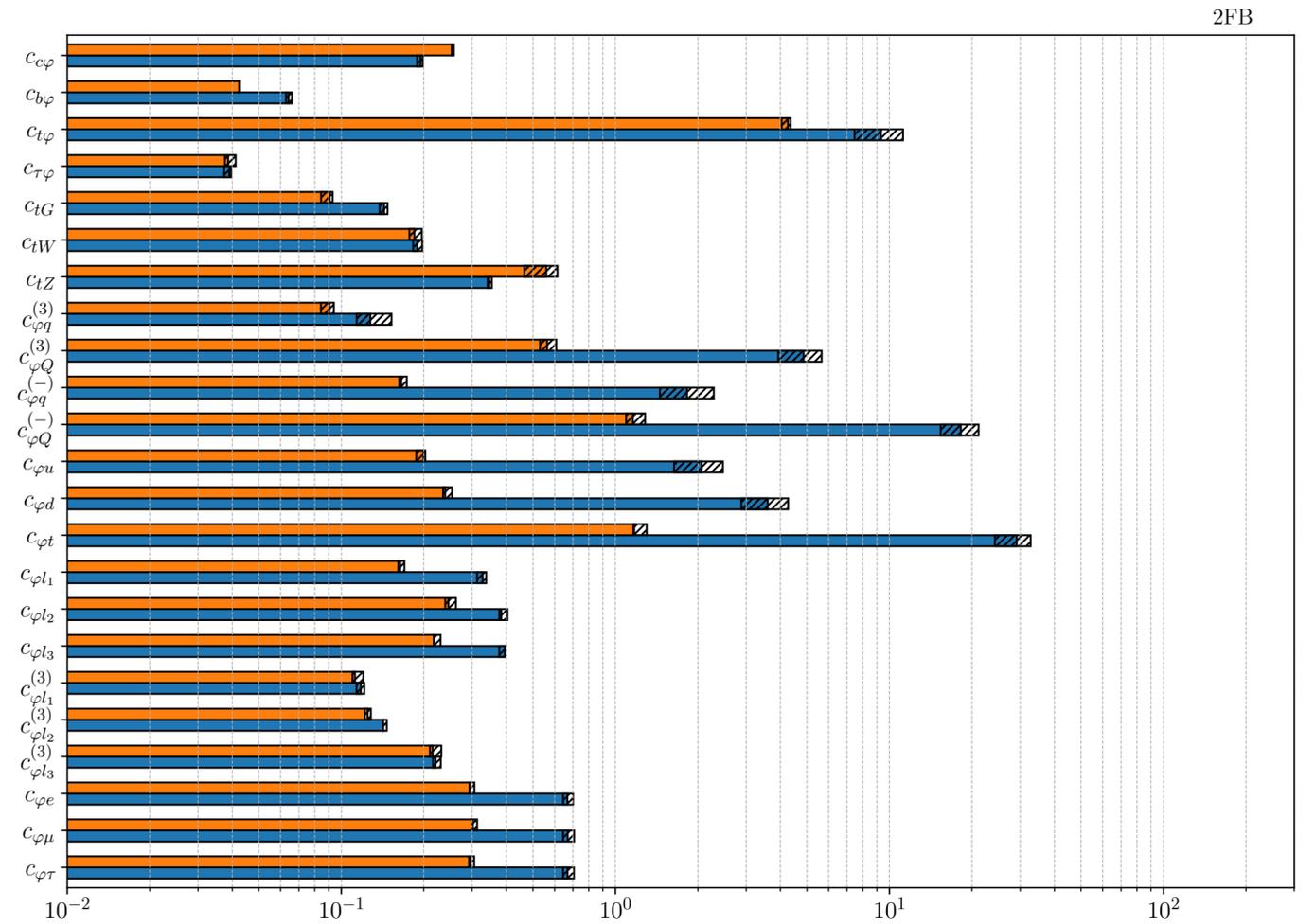
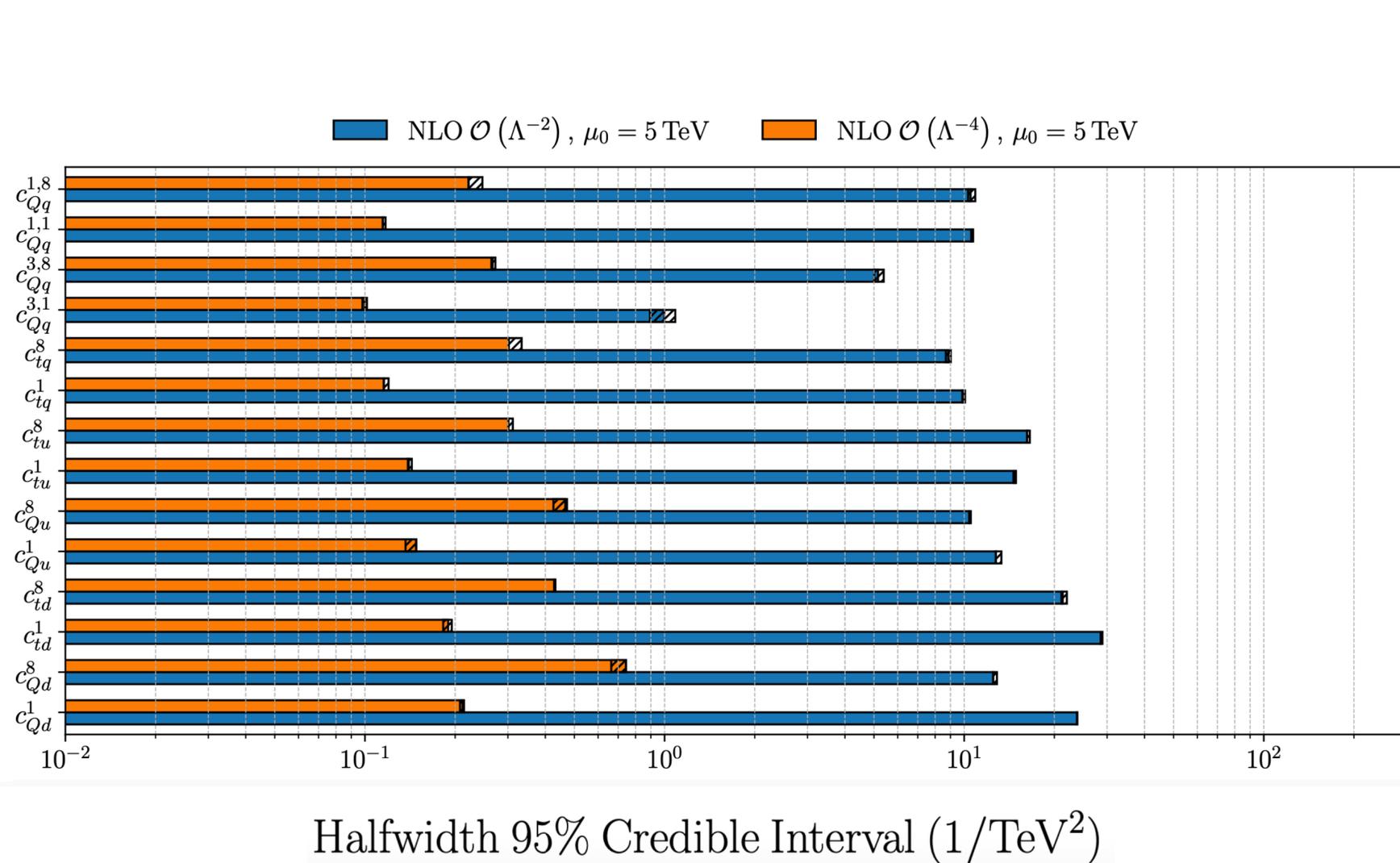
RG effects in the global fit

Global fits neglecting RGE effects can severely underestimate the bounds!



Impact of scale choice

Scale uncertainties are mild, as probed by varying the scale choice in each bin up and down by a factor 2, $\mu \rightarrow \tilde{\mu} = \kappa\mu$



SMEFT at future colliders

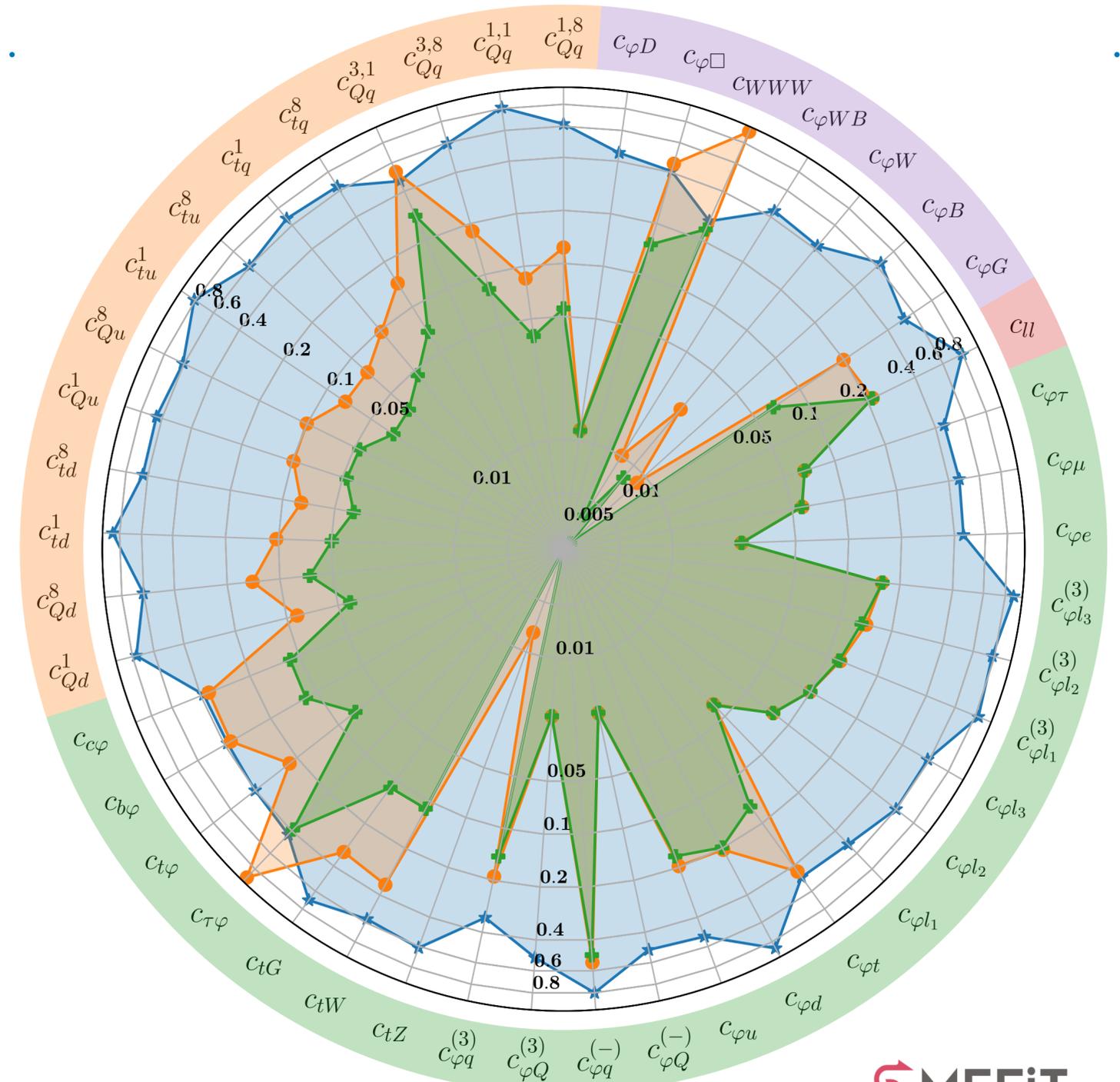


The SMEFT at HL-LHC

arXiv: 2404.12809

- ▶ We project all RunII datasets from the SMEFT3.0 baseline: one for each process and final state see backup for details
- ▶ We see an improvement ranging from 20 to 70 % in the marginalised fit
- ▶ EW operators improve only in the marginalised fit

Ratio of Uncertainties to SMEFT3.0 Baseline, $\mathcal{O}(\Lambda^{-2})$, Marginalised



MEFIT

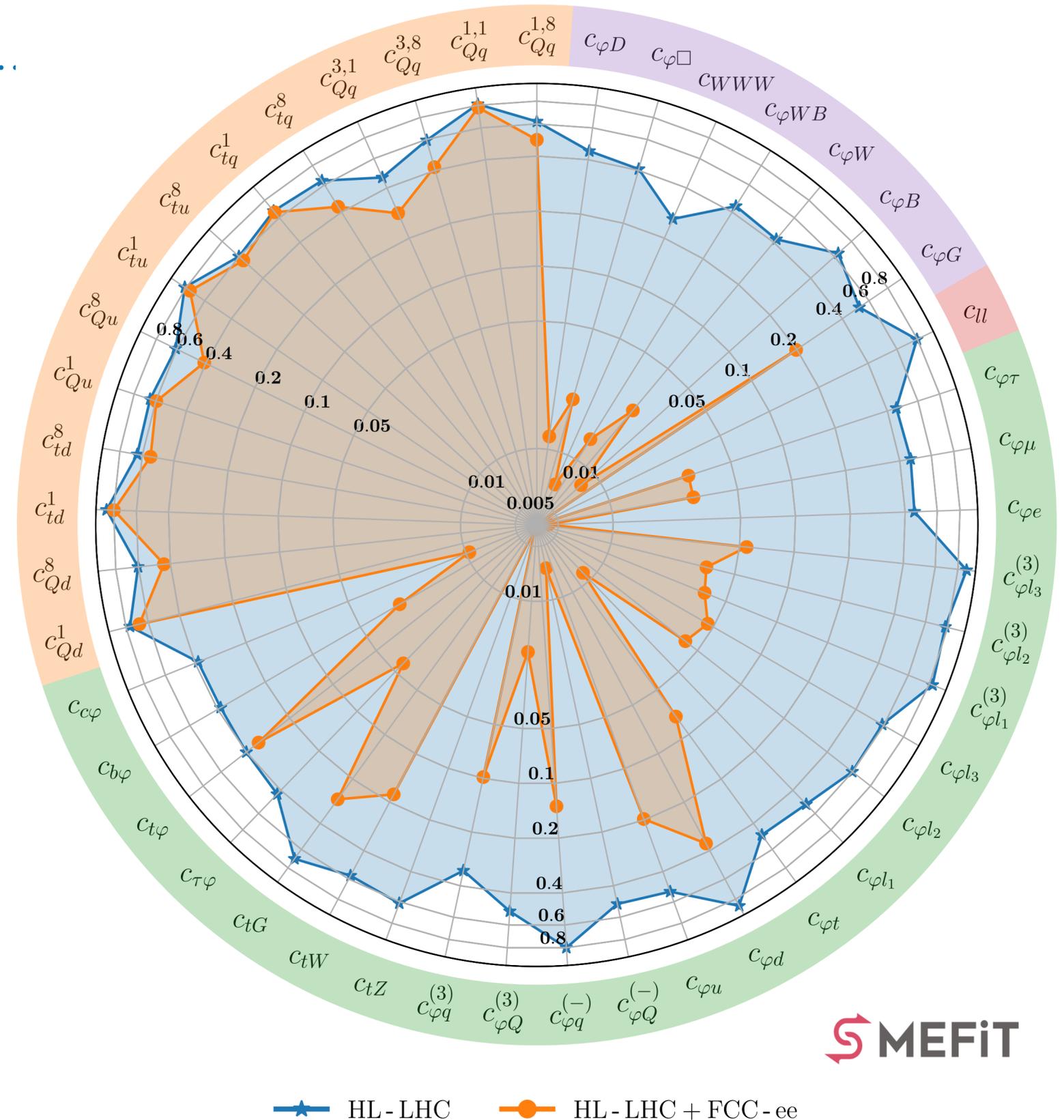
- ★ HL-LHC
- ✚ HL-LHC, individual
- SMEFT3.0, individual

The SMEFT at FCC-ee

Ratio of Uncertainties to SMEFiT3.0 Baseline, $\mathcal{O}(\Lambda^{-2})$, Marginalised

- EWPOs at the Z-pole
- Light fermion pair prediction
- Higgstrahlung and VBF
- Gauge boson pair production
- Top-quark pair production
- Optimal Observables

Energy (\sqrt{s})	\mathcal{L}_{int} (Run time)	
	FCC-ee	CEPC
91 GeV (Z-pole)	300 ab^{-1} (4 years)	100 ab^{-1} (2 years)
161 GeV ($2m_W$)	20 ab^{-1} (2 years)	6 ab^{-1} (1 year)
240 GeV	10 ab^{-1} (3 years)	20 ab^{-1} (10 years)
350 GeV	0.4 ab^{-1} (1 years)	-
365 GeV ($2m_t$)	3 ab^{-1} (4 years)	1 ab^{-1} (5 years)



The ultimate goal of the EFT program at the LHC is to bridge the gap to explicit UV models

UV model

SMEFT

Matching $\mu \leq \Lambda \simeq \min\{\mathbf{m}_{\text{UV}}\}$

$$\frac{c_{\phi\Box}}{\Lambda^2} = -\frac{g_1^4}{7680\pi^2} \frac{1}{m_\phi^2} - \frac{g_2^4}{2560\pi^2} \frac{1}{m_\phi^2} - \frac{3}{32\pi^2} \frac{\lambda_\phi^2}{m_\phi^2},$$

$$\frac{c_{t\phi}}{\Lambda^2} = -\frac{\lambda_\phi (y_\phi^u)_{33}}{m_\phi^2} - \frac{g_2^4 y_t^{\text{SM}}}{3840\pi^2} \frac{1}{m_\phi^2} + \frac{y_t^{\text{SM}} \lambda_\phi^2}{16\pi^2 m_\phi^2} + \frac{(4(y_b^{\text{SM}})^2 - 13(y_t^{\text{SM}})^2) \lambda_\phi (y_\phi^u)_{33}}{64\pi^2 m_\phi^2}$$

$$- \left(12\lambda_\phi^{\text{SM}} + (y_b^{\text{SM}})^2 - 11(y_t^{\text{SM}})^2\right) \frac{y_t^{\text{SM}} (y_\phi^u)_{33}^2}{64\pi^2 m_\phi^2} + \frac{3}{128\pi^2} \frac{\lambda_\phi (y_\phi^u)_{33}^3}{m_\phi^2},$$

$$\frac{\mathbf{c}}{\Lambda^n} = f(\mathbf{g}_{\text{UV}}, \mathbf{m}_{\text{UV}})$$

$$\sigma(\mathbf{c}) \longrightarrow \sigma(\mathbf{g}_{\text{UV}})$$

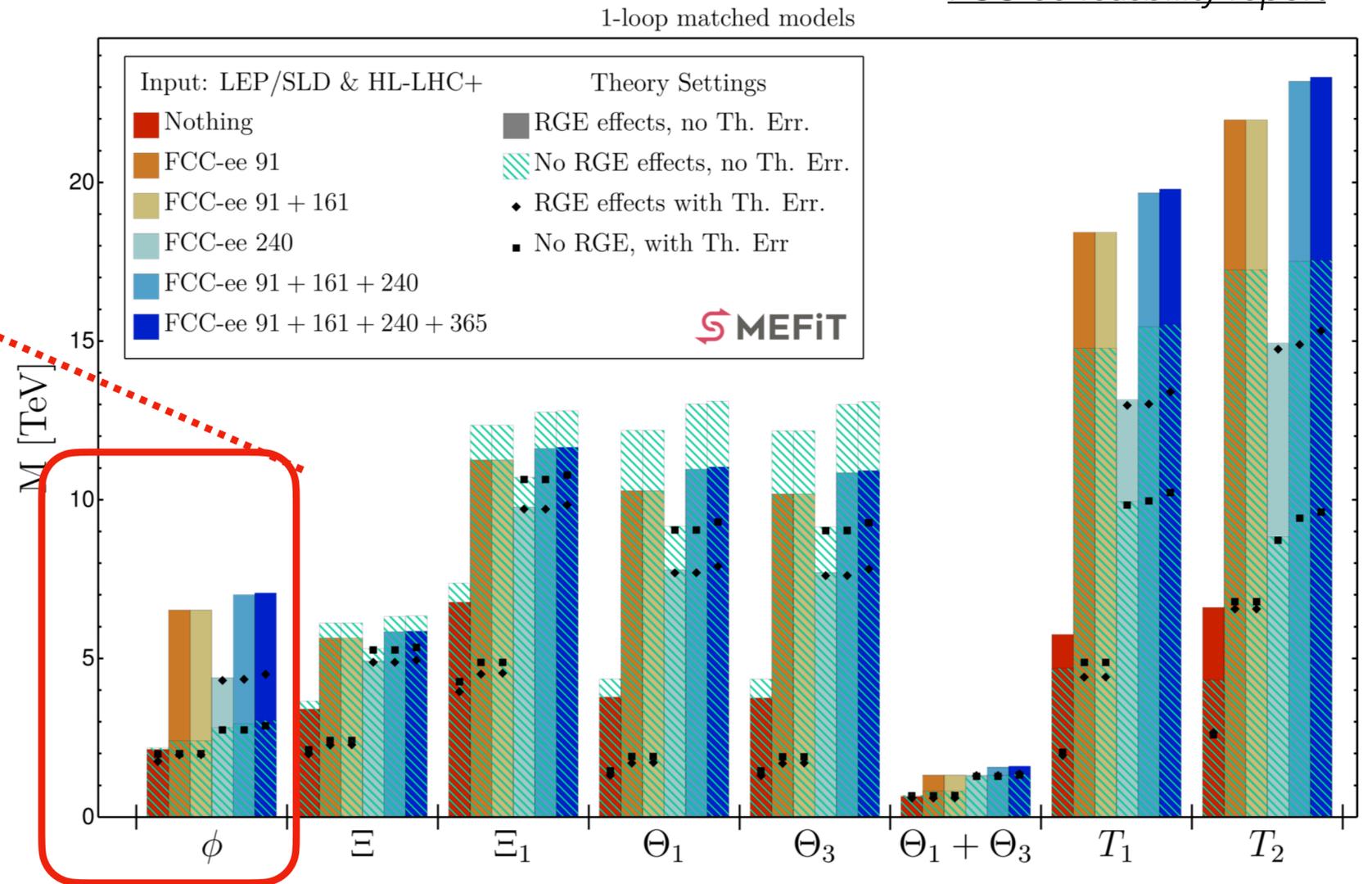
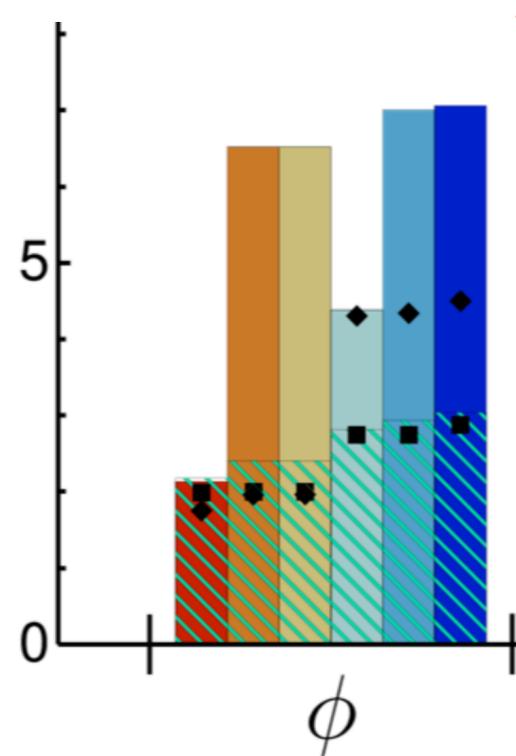
- Less parameters
- Stronger correlations
- Model dependent
- Sharper interpretation

Slide from A. Rossia at
SMEFT-tools 2025

UV complete models

RGE effects matter, and so do theory uncertainties!

FCC-ee feasibility report

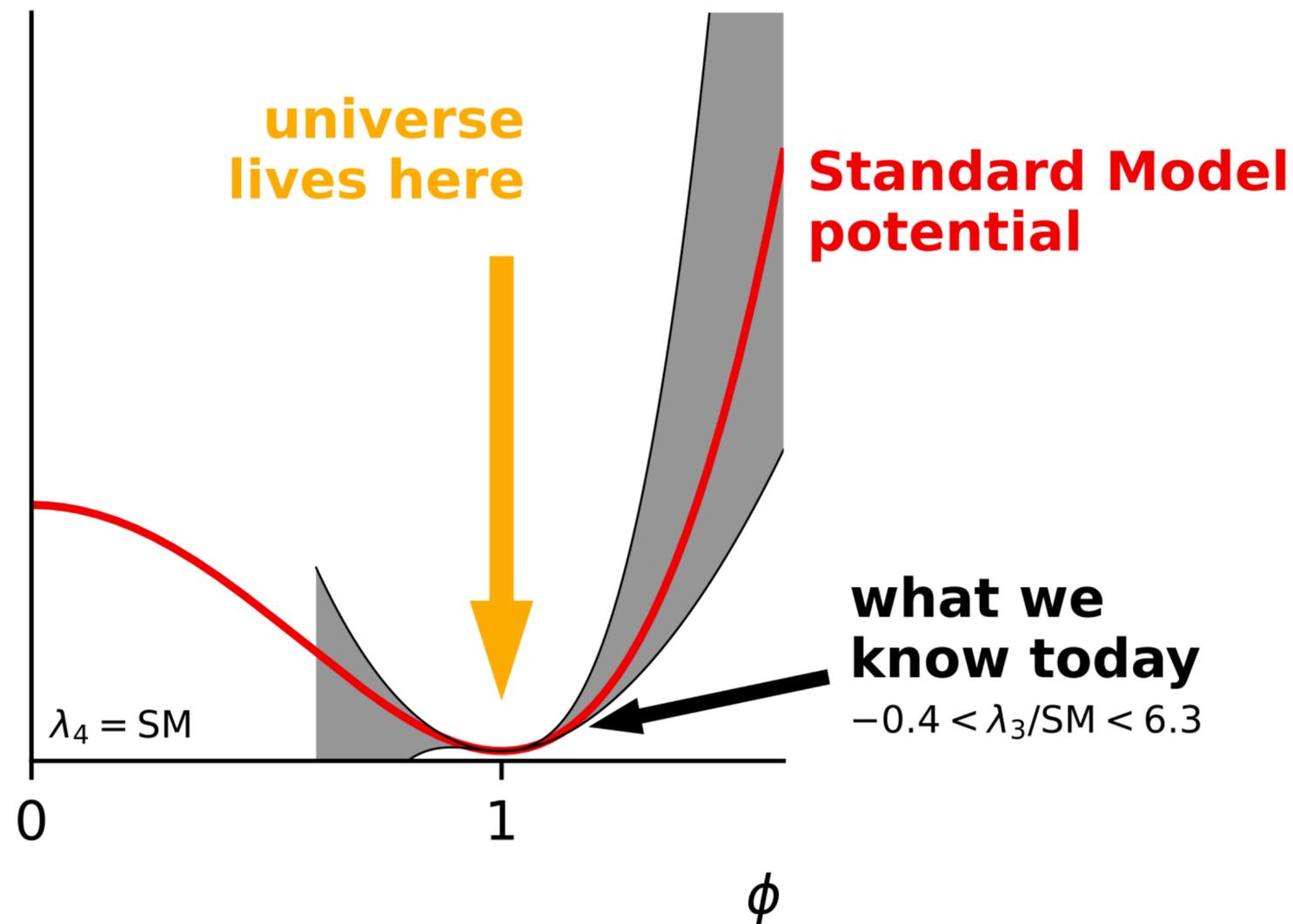


$$\mathcal{L}_{UV} = \mathcal{L}_{SM} + |D_\mu \phi|^2 - m_\phi^2 \phi^\dagger \phi - \left((y_\phi^e)_{ij} \phi^\dagger \bar{e}_R^i \ell_L^j + (y_\phi^d)_{ij} \phi^\dagger \bar{d}_R^i q_L^j + (y_\phi^u)_{ij} \phi^\dagger i\sigma_2 \bar{q}_L^{T,i} u_R^j + \lambda_\phi \phi^\dagger \phi |\phi|^2 + \text{h.c.} \right) - \text{scalar potential}$$

The Higgs self-coupling in the SMEFT

Mantani, Rojo, Rossia, Vryonidou, **JtH**
arXiv: 2504.05974, submitted to PRL

$V(\phi)$, today

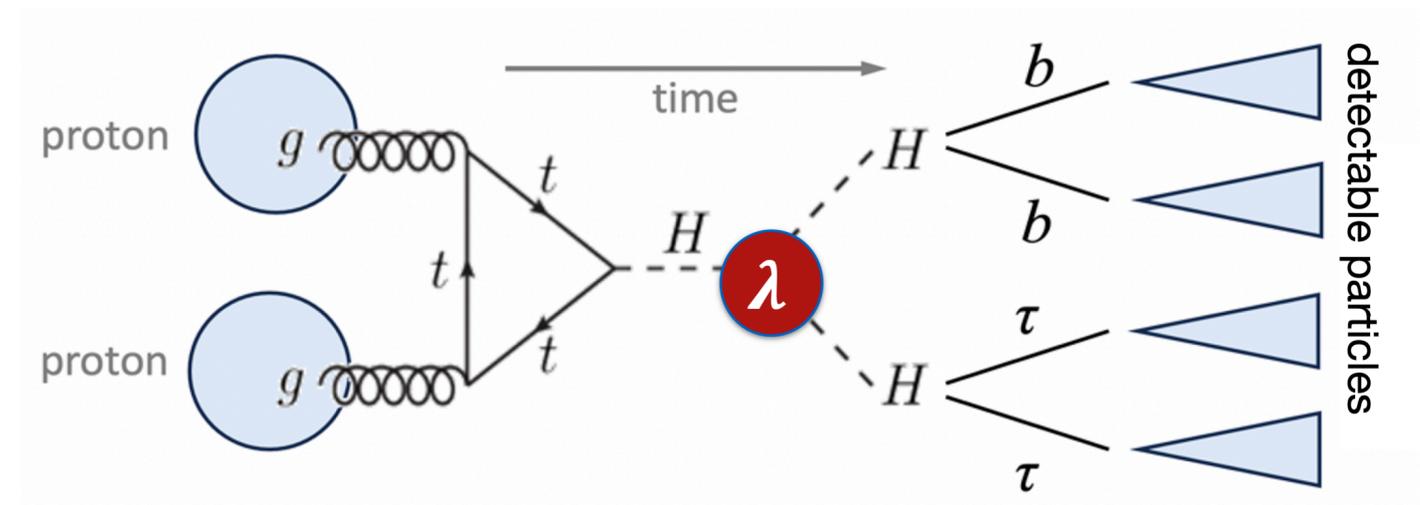


From Gavin Salam

Jaco ter Hoeve

The Higgs potential sheds light on ...

- the vacuum stability of our universe
- EW phase transition
- Matter anti-matter asymmetry

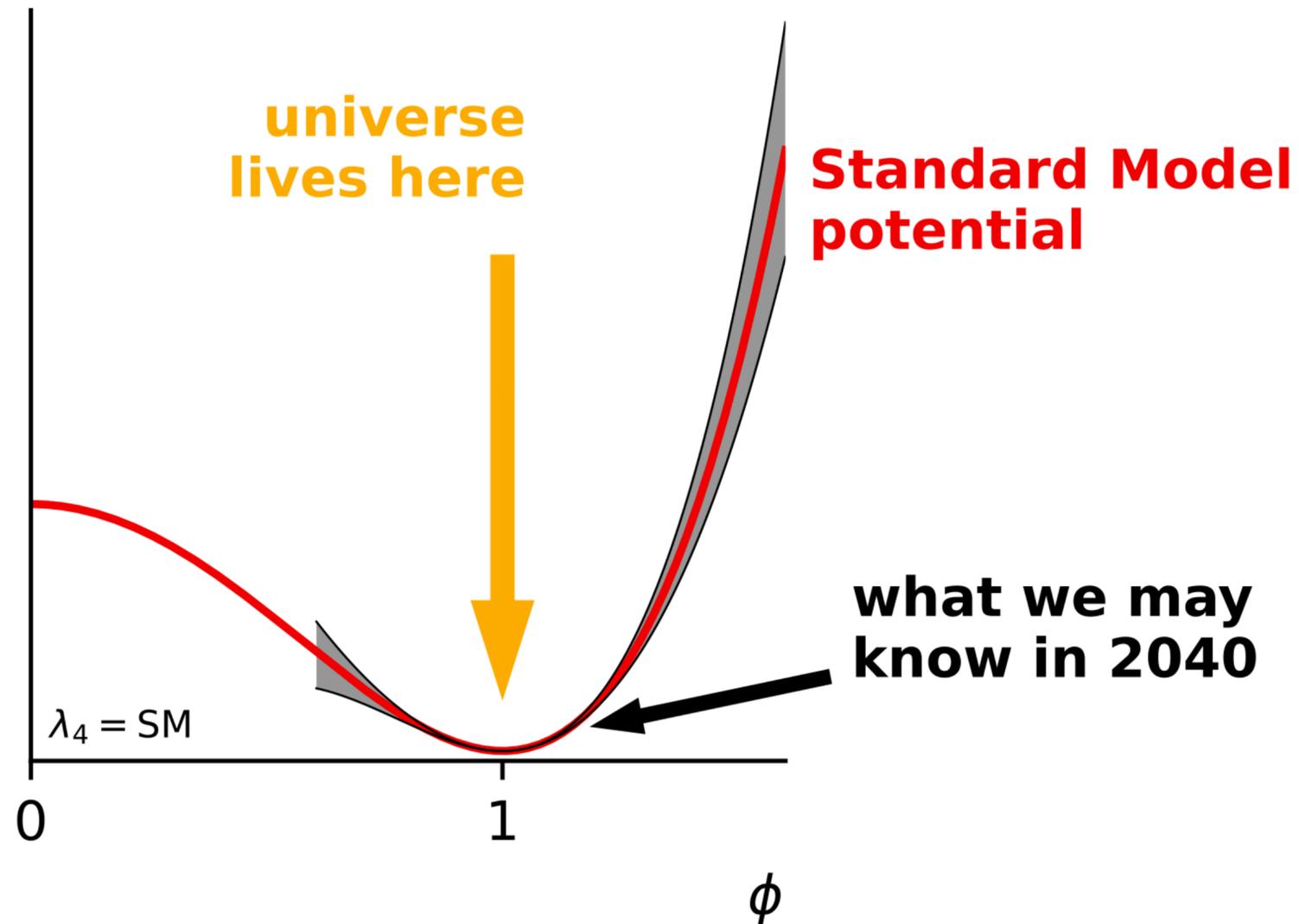


$$\delta\kappa_3 = -\frac{2v^4}{m_h^2} \frac{c_\varphi}{\Lambda^2} + \frac{3v^2}{\Lambda^2} \left(c_{\varphi\Box} - \frac{1}{4} c_{\varphi D} \right)$$

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$V(\phi)$, 2040 (HL-LHC)

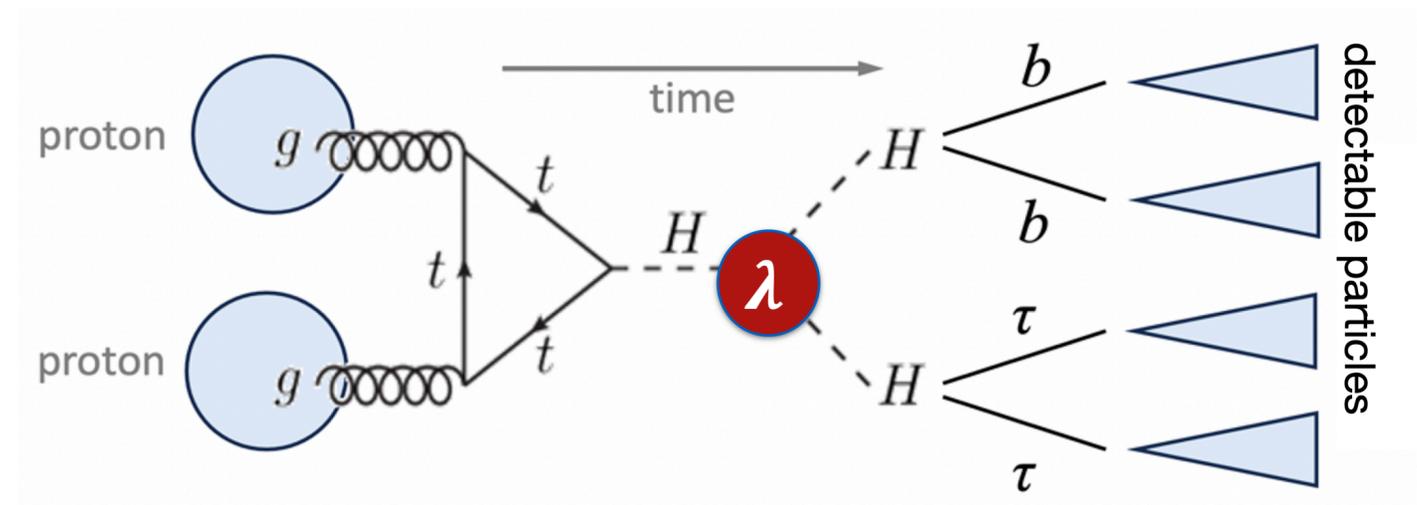


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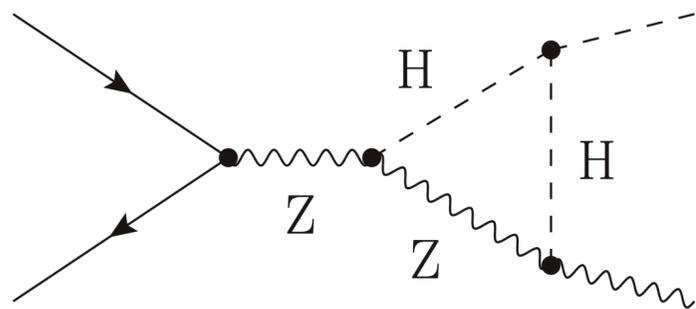
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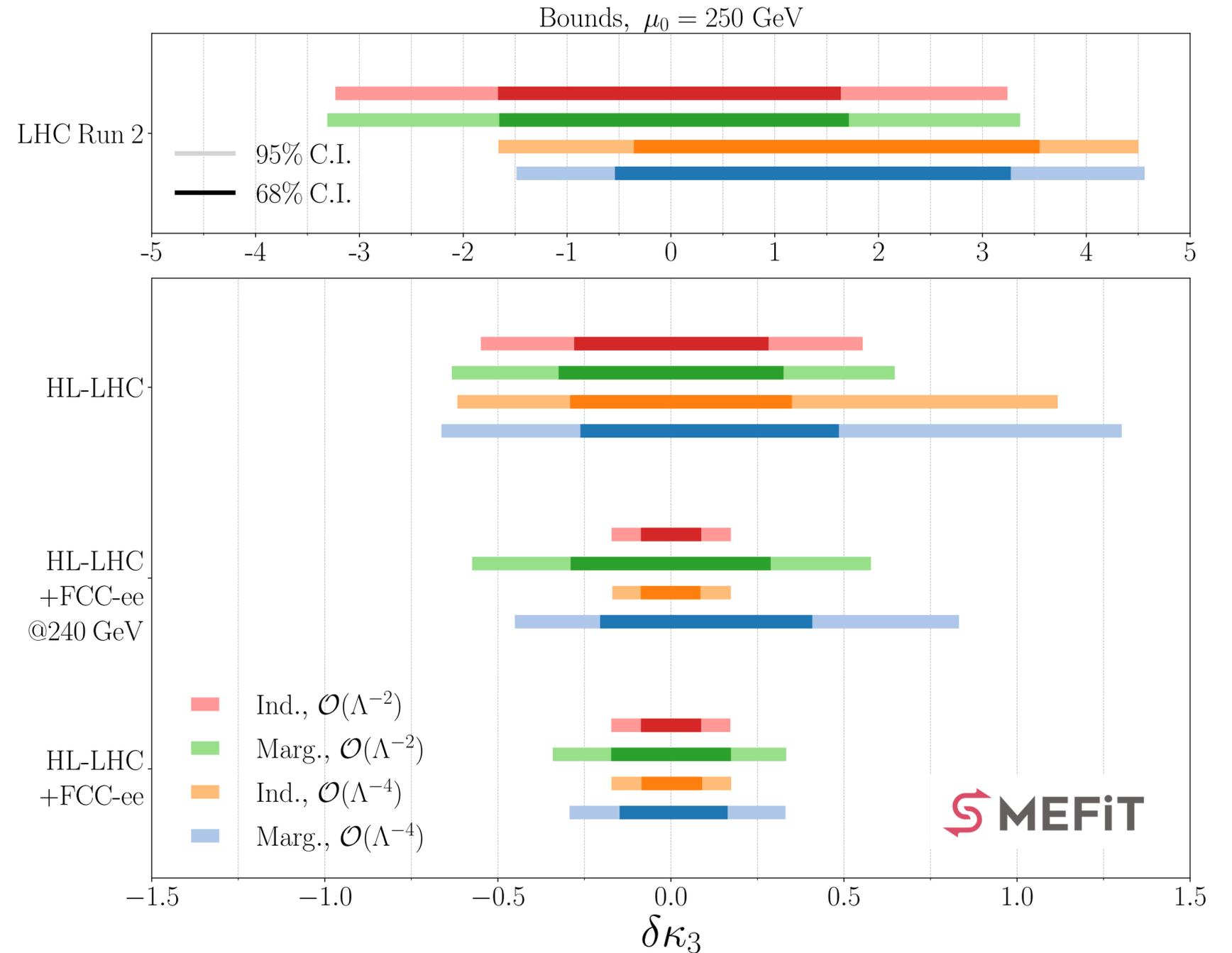
$$\delta\kappa_3 = -\frac{2v^4}{m_h^2} \frac{c_\varphi}{\Lambda^2} + \frac{3v^2}{\Lambda^2} \left(c_{\varphi\Box} - \frac{1}{4} c_{\varphi D} \right)$$

The Higgs self-coupling in the SMEFT

- ▶ HL-LHC + FCC-ee can pin down the Higgs self coupling to 15%
- ▶ Individual and marginalised bounds on κ_3 no longer agree at FCC-ee
- ▶ Significantly improving on HL-LHC is not possible without the 365 GeV run



arXiv: 2406.03557



Summary and conclusion

- The SMEFT provides a convenient tool to search for new physics in an agnostic way
- The impact of RGE effects on the SMEFT parameter space is non trivial and should be included
- Complementarity between the different FCC-ee runs is key to pin down the Higgs trilinear coupling to 15%
- The SMEFiT code is open source and can be installed from



lhcfithef.github.io/smefit_release

Summary and conclusion

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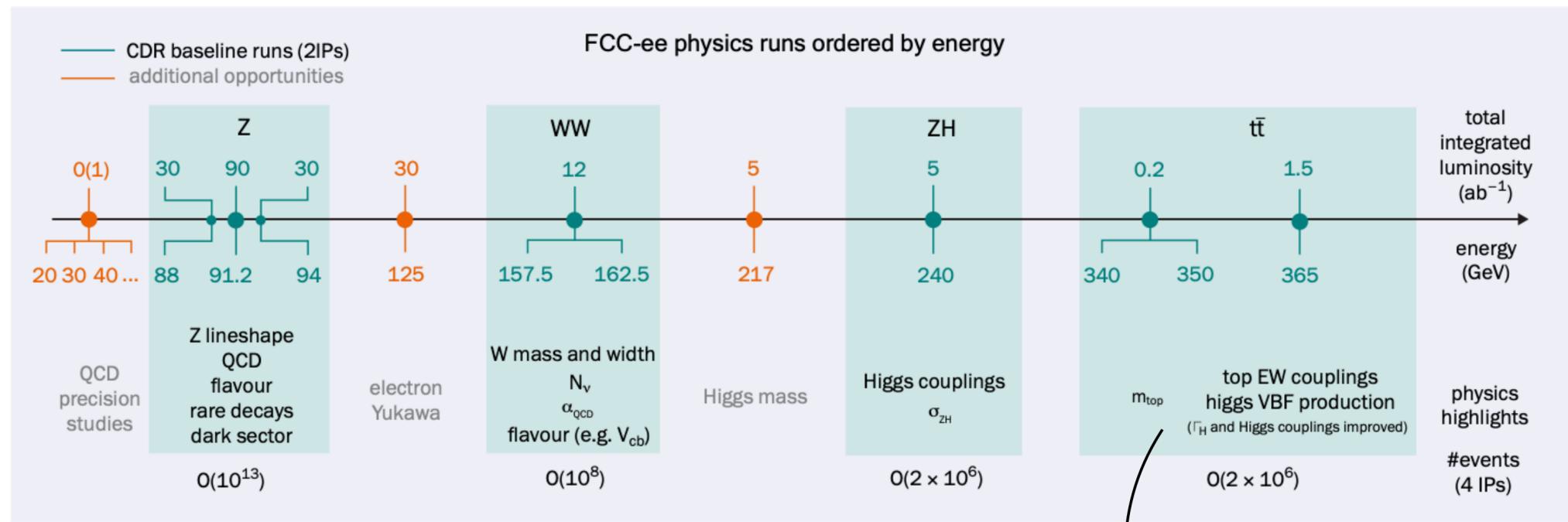
lhcfithef.github.io/smefit_release

Contact: jaco.ter.hoeve@ed.ac.uk **Thanks for your attention!**

Backup

Future colliders

- At the FCC-ee, we will separately target Z-pole EWPOs, diboson, Higgs and top pair production, each with an unprecedented luminosity!
- Expected production: $6 \cdot 10^{12}$ Z bosons, $2.4 \cdot 10^8$ W pairs, $2 \cdot 10^6$ H bosons and $2 \cdot 10^6$ top quark pairs within 16 years



Impossible to produce with an e^+e^- collider in the LHC tunnel!

Dataset upgrade

We extended SMEFiT2.0 with recent Run II datasets from top, diboson and Higgs production

Category	Processes	n_{dat}	
		SMEFiT2.0	SMEFiT3.0
Top quark production	$t\bar{t} + X$	94	115
	$t\bar{t}Z, t\bar{t}W$	14	21
	$t\bar{t}\gamma$	-	2
	single top (inclusive)	27	28
	tZ, tW	9	13
	$t\bar{t}t\bar{t}, t\bar{t}b\bar{b}$	6	12
	Total	150	189
Higgs production and decay	Run I signal strengths	22	22
	Run II signal strengths	40	40
	Run II, differential distributions & STXS	35	71
	Total	97	133
Diboson production	LEP-2	40	40
	LHC	30	41
	Total	70	81
Z-pole EWPOs	LEP-2	-	44
Baseline dataset	Total	317	449



Theory

SM: (N)NLO QCD + NLO EW

EFT: NLO QCD, linear and quadratics, with SMEFT@NLO

NNPDF4.0 no top

Data

447 measurements from **Higgs, top, diboson** and **EWPO**

Full experimental **correlations**



dataset_X.json

```
{
  "best_sm": [3.0],
  "scales": [91.0],
  "theory_cov": [[1.0]],
  "LO": {"SM": 1.0, "Op1": -0.2, "Op1*Op1": 0.4},
  "NLO_QCD": {"SM": 1.5, "Op1": -0.3, "Op1*Op1": 0.6}
}
```



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  "NLO_QCD": {"SM": 1.5, "Op1": -0.3, "Op1*Op1": 0.6}
}
```

dataset_X

```
dataset_name: ATLAS_ttW_13TeV_2016
doi: 10.1103/PhysRevD.99.072009
location: Figure 13 arxiv preprint
arxiv: 1901.03584
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units: fb
description: inclusive ttW cross-section
luminosity: 36.1
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systematics:
  - 0.2
sys_names: UNCORR
sys_type: MULT
```



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NNPDF4.0 no top

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Higgs, top, diboson and **EWPO**

Full experimental **correlations**

Projections

HL-LHC, FCCee, CEPC

+ Automatic projection module

dataset_X.json

```
{
  "best_sm": [3.0],
  "scales": [91.0],
  "theory_cov": [[1.0]],
  "LO": {"SM": 1.0, "Op1": -0.2, "Op1*Op1": 0.4},
  "NLO_QCD": {"SM": 1.5, "Op1": -0.3, "Op1*Op1": 0.6}
}
```



dataset_X

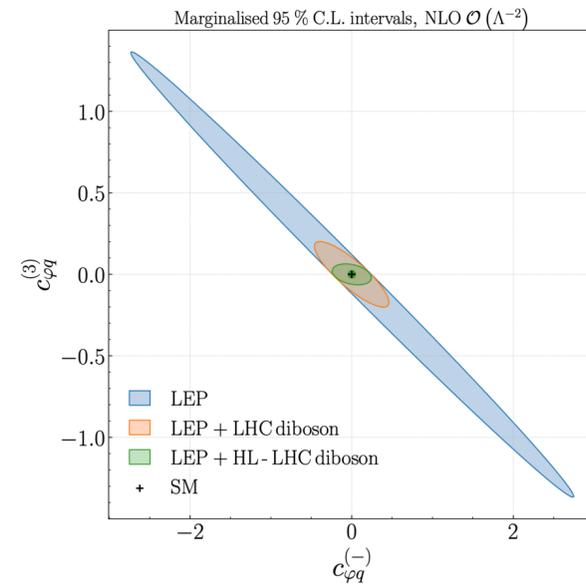
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units: fb
description: inclusive ttW cross-section
luminosity: 36.1
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  - 0.2
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```

Analysis tools

Report module

Automatised fit reports that analyse the SMEFiT results

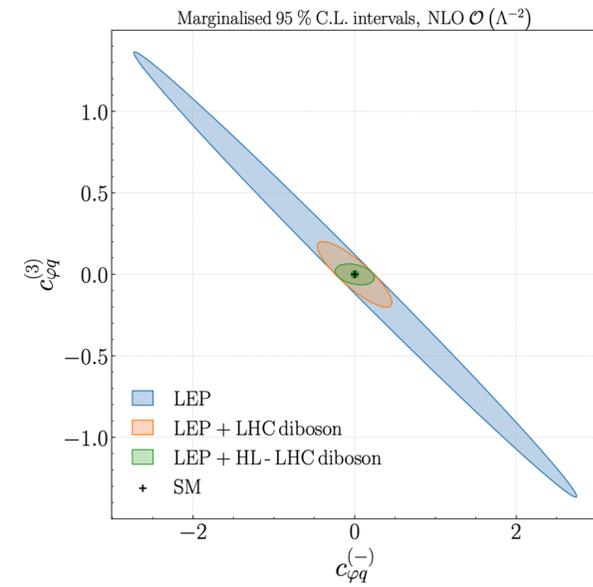
Analysis tools



Report module

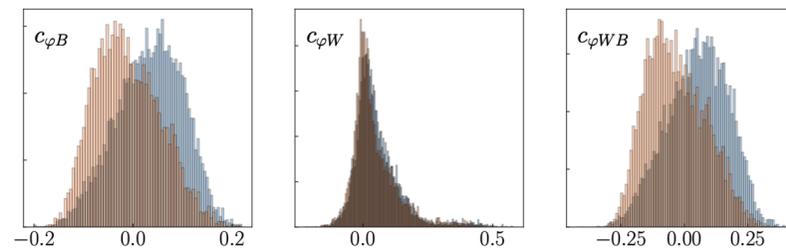
Automatised fit reports that analyse the SMEFiT results

Analysis tools

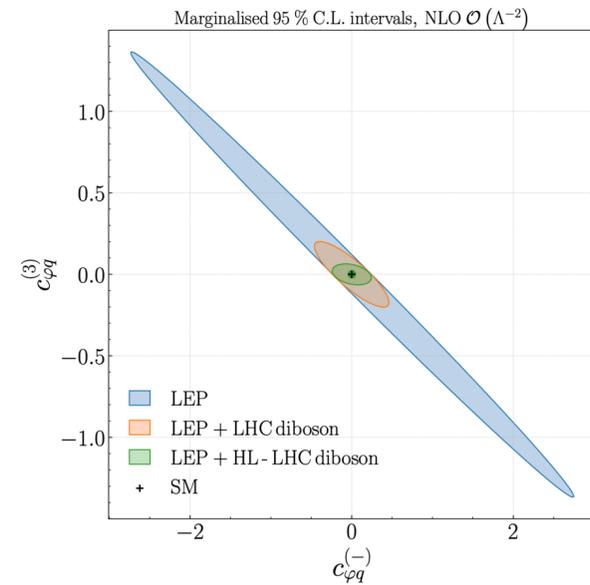


Report module

Automatised fit reports that analyse the SMEFiT results

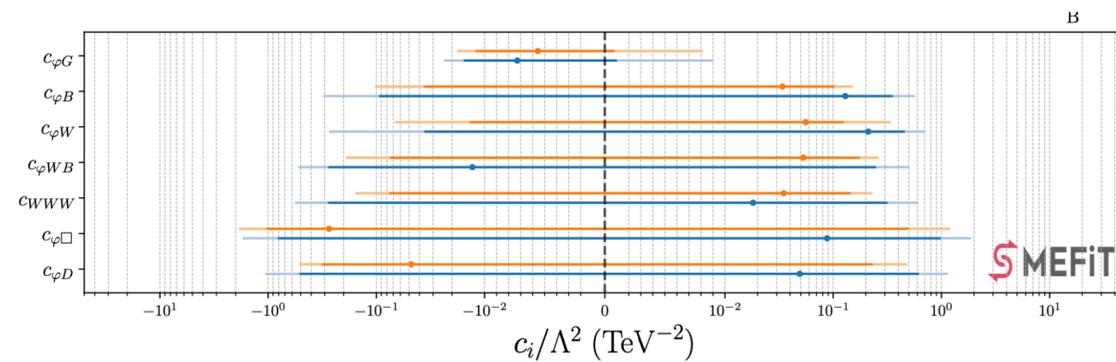
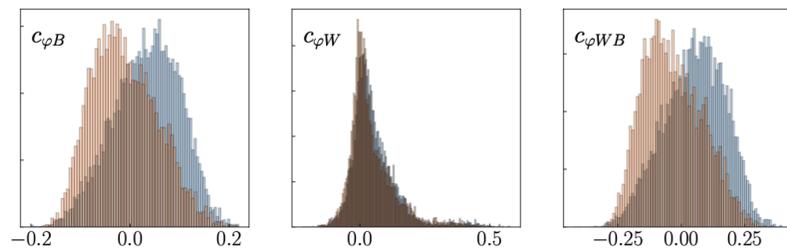


Analysis tools

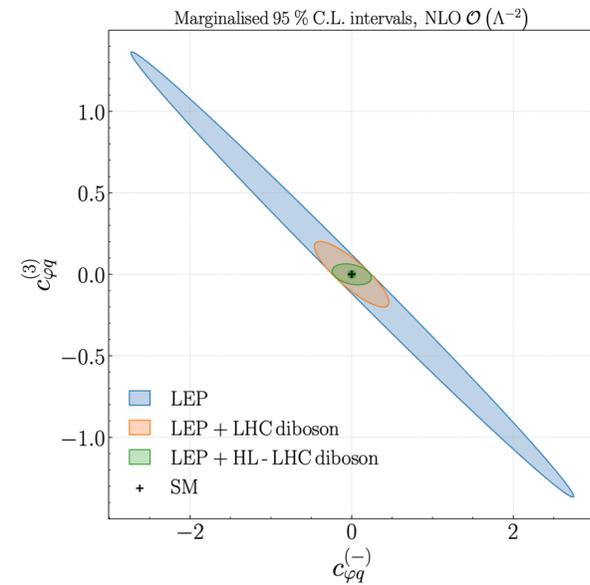


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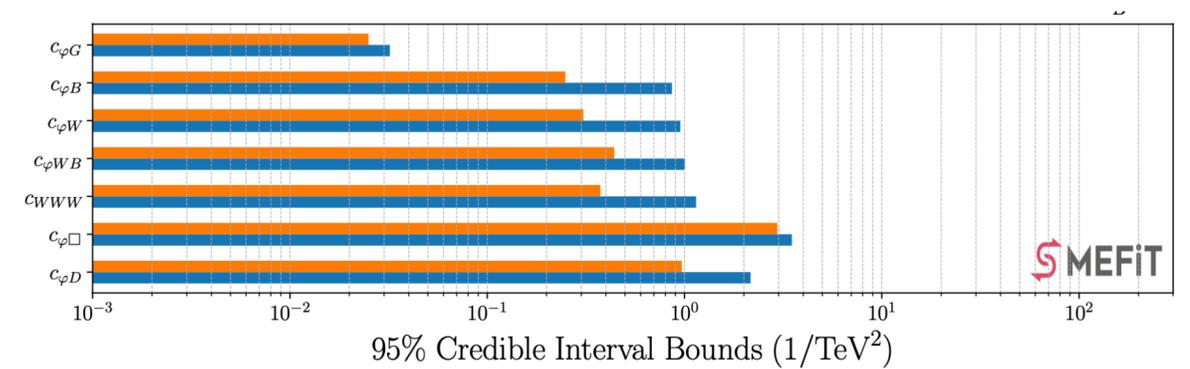
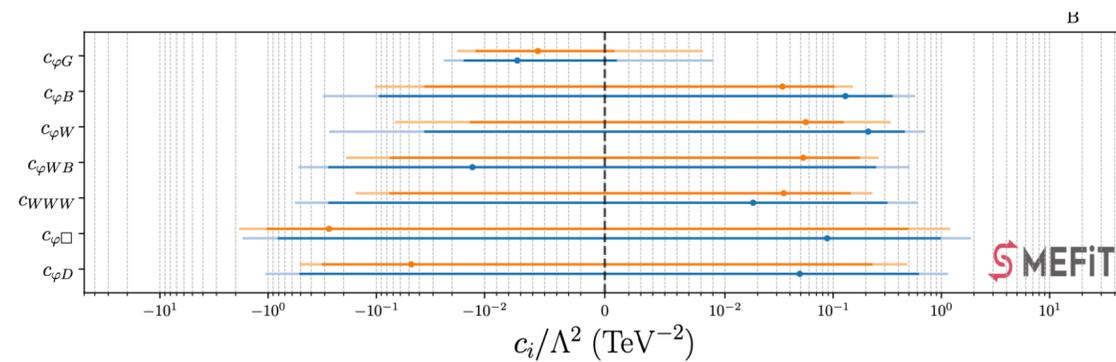
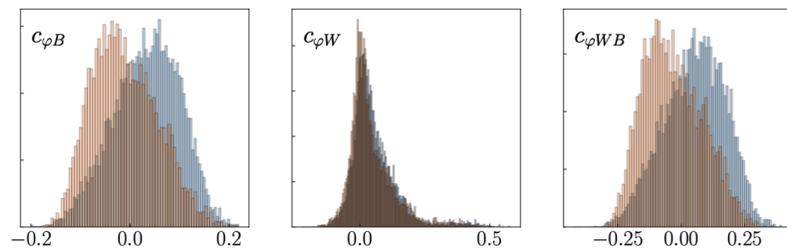


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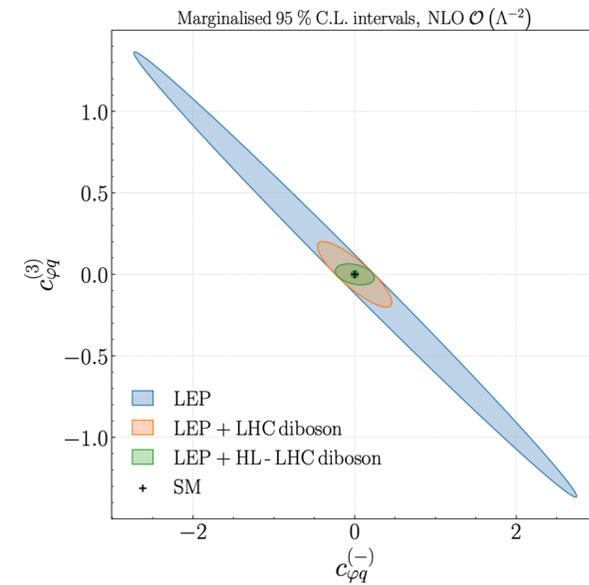


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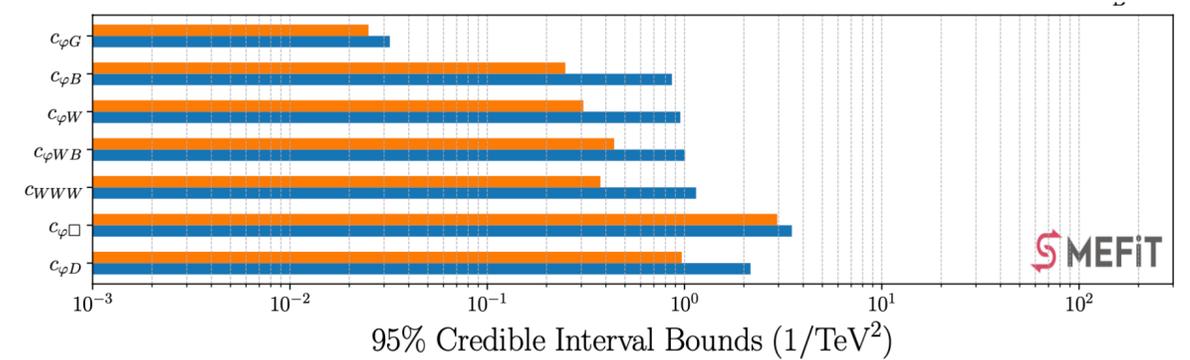
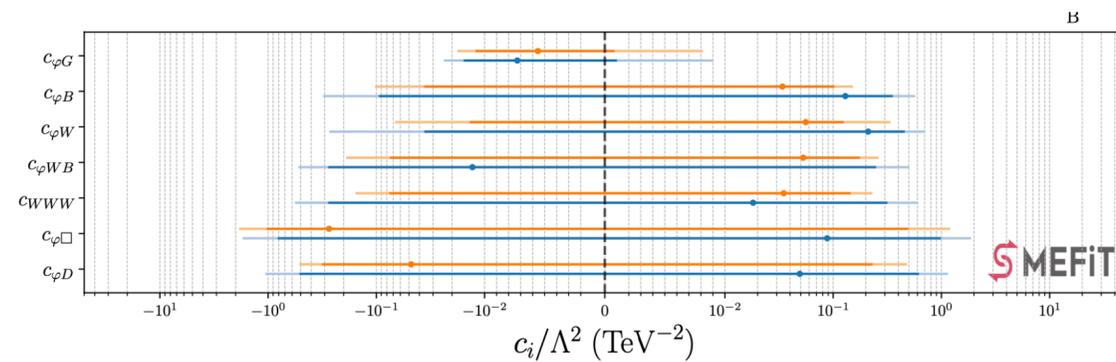
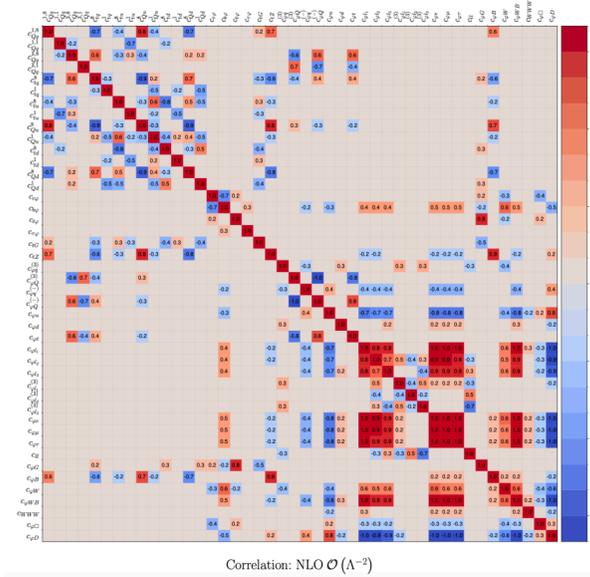
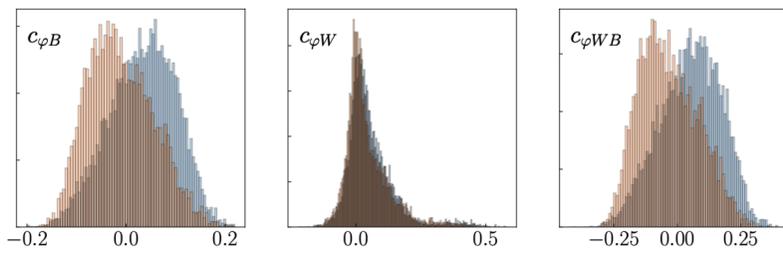


Analysis tools

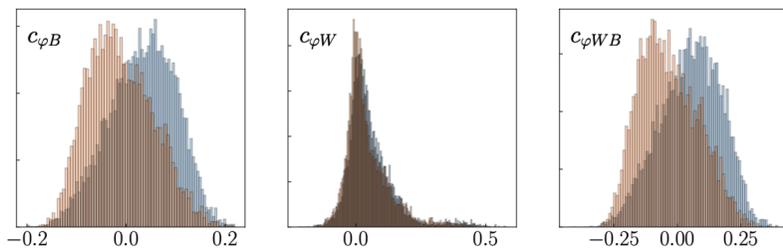
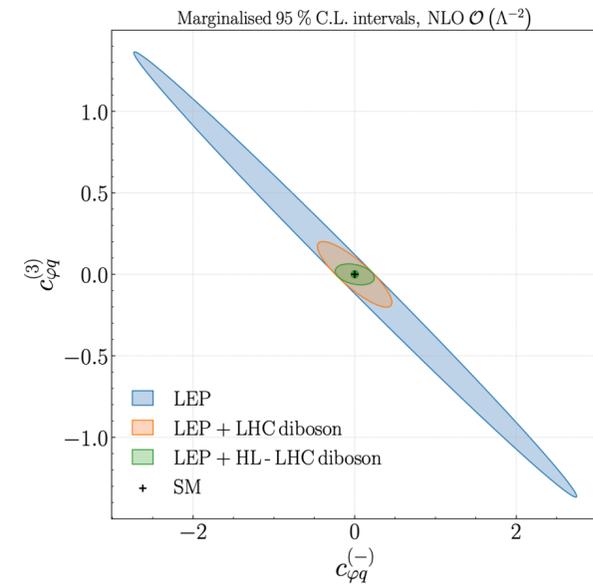


Report module

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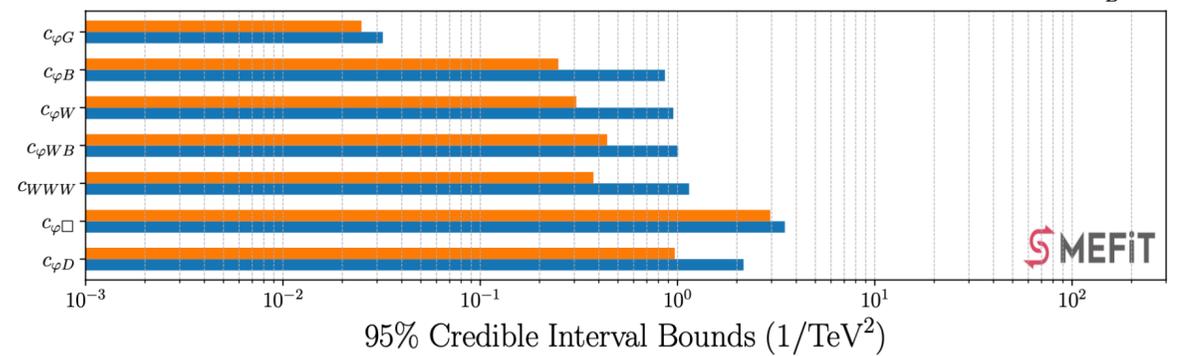
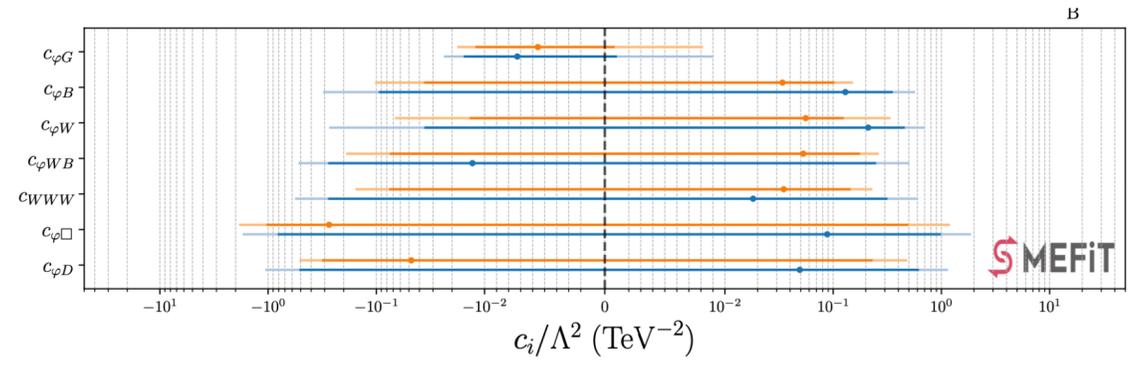
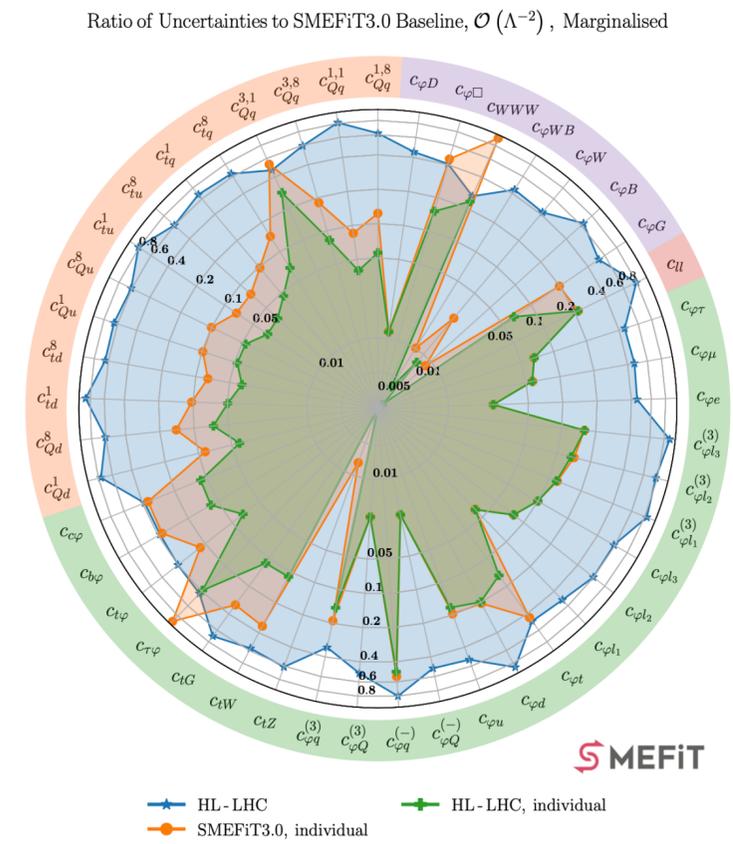
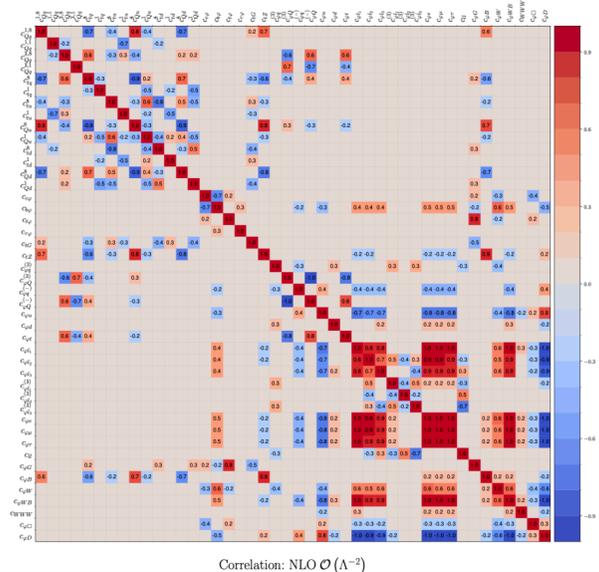


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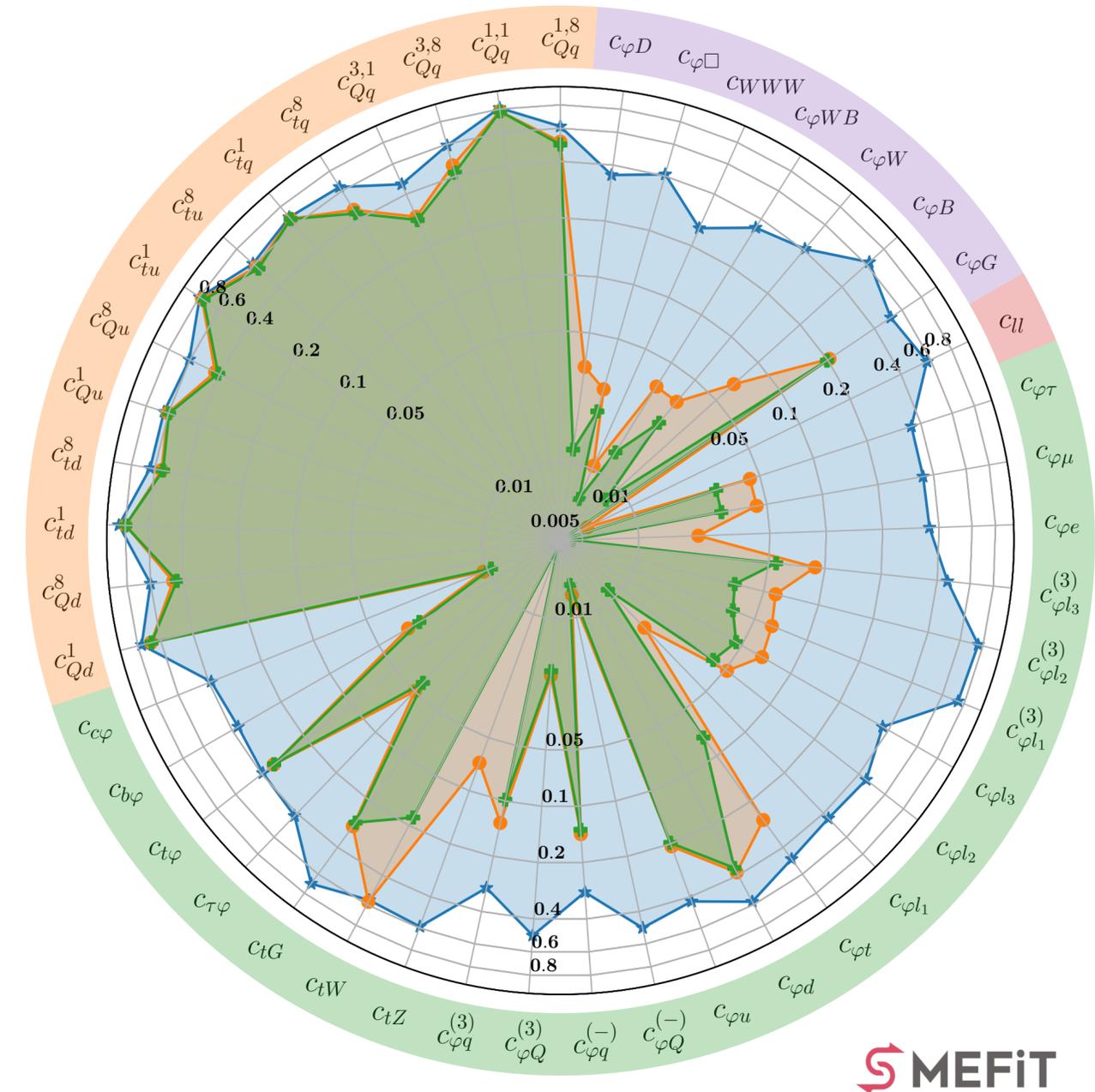
Automatised fit reports that analyse the SMEFiT results



Result: FCC-ee energy breakdown

Ratio of Uncertainties to SMEFIT3.0 Baseline, $\mathcal{O}(\Lambda^{-2})$, Marginalised

- ▶ The FCC-ee plans to operate **sequentially**, hence we need to study the impact at the various energies
- ▶ Largest impact for Z-pole at 91 GeV plus the Higgs factory run at 240 GeV
- ▶ We can try other combinations too in order to find the most optimal run order for the SMEFT

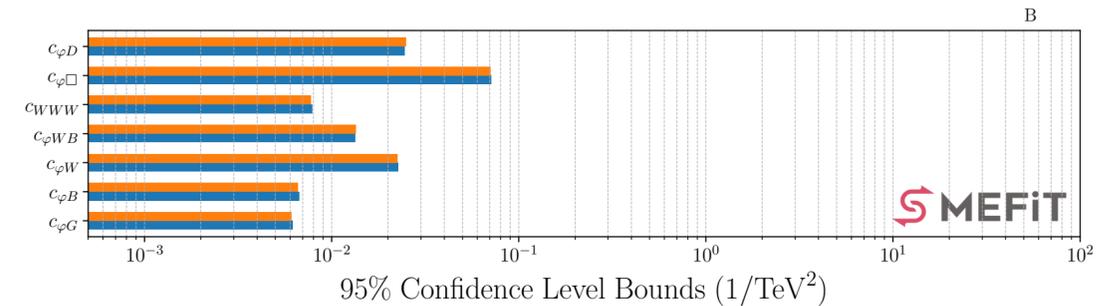
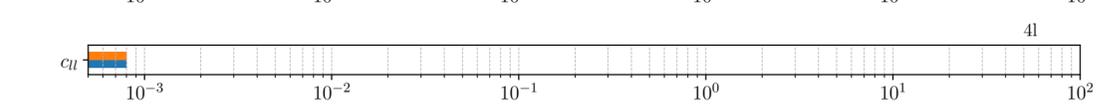
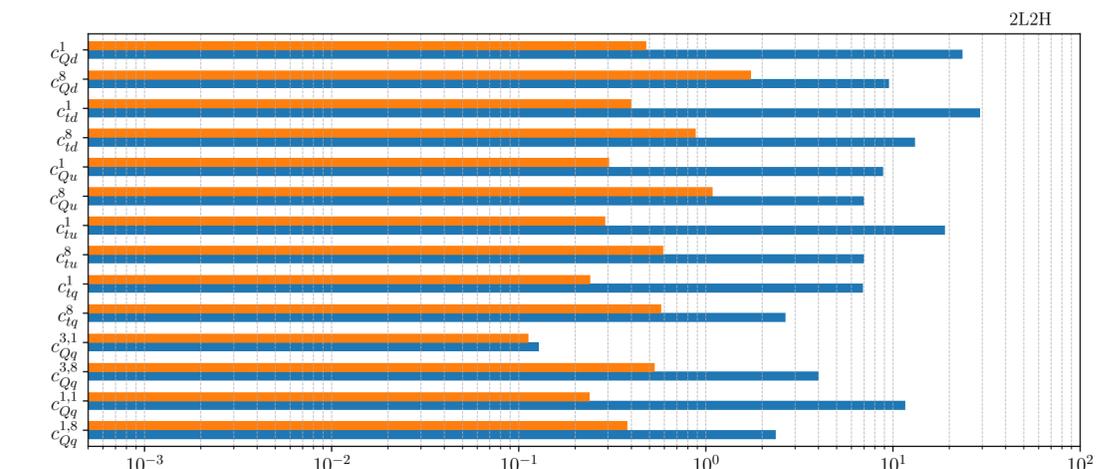
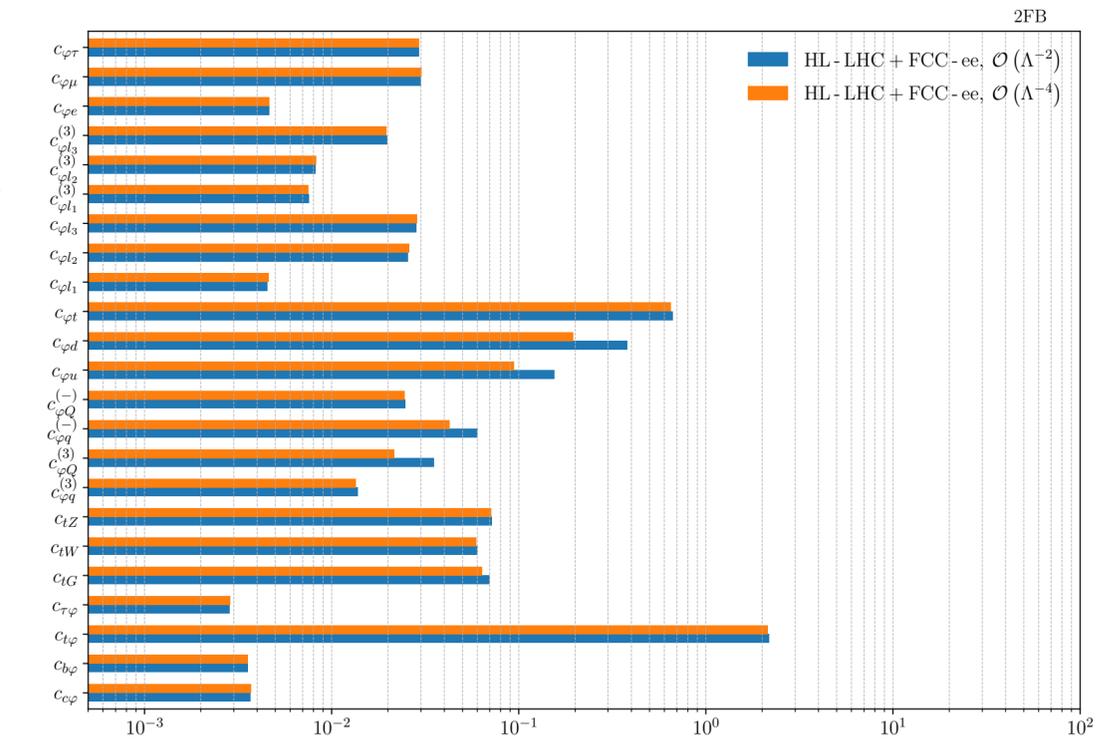
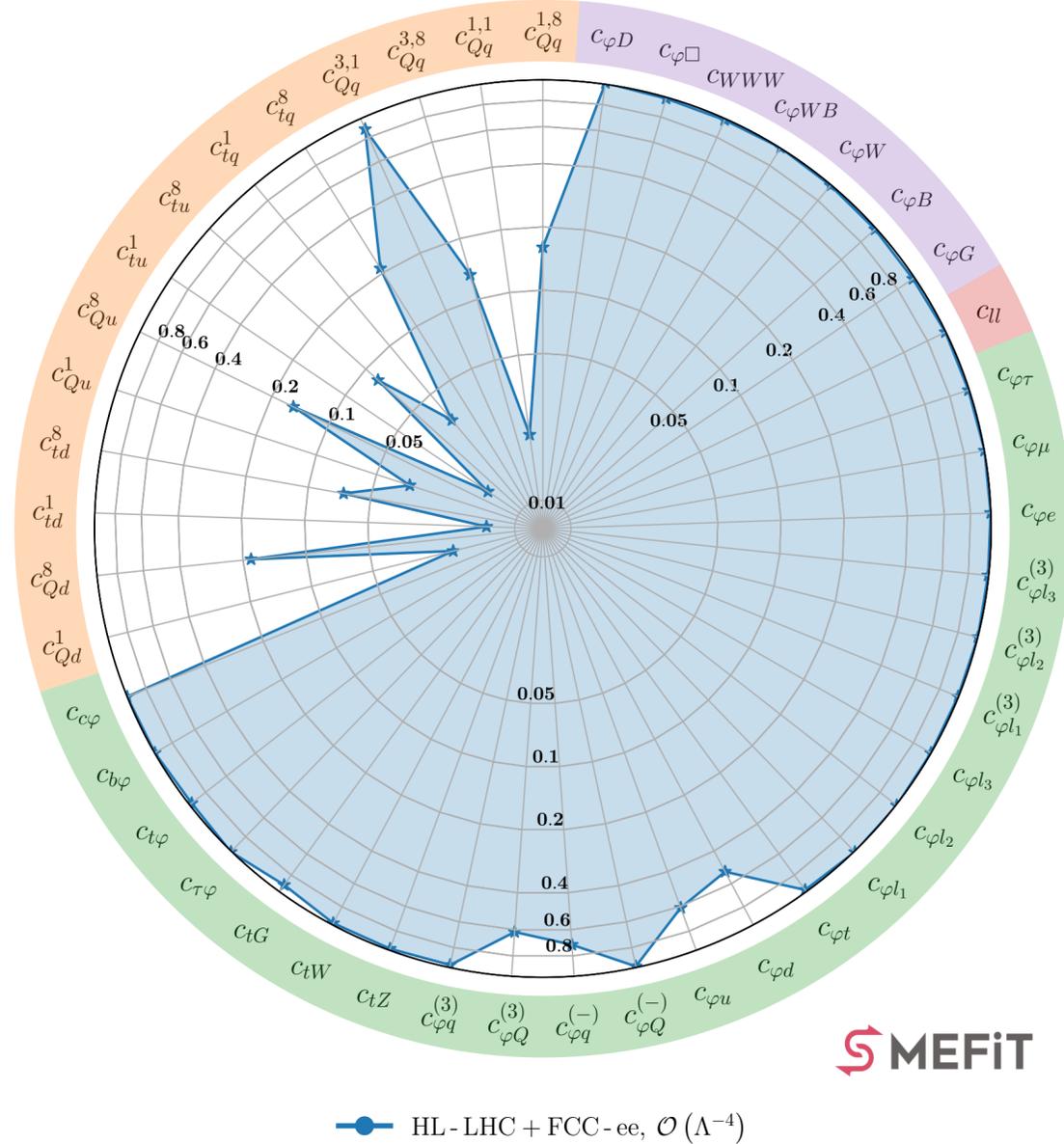


 MEFIT

- ★ HL-LHC + FCC-ee (91 GeV)
- ✚ HL-LHC + FCC-ee (91 + 161 + 240 + 365 GeV)
- HL-LHC + FCC-ee (91 + 240 GeV)

Impact of quadratics

Ratio of Uncertainties to HL - LHC + FCC-ee, $\mathcal{O}(\Lambda^{-2})$, Marginalised



Conclusion and outlook

- Presented **SMEFiT3.0**, a global fit of 50 Wilson coefficients to Higgs, top, diboson and EWPOs, including **quadratic corrections**
- We are becoming **increasingly sensitive** to possible new physics effects, both through (still) expanding LHC datasets, as well as through future colliders
- The FCC-ee offers an **unprecedented indirect mass reach** on new heavy particles
- RGE effects are crucial to include to connect experiments at widely separated scales
- Outlook: the inclusion of other proposed future colliders

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Contact: jaco.ter.hoeve@ed.ac.uk

Thanks for your attention!

SM predictions

Category	Process	SM	Code/Ref	SMEFT
Top quark production	$t\bar{t}$ (incl)	NNLO QCD	MG5_aMC NLO + NNLO K -fact	NLO QCD
	$t\bar{t} + V$	NLO QCD	MG5_aMC NLO	LO QCD + NLO SM K -fact
	single- t (incl)	NNLO QCD	MG5_aMC NLO + NNLO K -fact	NLO QCD
	$t + V$	NLO QCD	MG5_aMC NLO	LO QCD + NLO SM K -fact
	$t\bar{t}t, t\bar{t}b$	NLO QCD	MG5_aMC NLO	LO QCD + NLO SM K -fact
Higgs production and decay	$gg \rightarrow h$	NNLO QCD + NLO EW	HXSWG	NLO QCD
	VBF	NNLO QCD + NLO EW	HXSWG	LO QCD
	$h + V$	NNLO QCD + NLO EW	HXSWG	NLO QCD
	$h t\bar{t}$	NNLO QCD + NLO EW	HXSWG	NLO QCD
	$h \rightarrow X$	NNLO QCD + NLO EW	HXSWG	NLO QCD ($X = b\bar{b}$) LO QCD ($X \neq b\bar{b}$)
Diboson production	$e^+e^- \rightarrow W^+W^-$	NNLO QCD + NLO EW	LEP EWWG	LO QCD
	$pp \rightarrow VV'$	NNLO QCD	MATRIX	NLO QCD

HL-LHC projected datasets

Dataset	\mathcal{L} (fb ⁻¹)	Info	Observables	n_{dat}	Ref.
ATLAS_STXS_RunII_13TeV_2022	139	$ggF, \text{VBF}, Vh, t\bar{t}h, th$	$d\sigma/dp_T^h$ $d\sigma/dm_{jj}$ $d\sigma/dp_T^V$	36	[55]
CMS_ggF_aa_13TeV	77.4	$ggF, h \rightarrow \gamma\gamma$	$\sigma_{ggF}(p_T^h, N_{\text{jets}})$	6	[83]
ATLAS_ggF_ZZ_13TeV	79.8	$ggF, h \rightarrow ZZ$	$\sigma_{ggF}(p_T^h, N_{\text{jets}})$	6	[84]
ATLAS_ggF_13TeV_2015	36.1	$ggF, h \rightarrow ZZ, h \rightarrow \gamma\gamma$	$d\sigma(ggF)/dp_T^h$	9	[85]
ATLAS_WH_Hbb_13TeV	79.8	$Wh, h \rightarrow b\bar{b}$	$d\sigma^{(\text{fid})}/dp_T^W$ (stage 1 STXS)	2	[86]
ATLAS_ZH_Hbb_13TeV	79.8	$Zh, h \rightarrow b\bar{b}$	$d\sigma^{(\text{fid})}/dp_T^Z$ (stage 1 STXS)	2	[86]
CMS_H_13TeV_2015_pTH	35.9	$h \rightarrow b\bar{b}, h \rightarrow \gamma\gamma, h \rightarrow ZZ$	$d\sigma/dp_T^h$	9	[87]
ATLAS_WW_13TeV_2016_memu	36.1	fully leptonic	$d\sigma^{(\text{fid})}/dm_{e\mu}$	13	[88]
ATLAS_WZ_13TeV_2016_mTWZ	36.1	fully leptonic	$d\sigma^{(\text{fid})}/dm_T^{WZ}$	6	[89]
CMS_WZ_13TeV_2016_pTZ	35.9	fully leptonic	$d\sigma^{(\text{fid})}/dp_T^Z$	11	[90]
CMS_WZ_13TeV_2022_pTZ	137	fully leptonic	$d\sigma/dp_T^Z$	11	[56]

Dataset	\mathcal{L} (fb ⁻¹)	Info	Observables	n_{dat}	Ref.
ATLAS_tt_13TeV_ljets_2016_Mtt	36.1	ℓ +jets	$d\sigma/dm_{t\bar{t}}$	7	[91]
CMS_tt_13TeV_dilep_2016_Mtt	35.9	dilepton	$d\sigma/dm_{t\bar{t}}$	7	[92]
CMS_tt_13TeV_Mtt	137	ℓ +jets	$1/\sigma d\sigma/dm_{t\bar{t}}$	14	[57]
CMS_tt_13TeV_ljets_inc	137	ℓ +jets	$\sigma(t\bar{t})$	1	[57]
ATLAS_tt_13TeV_asy_2022	139	ℓ + jets	A_C	5	[59]
CMS_tt_13TeV_asy	138	ℓ + jets	A_C	3	[58]
ATLAS_Whel_13TeV	139	W -helicity fraction	F_0, F_L	2	[60]
ATLAS_ttbb_13TeV_2016	36.1	lepton + jets	$\sigma_{\text{tot}}(t\bar{t}b\bar{b})$	1	[93]
CMS_ttbb_13TeV_2016	35.9	all-jets	$\sigma_{\text{tot}}(t\bar{t}b\bar{b})$	1	[94]
CMS_ttbb_13TeV_dilepton_inc	35.9	dilepton	$\sigma_{\text{tot}}(t\bar{t}b\bar{b})$	1	[68]
CMS_ttbb_13TeV_ljets_inc	35.9	lepton + jets	$\sigma_{\text{tot}}(t\bar{t}b\bar{b})$	1	[68]
ATLAS_tttt_13TeV_run2	139	multi-lepton	$\sigma_{\text{tot}}(t\bar{t}t\bar{t})$	1	[95]
CMS_tttt_13TeV_run2	137	same-sign or multi-lepton	$\sigma_{\text{tot}}(t\bar{t}t\bar{t})$	1	[96]
ATLAS_tttt_13TeV_slep_inc	139	single-lepton	$\sigma_{\text{tot}}(t\bar{t}t\bar{t})$	1	[64]
CMS_tttt_13TeV_slep_inc	35.8	single-lepton	$\sigma_{\text{tot}}(t\bar{t}t\bar{t})$	1	[65]
ATLAS_tttt_13TeV_2023	139	multi-lepton	$\sigma_{\text{tot}}(t\bar{t}t\bar{t})$	1	[66]
CMS_tttt_13TeV_2023	139	same-sign or multi-lepton	$\sigma_{\text{tot}}(t\bar{t}t\bar{t})$	1	[67]
CMS_ttZ_13TeV_pTZ	77.5	$t\bar{t}Z$	$d\sigma(t\bar{t}Z)/dp_T^Z$	4	[97]
ATLAS_ttZ_13TeV_pTZ	139	$t\bar{t}Z$	$d\sigma(t\bar{t}Z)/dp_T^Z$	7	[61]
ATLAS_ttW_13TeV_2016	36.1	$t\bar{t}W$	$\sigma_{\text{tot}}(t\bar{t}W)$	1	[98]
CMS_ttW_13TeV	35.9	$t\bar{t}W$	$\sigma_{\text{tot}}(t\bar{t}W)$	1	[99]
ATLAS_t_tch_13TeV_inc	3.2	t -channel	$\sigma_{\text{tot}}(tq), \sigma_{\text{tot}}(\bar{t}q)$	2	[100]
CMS_t_tch_13TeV_2019_diff_Yt	35.9	t -channel	$d\sigma/d y_t $	5	[101]
ATLAS_t_sch_13TeV_inc	139	s -channel	$\sigma(t + \bar{t})$	1	[69]
ATLAS_tW_13TeV_inc	3.2	multi-lepton	$\sigma_{\text{tot}}(tW)$	1	[102]
CMS_tW_13TeV_inc	35.9	multi-lepton	$\sigma_{\text{tot}}(tW)$	1	[103]
CMS_tW_13TeV_slep_inc	36	single-lepton	$\sigma_{\text{tot}}(tW)$	1	[71]
ATLAS_tZ_13TeV_run2_inc	139	multi-lepton + jets	$\sigma_{\text{fid}}(t\ell^+\ell^-q)$	1	[104]
CMS_tZ_13TeV_pTt	138	multi-lepton + jets	$d\sigma_{\text{fid}}(tZj)/dp_T^t$	3	[70]

FCC-ee and CEPC datasets

Zh and VBF (*hνν*)

EWPOs

Z-pole EWPOs ($\sqrt{s} = 91.2$ GeV)		
\mathcal{O}_i	$\delta/\Delta \mathcal{O}_i$	
	FCC-ee	CEPC
$\alpha(m_Z)^{-1} (\times 10^3)$	$\Delta = 2.7$ (1.2)	$\Delta = 17.8$
Γ_W (MeV)	$\Delta = 0.85$ (0.3)	$\Delta = 1.8$ (0.9)
Γ_Z (MeV)	$\Delta = 0.0028$ (0.025)	$\Delta = 0.005$ (0.025)
$A_e (\times 10^5)$	$\Delta = 0.5$ (2)	$\Delta = 1.5$
$A_\mu (\times 10^5)$	$\Delta = 1.6$ (2.2)	$\Delta = 3.0$ (1.8)
$A_\tau (\times 10^5)$	$\Delta = 0.35$ (20)	$\Delta = 1.2$ (6.9)
$A_b (\times 10^5)$	$\Delta = 1.7$ (21)	$\Delta = 3$ (21)
$A_c (\times 10^5)$	$\Delta = 14$ (15)	$\Delta = 6$ (30)
σ_{had}^0 (pb)	$\Delta = 0.025$ (4)	$\Delta = 0.05$ (2)
$R_e (\times 10^3)$	$\delta = 0.0028$ (0.3)	$\delta = 0.003$ (0.2)
$R_\mu (\times 10^3)$	$\delta = 0.0021$ (0.05)	$\delta = 0.003$ (0.1)
$R_\tau (\times 10^3)$	$\delta = 0.0021$ (0.1)	$\delta = 0.003$ (0.1)
$R_b (\times 10^3)$	$\delta = 0.001$ (0.3)	$\delta = 0.005$ (0.2)
$R_c (\times 10^3)$	$\delta = 0.011$ (1.5)	$\delta = 0.02$ (1)

$e^+e^- \rightarrow Zh$				
	$\sqrt{s} = 240$ GeV		$\sqrt{s} = 365$ GeV	
\mathcal{O}_i	$\delta_{\text{exp}} \mathcal{O}_i$ (FCC-ee)	$\delta_{\text{exp}} \mathcal{O}_i$ (CEPC)	$\delta_{\text{exp}} \mathcal{O}_i$ (FCC-ee)	$\delta_{\text{exp}} \mathcal{O}_i$ (CEPC)
σ_{Zh}	0.0035	0.0026	0.0064	0.014
$\sigma_{Zh} \times \text{BR}_{b\bar{b}}$	0.0021	0.0014	0.0035	0.009
$\sigma_{Zh} \times \text{BR}_{c\bar{c}}$	0.0156	0.0202	0.046	0.088
$\sigma_{Zh} \times \text{BR}_{gg}$	0.0134	0.0081	0.0247	0.034
$\sigma_{Zh} \times \text{BR}_{ZZ}$	0.0311	0.0417	0.0849	0.2
$\sigma_{Zh} \times \text{BR}_{WW}$	0.0085	0.0053	0.0184	0.028
$\sigma_{Zh} \times \text{BR}_{\tau^+\tau^-}$	0.0064	0.0042	0.0127	0.021
$\sigma_{Zh} \times \text{BR}_{\gamma\gamma}$	0.0636	0.0302	0.127	0.11
$\sigma_{Zh} \times \text{BR}_{\gamma Z}$	0.12	0.085	-	-
$e^+e^- \rightarrow h\nu\nu$				
	$\sqrt{s} = 240$ GeV		$\sqrt{s} = 365$ GeV	
\mathcal{O}_i	$\delta_{\text{exp}} \mathcal{O}_i$ (FCC-ee)	$\delta_{\text{exp}} \mathcal{O}_i$ (CEPC)	$\delta_{\text{exp}} \mathcal{O}_i$ (FCC-ee)	$\delta_{\text{exp}} \mathcal{O}_i$ (CEPC)
$\sigma_{h\nu\nu} \times \text{BR}_{b\bar{b}}$	0.0219	0.0159	0.0064	0.011
$\sigma_{h\nu\nu} \times \text{BR}_{c\bar{c}}$	-	-	0.0707	0.16
$\sigma_{h\nu\nu} \times \text{BR}_{gg}$	-	-	0.0318	0.045
$\sigma_{h\nu\nu} \times \text{BR}_{ZZ}$	-	-	0.0707	0.21
$\sigma_{h\nu\nu} \times \text{BR}_{WW}$	-	-	0.0255	0.044
$\sigma_{h\nu\nu} \times \text{BR}_{\tau^+\tau^-}$	-	-	0.0566	0.042
$\sigma_{h\nu\nu} \times \text{BR}_{\gamma\gamma}$	-	-	0.156	0.16

FCC-ee and CEPC datasets

$e^+e^- \rightarrow f\bar{f}$					Light fermion production	
	$\sqrt{s} = 240 \text{ GeV}$		$\sqrt{s} = 365 \text{ GeV}$			
\mathcal{O}_i	$\Delta_{\text{exp}} \mathcal{O}_i$ (FCC-ee)	$\Delta_{\text{exp}} \mathcal{O}_i$ (CEPC)	$\Delta_{\text{exp}} \mathcal{O}_i$ (FCC-ee)	$\Delta_{\text{exp}} \mathcal{O}_i$ (CEPC)		
$\sigma_{\text{tot}}(e^+e^-)$ [fb]	2.29	1.62	2.74	4.68		
$A_{\text{FB}}(e^+e^-)$	$9.79 \cdot 10^{-6}$	$6.92 \cdot 10^{-6}$	$2.83 \cdot 10^{-5}$	$4.83 \cdot 10^{-5}$		
$\sigma_{\text{tot}}(\mu^+\mu^-)$ [fb]	0.405	0.287	0.48	0.82		
$A_{\text{FB}}(\mu^+\mu^-)$	$1.98 \cdot 10^{-4}$	$1.397 \cdot 10^{-4}$	$5.69 \cdot 10^{-4}$	$9.7 \cdot 10^{-4}$		
$\sigma_{\text{tot}}(\tau^+\tau^-)$ [fb]	0.374	0.264	0.443	0.756		
$A_{\text{FB}}(\tau^+\tau^-)$	$2.17 \cdot 10^{-4}$	$1.53 \cdot 10^{-4}$	$6.24 \cdot 10^{-4}$	0.00106		
$\sigma_{\text{tot}}(c\bar{c})$ [fb]	0.088	0.062	0.102	0.175		
$A_{\text{FB}}(c\bar{c})$	0.000813	$5.74 \cdot 10^{-4}$	0.00238	0.00405		
$\sigma_{\text{tot}}(b\bar{b})$ [fb]	0.151	0.107	0.171	0.29		
$A_{\text{FB}}(b\bar{b})$	$4.86 \cdot 10^{-4}$	$3.44 \cdot 10^{-4}$	0.00142	0.00243		

$e^+e^- \rightarrow W^+W^-$						
\mathcal{O}_i	$\sqrt{s} = 161 \text{ GeV}$		$\sqrt{s} = 240 \text{ GeV}$		$\sqrt{s} = 365 \text{ GeV}$	
	δ_{exp} (FCC-ee)	δ_{exp} (CEPC)	δ_{exp} (FCC-ee)	δ_{exp} (CEPC)	δ_{exp} (FCC-ee)	δ_{exp} (CEPC)
σ_{WW}	$1.36 \cdot 10^{-4}$	$2.48 \cdot 10^{-4}$	$1.22 \cdot 10^{-4}$	$8.63 \cdot 10^{-5}$	$2.81 \cdot 10^{-4}$	$4.87 \cdot 10^{-4}$
$\text{BR}_{W \rightarrow \ell_i \nu_i}$	$2.72 \cdot 10^{-4}$	$4.95 \cdot 10^{-4}$	$2.44 \cdot 10^{-4}$	$1.73 \cdot 10^{-4}$	$5.63 \cdot 10^{-4}$	$9.75 \cdot 10^{-4}$

HL-LHC projections

- ▶ The central values of the pseudo data are fluctuated around the SM

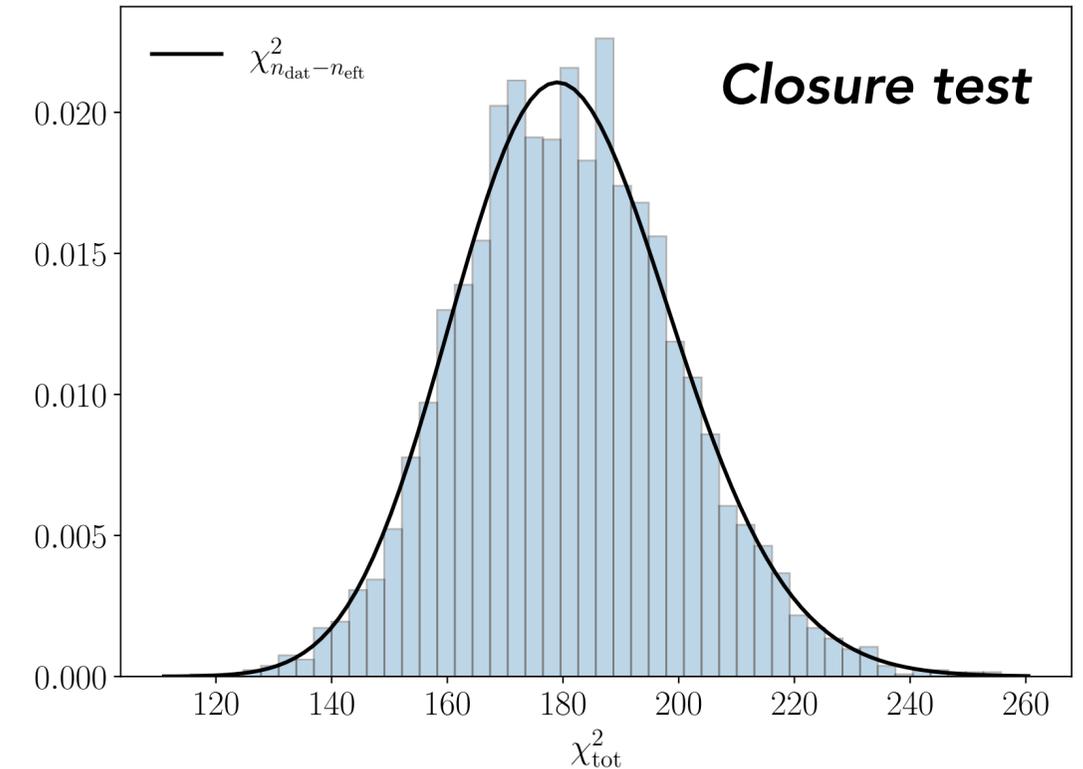
$$\mathcal{O}_i^{(\text{exp})} = \mathcal{O}_i^{(\text{th})} \left(1 + r_i \delta_i^{(\text{stat})} + \sum_{k=1}^{n_{\text{sys}}} r_{k,i} \delta_{k,i}^{(\text{sys})} \right)$$

- ▶ Statistical uncertainties we rescale according to the improved luminosity

$$\delta_i^{(\text{stat})} = \tilde{\delta}_i^{(\text{stat})} \sqrt{\frac{\mathcal{L}_{\text{Run2}}}{\mathcal{L}_{\text{HLLHC}}}}$$

$$\delta_{k,i}^{(\text{sys})} = \tilde{\delta}_{k,i}^{(\text{sys})} \times f_{\text{red}}^{(k)} \quad k = 1, \dots, n_{\text{sys}}$$

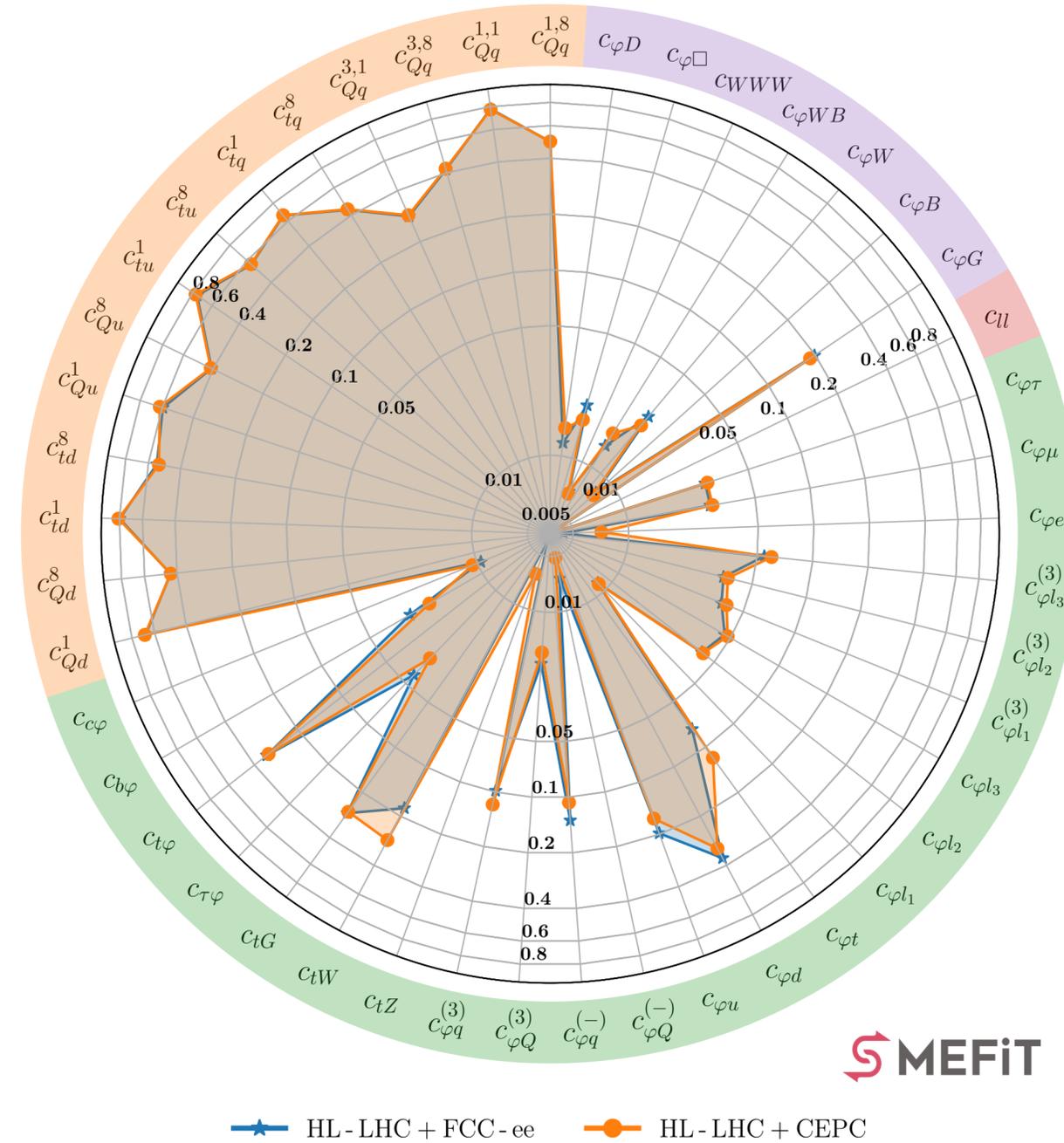
- ▶ While systematics are rescaled by an overall factor, namely 1/2 for all datasets



- + flexible framework that can project any Run II dataset
- + SMEFT predictions can be recycled
- No additional bins in the tails

FCC-ee and CEPC

Ratio of Uncertainties to SMEFiT3.0 Baseline, $\mathcal{O}(\Lambda^{-2})$, Marginalised



1-loop & multi-particle matching

